Cosmological Stasis and its Realization from Dynamical Scalars



Based on work done in collaboration with:

Keith Dienes, Fei Huang, Lucien Heurtier, and Tim Tait [arXiv:2406.06830]

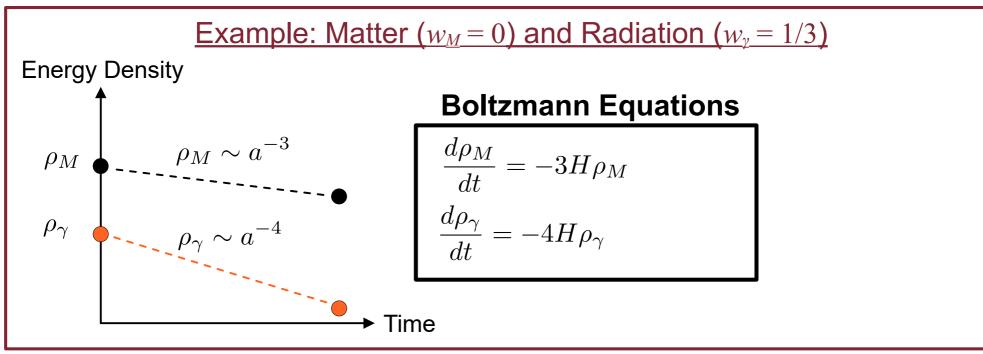
Keith Dienes, Fei Huang, Lucien Heurtier, Doojin Kim and Tim Tait [arXiv:2111.04753] Keith Dienes, Fei Huang, Lucien Heurtier, Doojin Kim and Tim Tait [arXiv:2212.01369] Keith Dienes, Fei Huang, Lucien Heurtier, and Tim Tait [arXiv:2406.06830]

CETUP* Workshop, Lead, South Dakota, June 27th, 2024

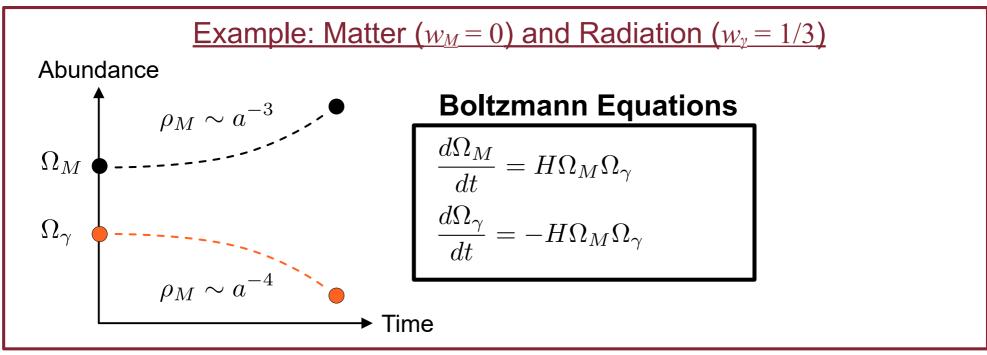
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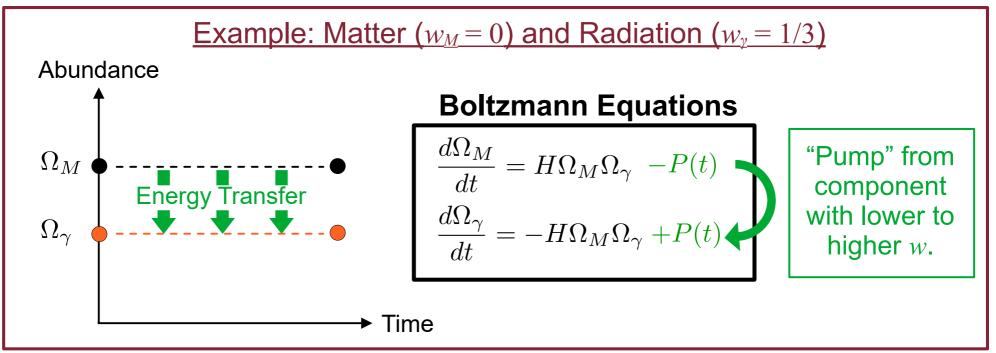
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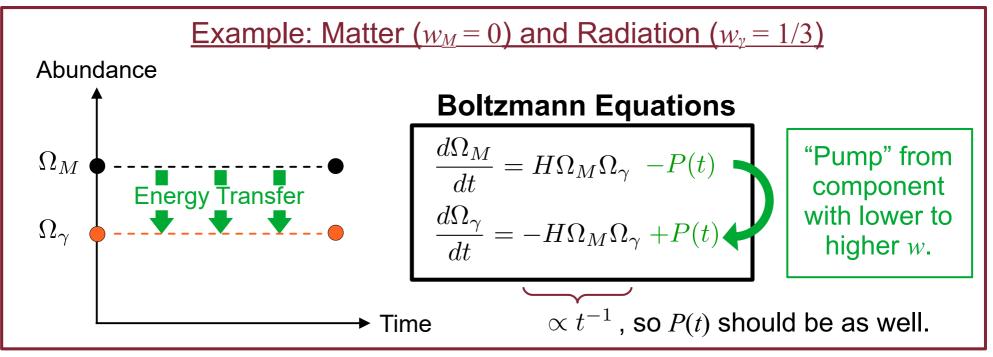
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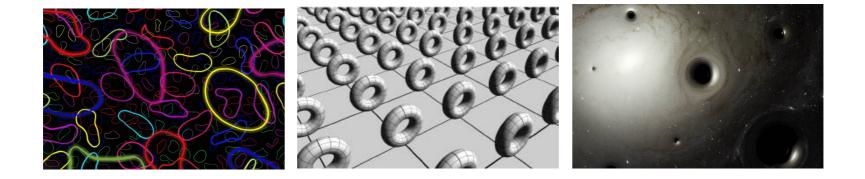


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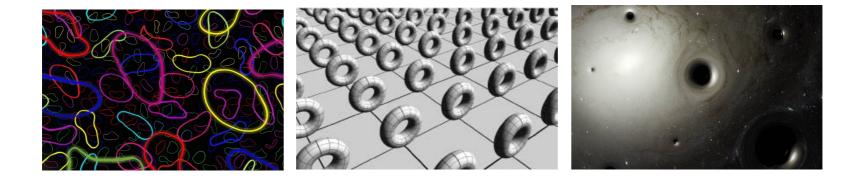
- The exponential decay of a single matter species, which occurs over a relatively short time period, is insufficient for achieving stasis.
- However, a <u>tower of matter states</u> ϕ_{ℓ} , where $\ell = 0, 1, 2, ..., N 1$, whose decay widths Γ_{ℓ} and initial abundances $\Omega_{\ell}^{(0)}$ scale across the tower as a function of their mass m_{ℓ} can indeed give rise to a pump that compensate for the effect of cosmic expansion over a extended period.

 Pump terms with the right time-dependence for stasis emerge naturally in scenarios involving <u>towers of states</u> with broad spectra of masses, cosmological abundances, lifetimes, etc.

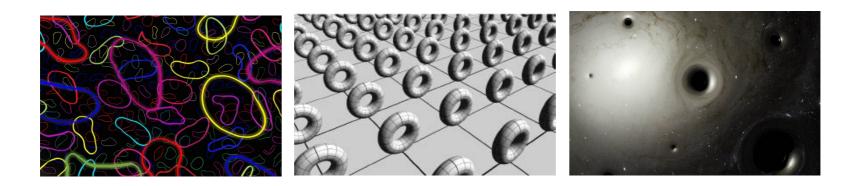
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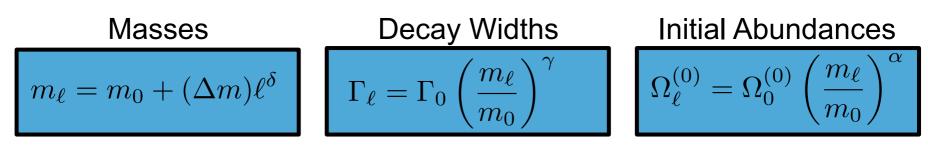


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- The modified cosmological histories associated with stasis can affect the evolution of **scalar and tensor perturbations**.



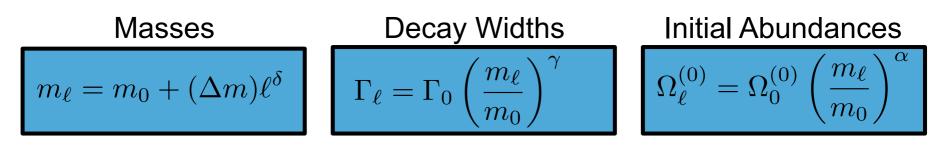
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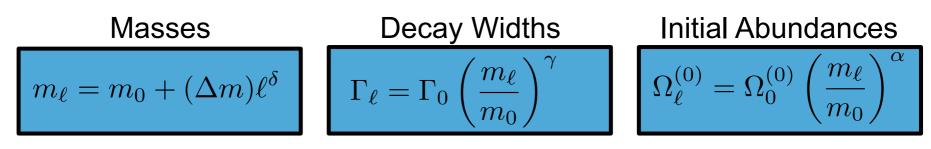
- Towers of states with mass spectra of this form arise naturally in many extensions of the Standard Model.
 - KK excitations of a 5D scalar:
 - Bound states of a stronglycoupled gauge theory:

$$\begin{cases} mR \ll 1 & \longrightarrow & \delta \sim 1 \\ mR \gg 1 & \longrightarrow & \delta \sim 2 \end{cases}$$

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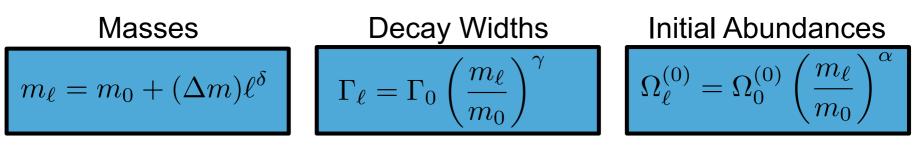
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- Scaling of initial abundances depends on how they're generated:

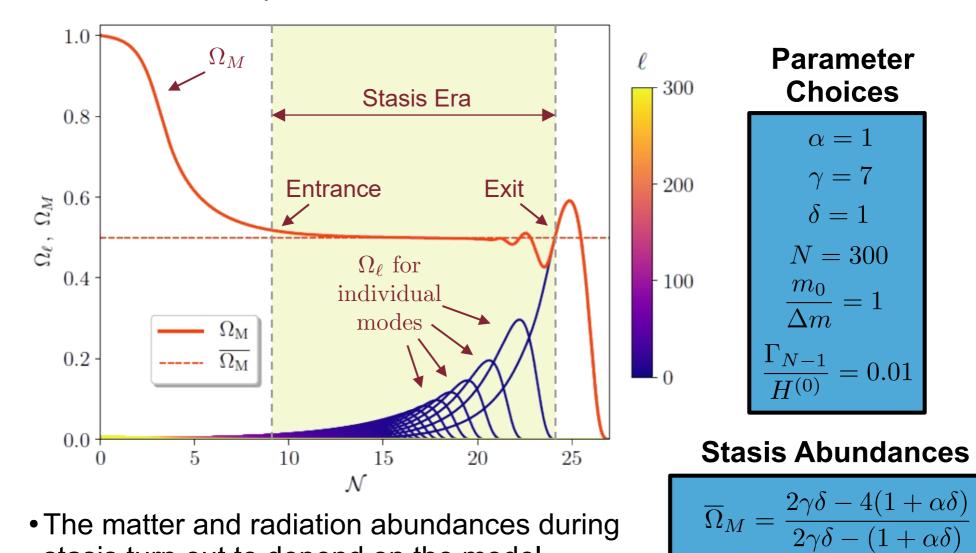
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The Emergence of Stasis

 In BSM scenarios of this sort, stasis emerges generically, with minimal additional assumptions.

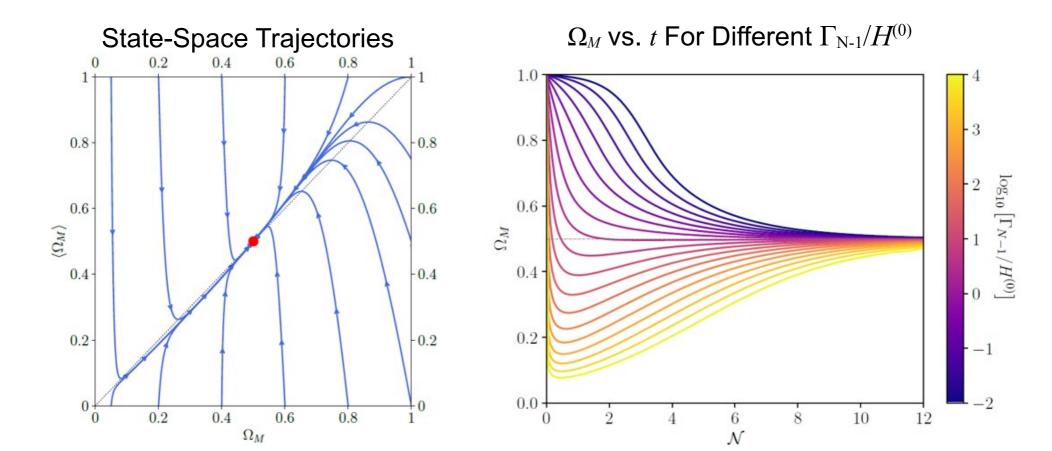


 $\overline{\Omega}_{\gamma} = 1 - \overline{\Omega}_M$

• The matter and radiation abundances during stasis turn out to depend on the model parameters α , γ , and δ .

Stasis as a Global Attractor

- Perhaps even more importantly, achieving cosmological stasis does not require a fine-tuning of the initial conditions for Ω_M and H – or, alternatively, for Ω_m and its time-average $\langle \Omega_M \rangle$ – or for the ratio $\Gamma_{N-1}/H^{(0)}$.
- In fact, stasis is a **global attractor** in the sense that Ω_M and Ω_γ will **evolve toward their stasis values** regardless of what these initial conditions are.



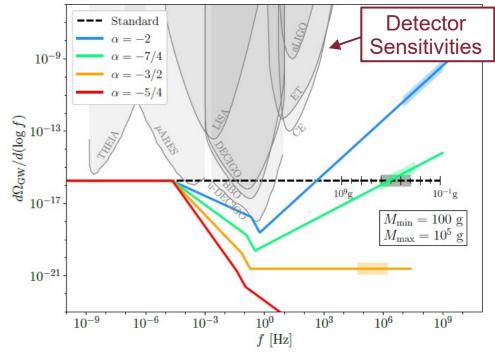
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- <u>Gravitational waves</u>: Such modifications of the cosmological timeline also alter the contribution to the graviational-wave background generated by other sources.

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GW Spectra in Primordial-Black-Hole-Induced Stasis



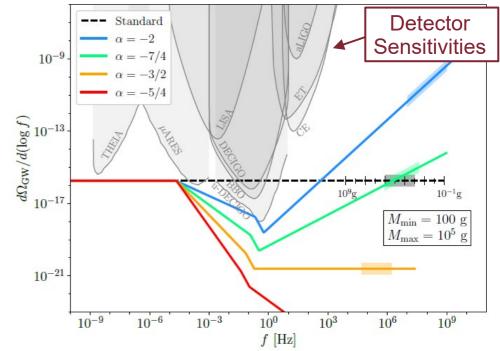
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• **Density perturbations**: Such perturbations evolve differently than they do in the standard cosmology, since $w_{\text{eff}} \neq 1/3$ during stasis. Possible implications for small-scale structure.

[Dienes, Huang, Heurtier, Hoover, Paulsen, Tait, BT '24]

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Other Kinds of Stasis

 Scenarios for BSM physics can give rise to stases involving <u>other</u> <u>cosmological components</u> beyond matter and radiation.

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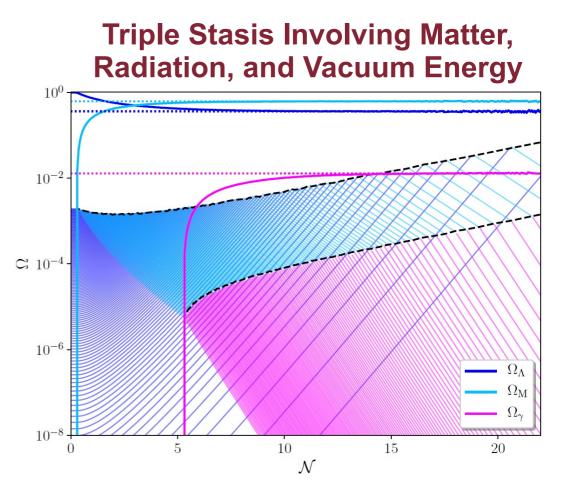
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- A stasis can also arise involving <u>more than two components</u> – for example, vacuum energy, matter, and radiation – can also arise in BSM contexts. Such stases also turn out to be global attractors. [Dienes, Huang, Heurtier, Tait, BT '23]





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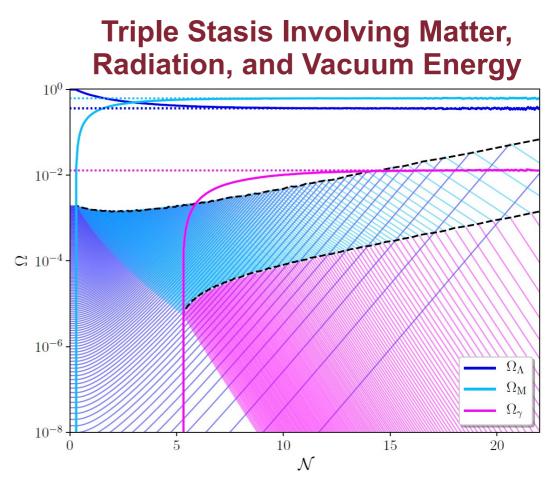
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Lucien Heurtier

Fei Huang



• A set of **axion-like scalar fields** whose background values transition from overdamped to underdamped oscillation can give rise to stasis... and *that's* what I'll be focusing on for the remainder of this talk!

[Dienes, Huang, Heurtier, Tait, BT '24]

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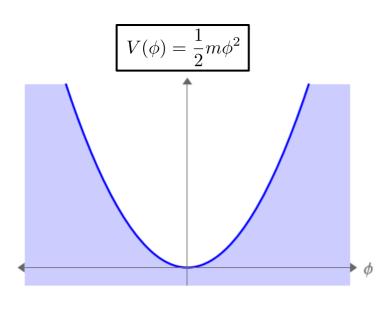
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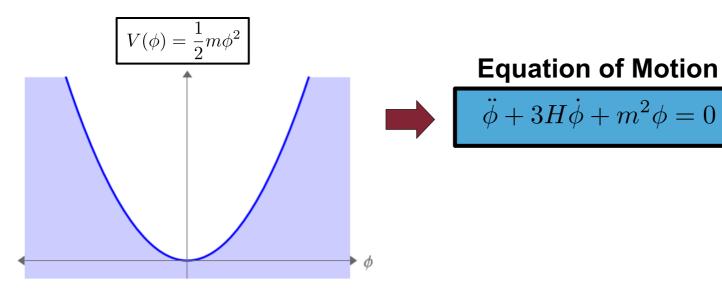
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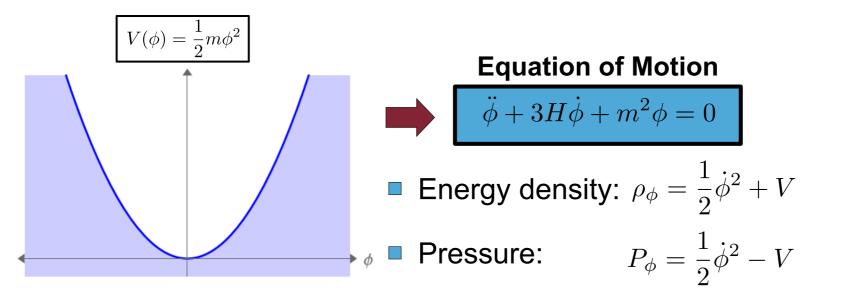
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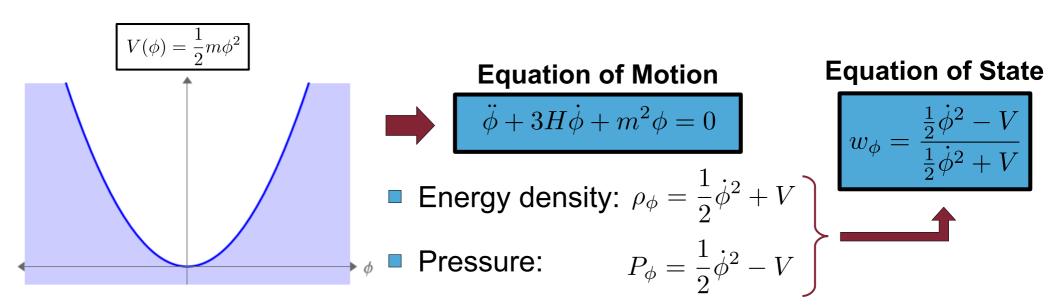
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- Such a stasis, as we'll see, would be characterized by an effective equation-of-state parameter <u>between that of vacuum energy</u> ($w_{\Lambda} = -1$) and matter ($w_{M} = 0$).
- Moreover, stases involving dynamical scalars give rise to some phenomena not seen in other realizations of stasis which could potentially useful for addressing fundamenal questions in cosmology.

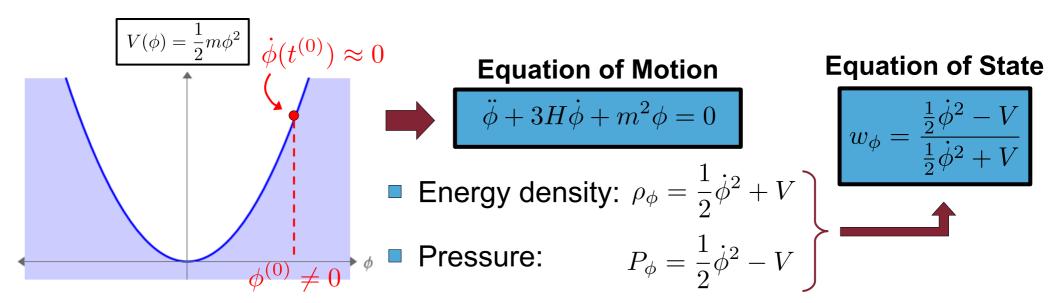




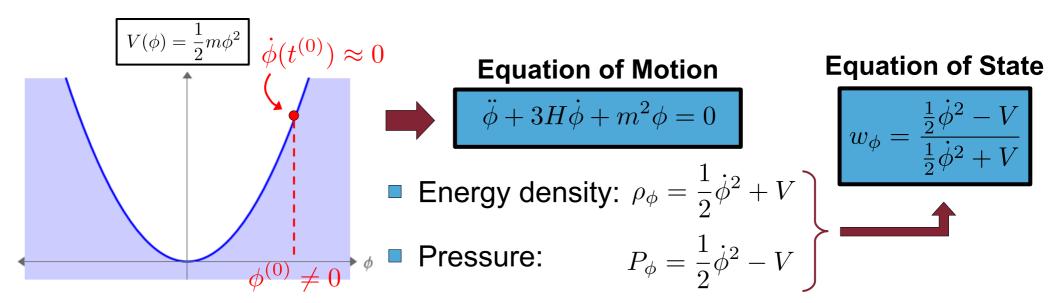




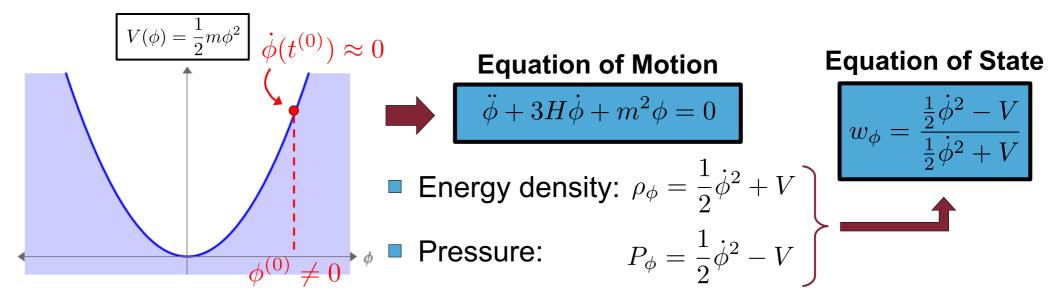
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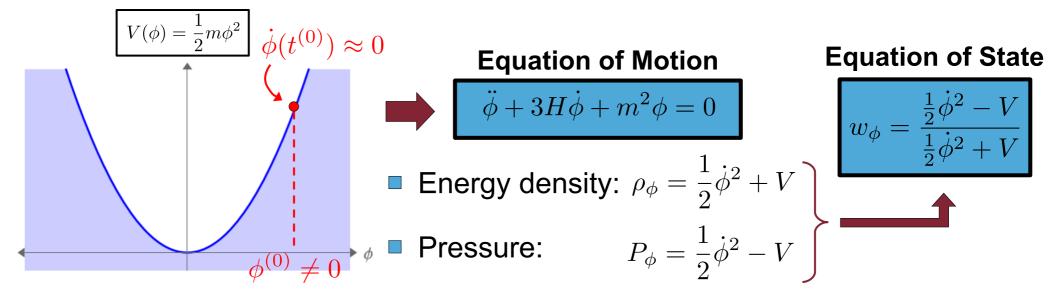


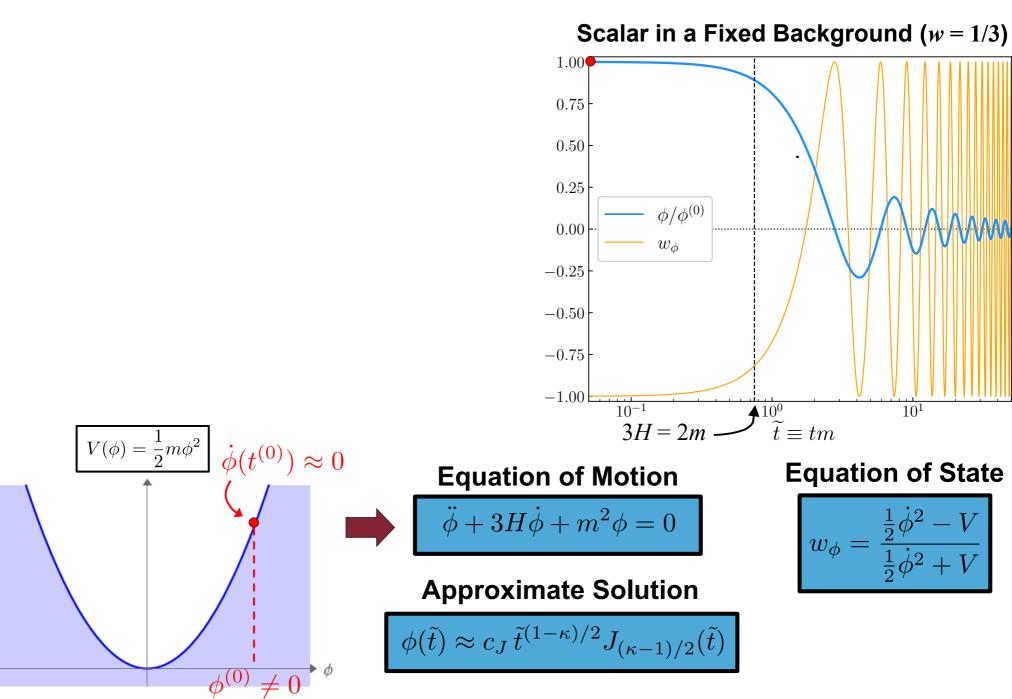
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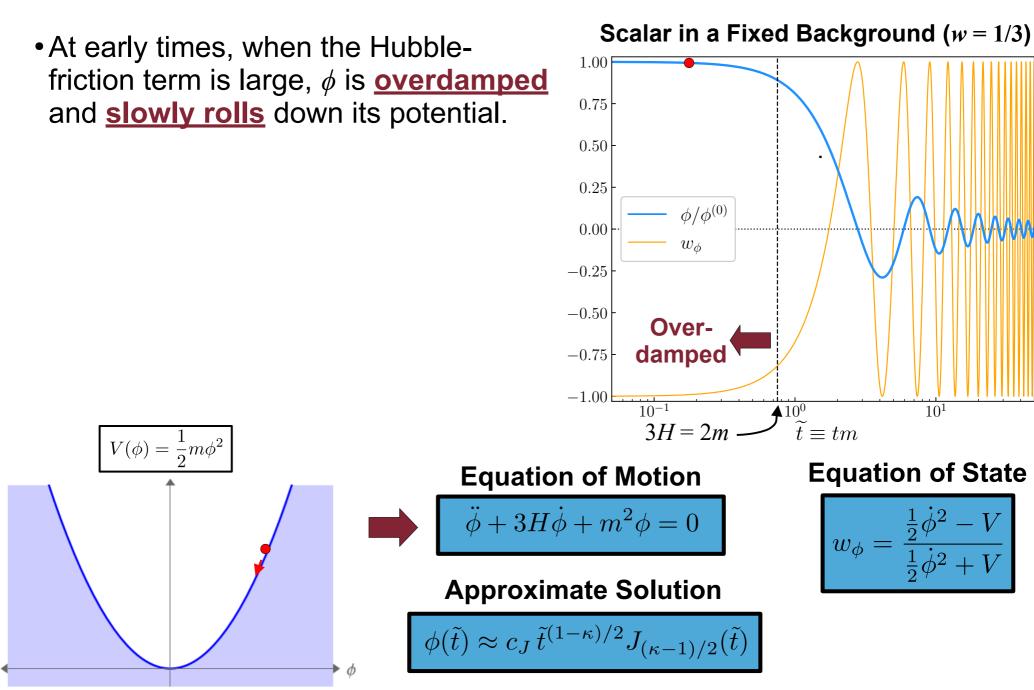


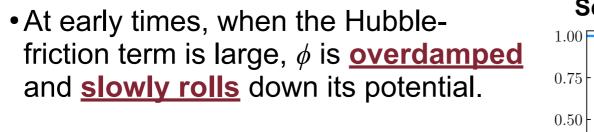
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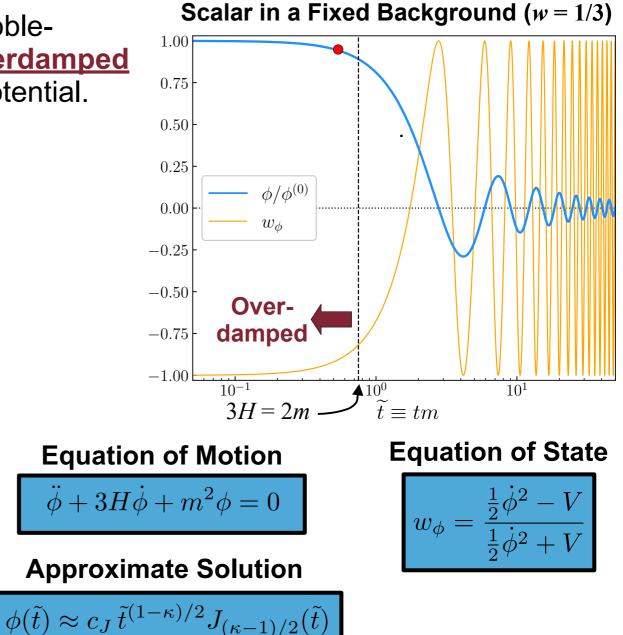
$$H \approx rac{\kappa}{3t}$$
 , where $\kappa \equiv rac{2}{1+w}$

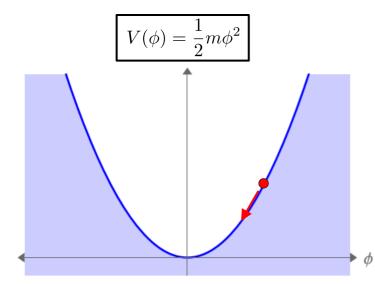


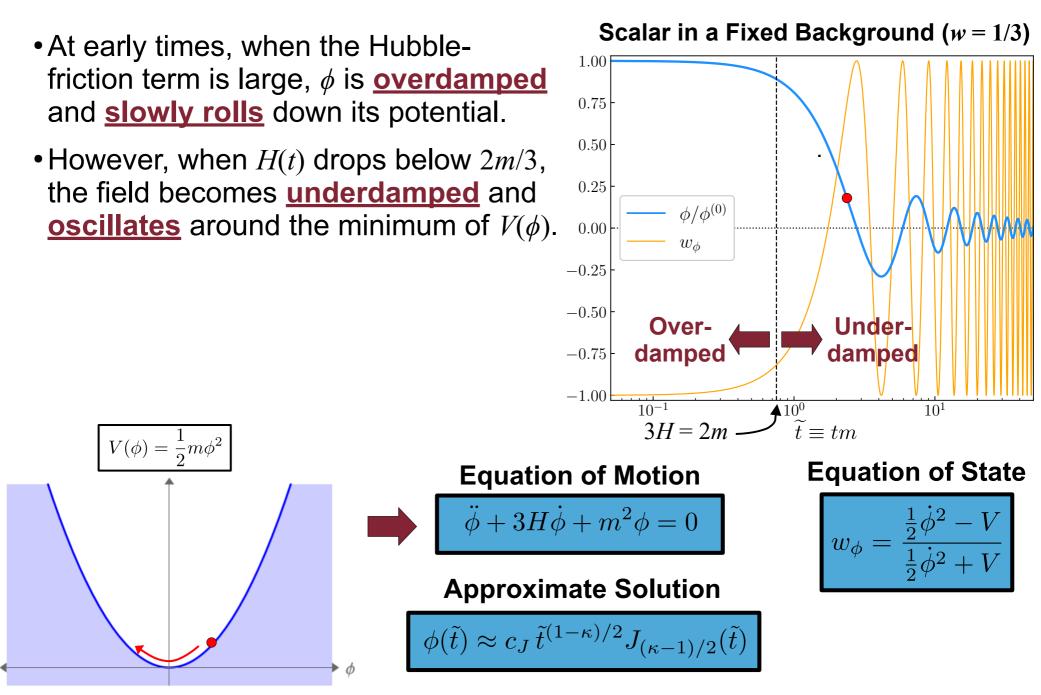


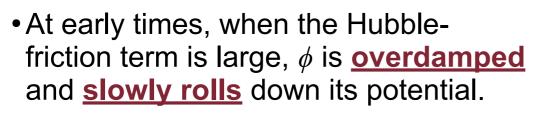




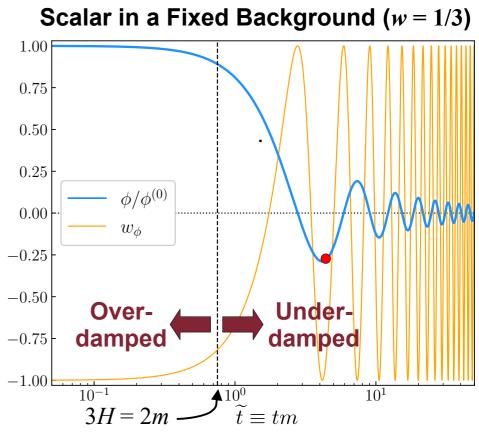


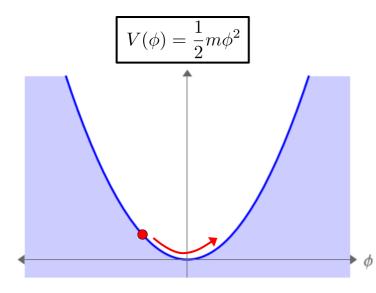


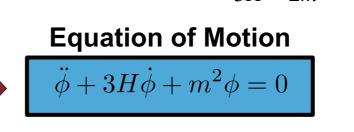




- However, when *H*(*t*) drops below 2*m*/3, the field becomes <u>underdamped</u> and <u>oscillates</u> around the minimum of *V*(φ).
- $w_{\phi}(t)$ oscillates rapidly at late times, but averages to $\langle w \rangle_t \approx 0$ over sufficiently long timescales.



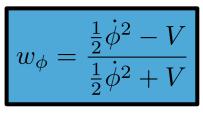


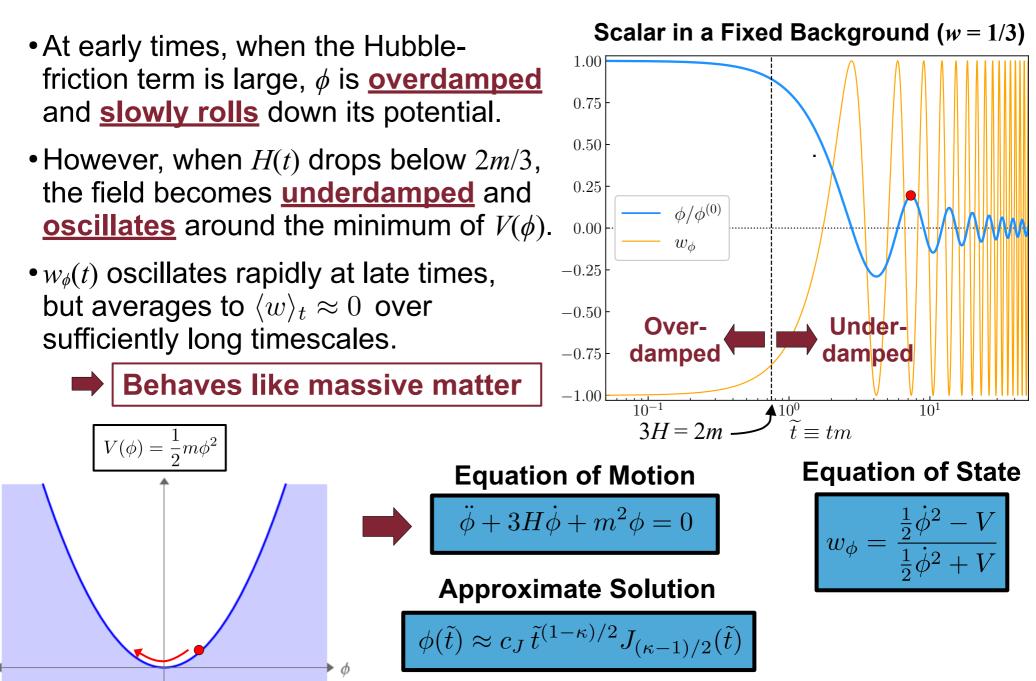


Approximate Solution

$$\phi(\tilde{t}) \approx c_J \, \tilde{t}^{(1-\kappa)/2} J_{(\kappa-1)/2}(\tilde{t})$$

Equation of State

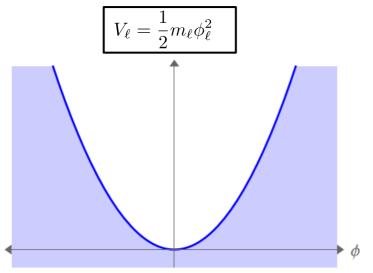




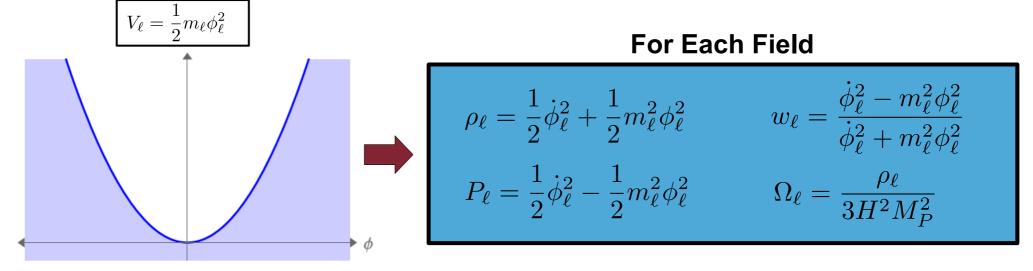
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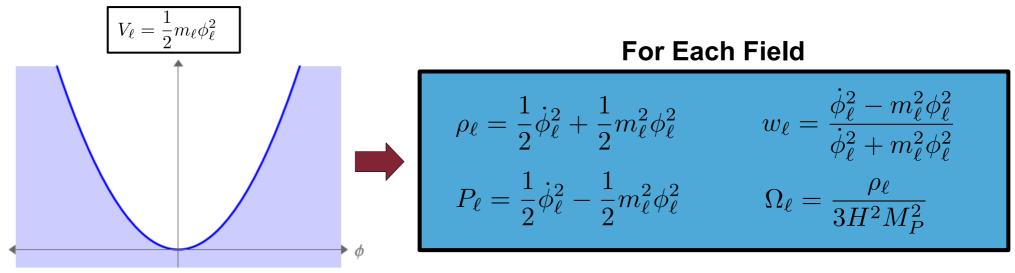
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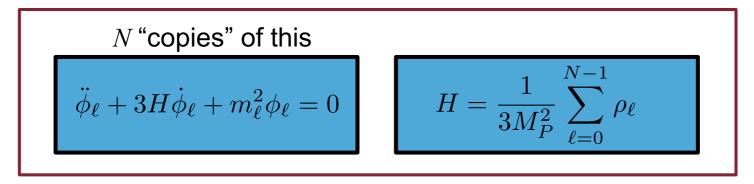


 We'll also assume (for the moment) that there's <u>no background energy</u> **<u>component</u>**: the collective energy density of the ϕ_{ℓ} dominates the universe.

- Now let's consider the case in which the universe comprises a <u>tower</u> of N such scalars ϕ_{ℓ} , where the index $\ell = 0, 1, 2, ..., N-1$ labels these states in order of increasing mass.
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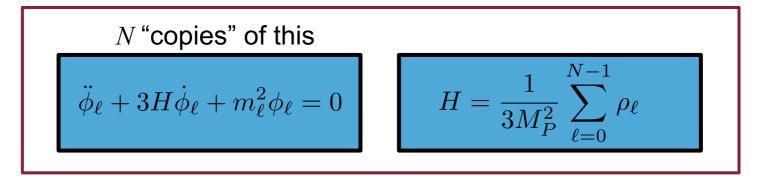
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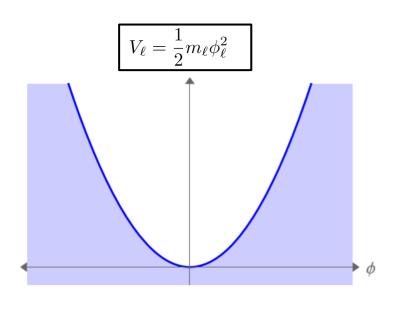
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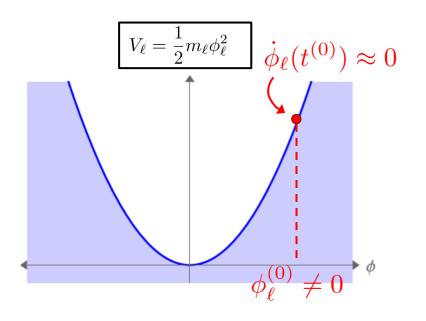


Let's see what the cosmology of such a tower of scalar-field zero modes looks like!

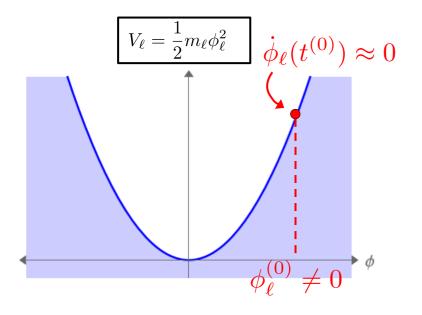
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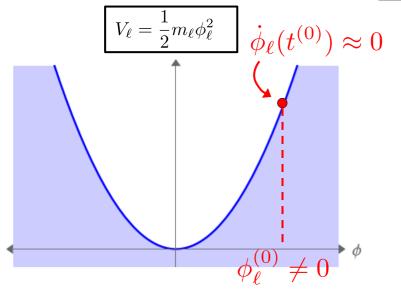


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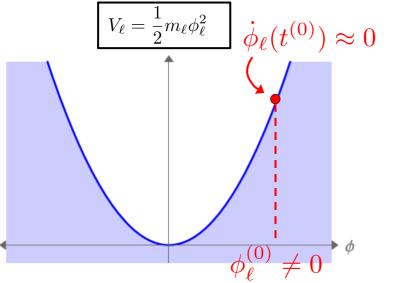
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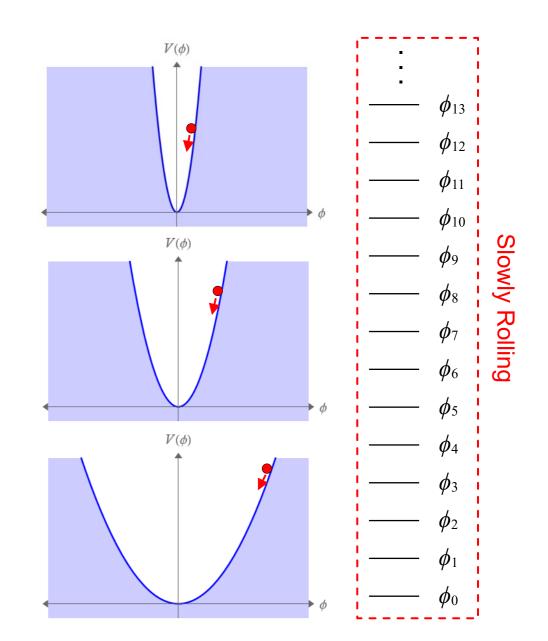
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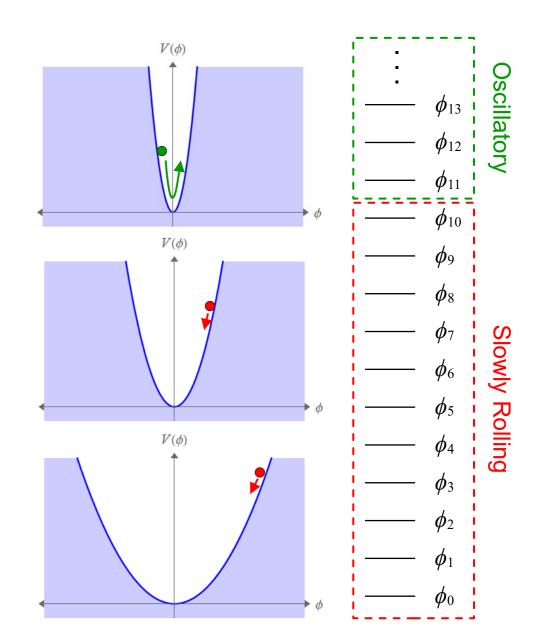
• For a given mass spectrum, the overall scale of the abundances can be parameterized by the ratio $\phi_0^{(0)}/M_P$, or, equivalently, by the ratio $H^{(0)}/m_{N-1}$.

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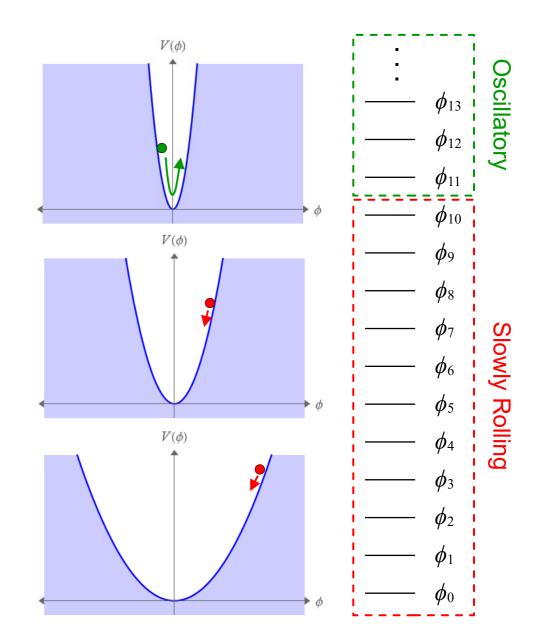
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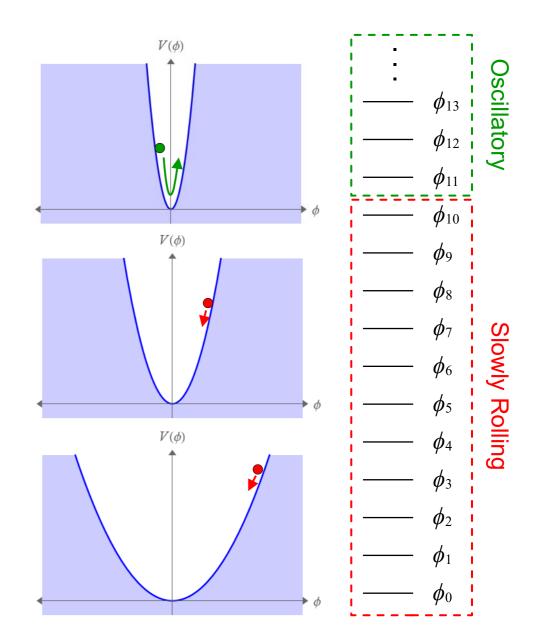
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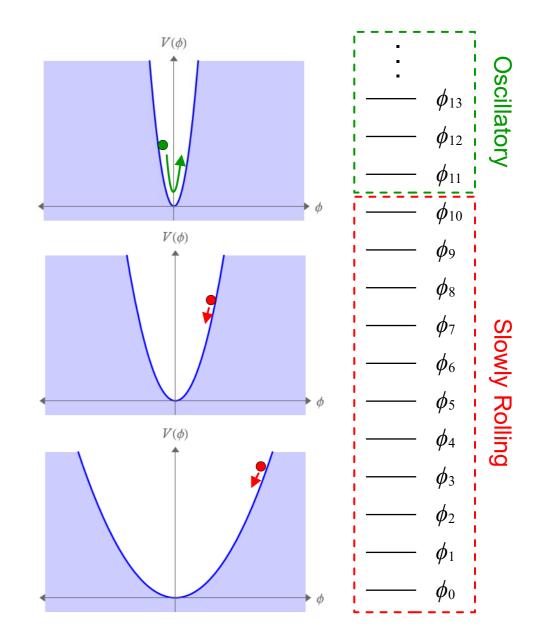


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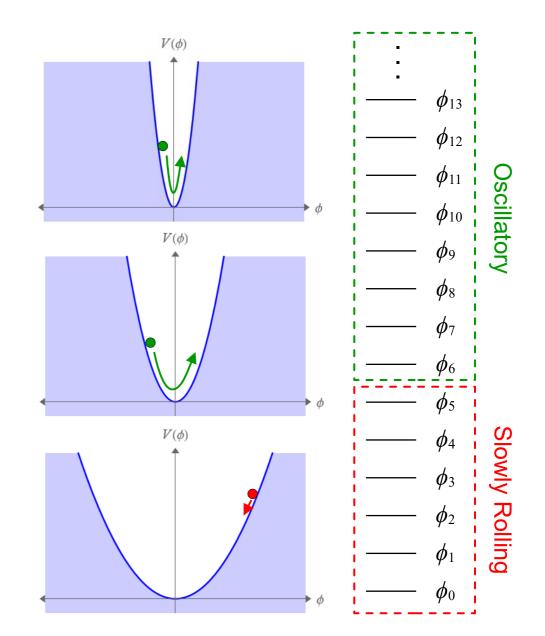


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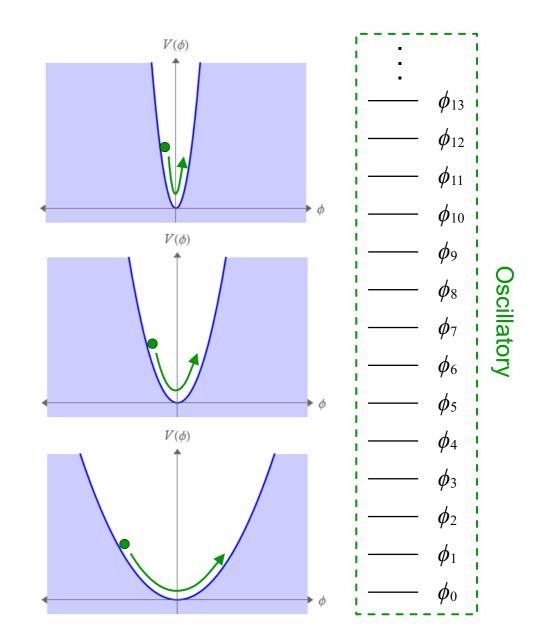


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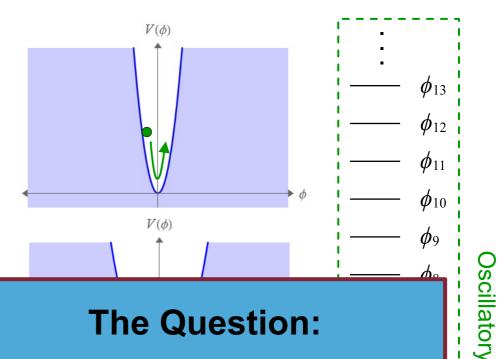
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The Question:

Can we achieve a stasis between these slow-roll and oscillatory cosmological energy components, which act like vacuum energy and matter, respectively?

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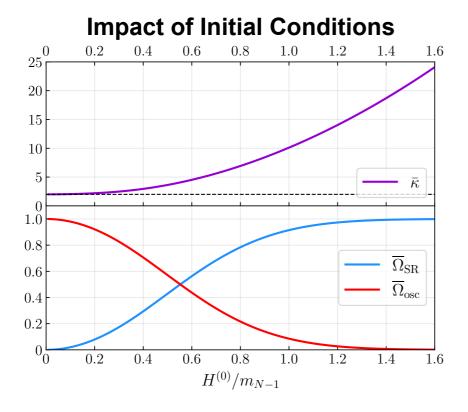
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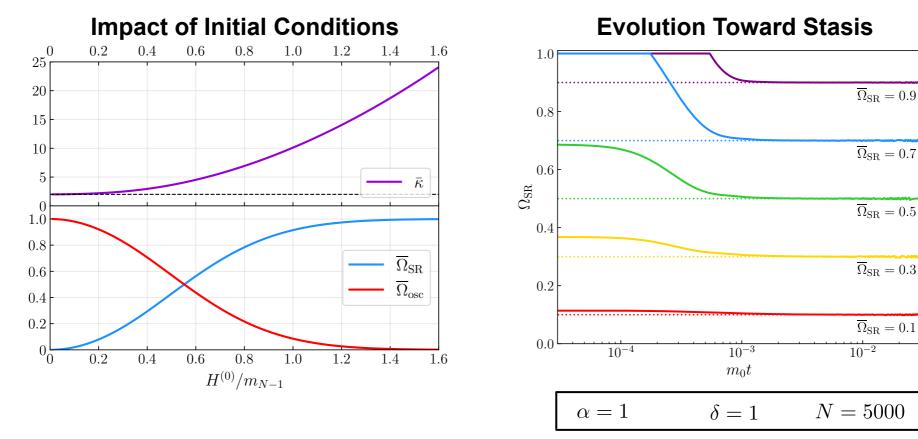
Towers which satisfy this relation give rise to stasis. For any δ , this corresponds to $\phi_\ell^{(0)} \sim \ell^{-1/2}$

• Unlike in previous realizations of stasis, the stasis abundances $\overline{\Omega}_{sR}$ and $\overline{\Omega}_{osc}$ depend on the **initial conditions** for the scalar tower.

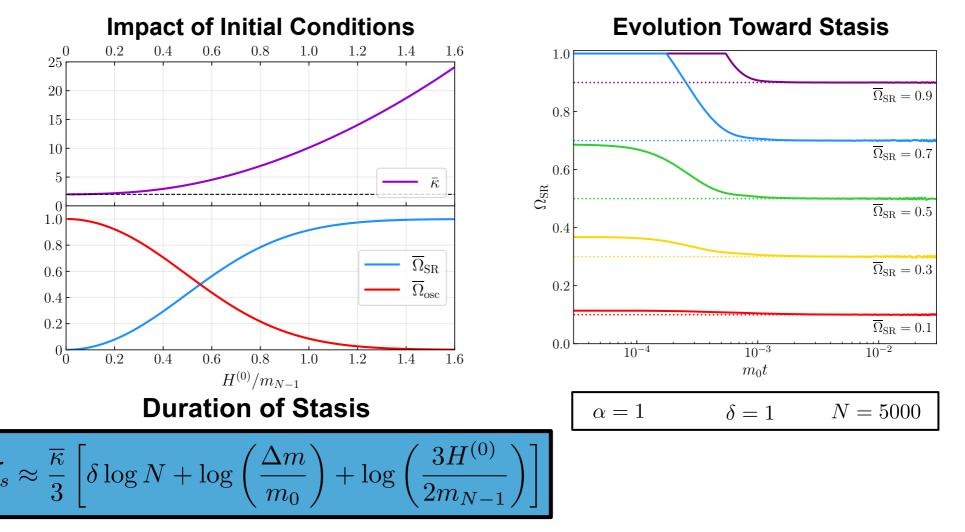
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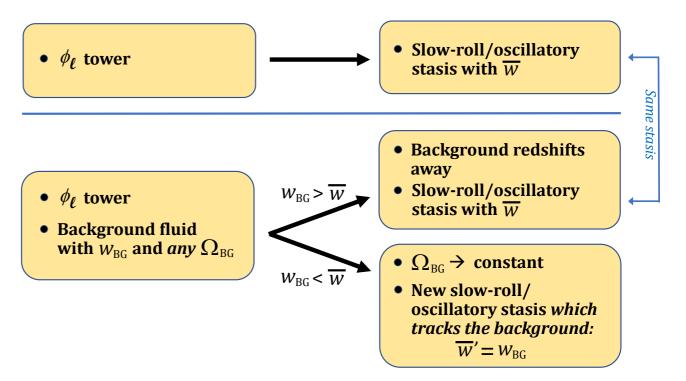


• Let's now consider how the cosmological dynamics is modified if we include a **<u>background energy component</u>** with a constant equation-of-state parameter w_{BG} in addition to the tower.

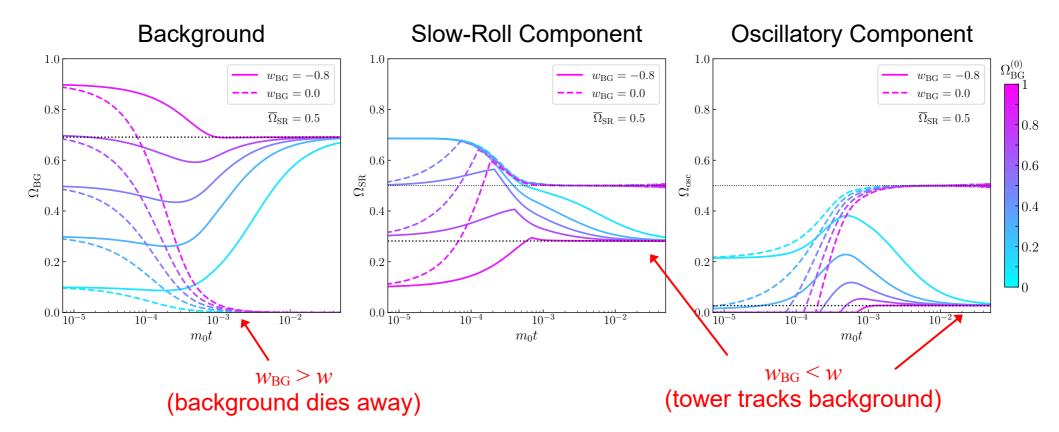
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- It turns out that in the presence of such an energy component, the universe still evolves toward stasis (or something like it).
- However, the outcome depends on the relationship between w_{BG} and the equation-of-state parameter \overline{w} the tower *would* have had during stasis if the background component weren't present.



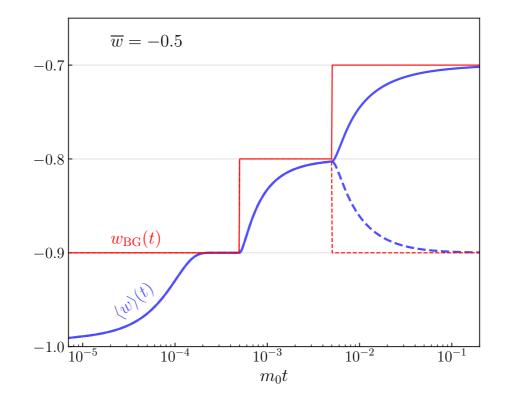


• The <u>tracking phenomenon</u> which arises in situations in which $w_{BG} < \overline{w}$ has not been observed in other realizations of stasis.

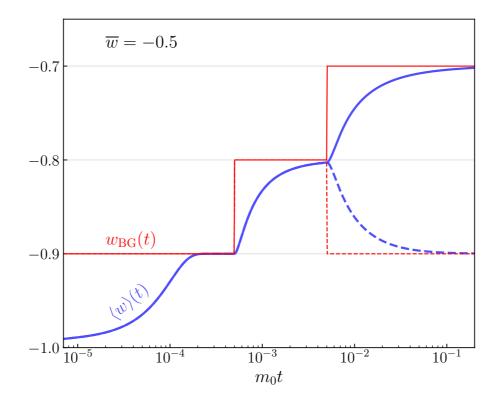


 These results provide insight about how the universe might <u>enter into</u> – or exit from – an stasis epoch involving dynamical scalars.

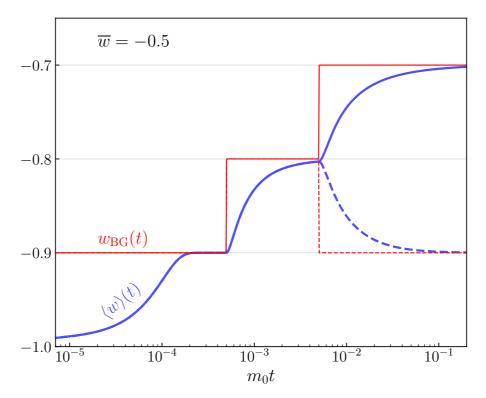
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• Moreover, $\langle w \rangle$ tracks w_{BG} even in the regime in which w_{BG} evolves continuously with time.

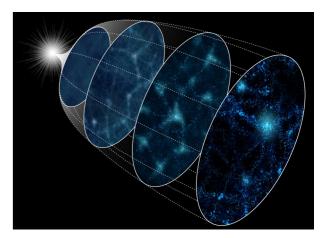
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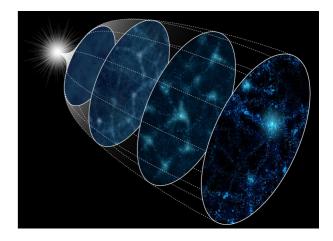
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Can such a stasis furnish a framework for cosmic inflation?

• This is an intriguing possibility for a number of reasons.

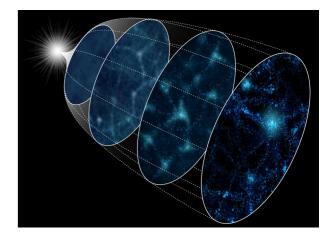


- Since values of \overline{w} within the range $-1 \le \overline{w} \le -1/3$ can be achieved during a stasis epoch involving dynamical scalars, the universe can undergo **accelerated expansion** during stasis.
- Since such a stasis epoch can endure for well over $N_e \sim 60$ *e*-folds of expansion, it can in principle solve the <u>horizon and flatness problems</u>.

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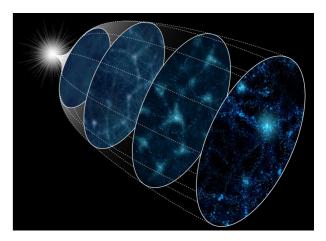
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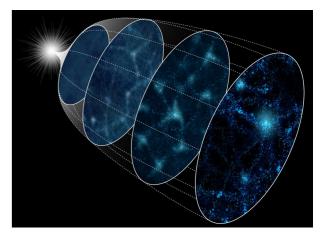
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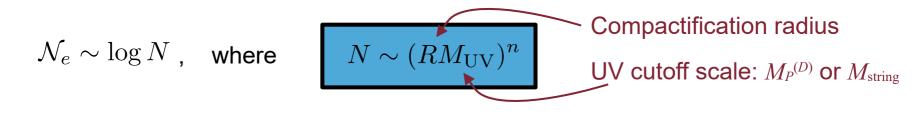


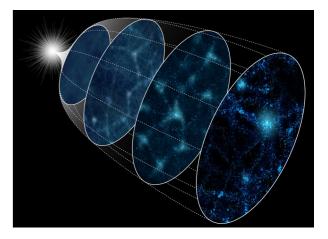
A "graceful exit" from inflation is built into this scenario. It ends with the ϕ_{ℓ} behaving like massive matter. Reheating can be achieved in principle via their subsequent decays.





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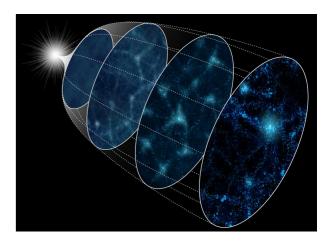


Compactification radius

UV cutoff scale: $M_P^{(D)}$ or M_{string}



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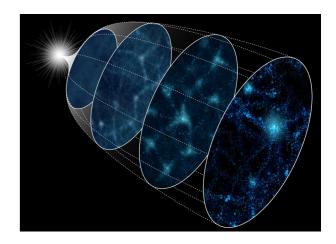
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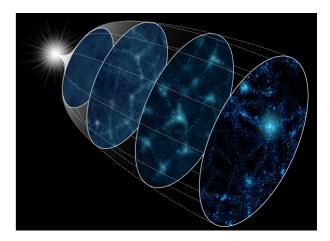
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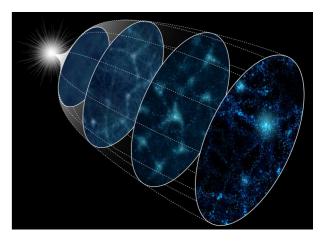
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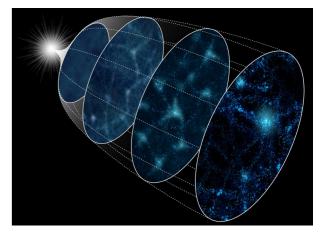
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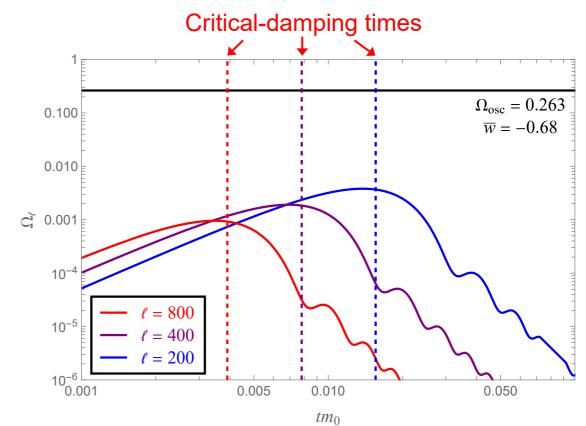
This is an intriguing possibility – and one that warrants further exploration!



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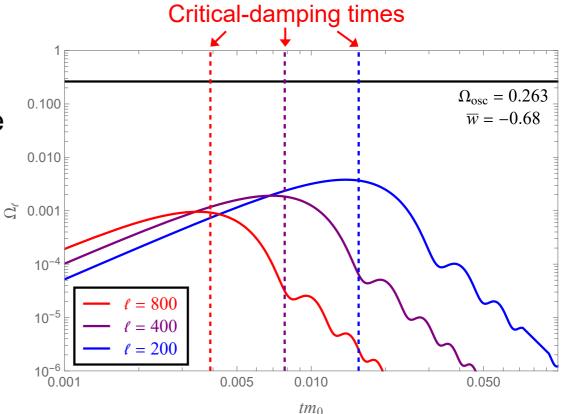
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- The abundance Ω_{ℓ} of each ϕ_{ℓ} initially rises while it's slowly rolling, but then **falls precipitiously** once its critical-damping time is reached.
- In this way, stasis inflation features an "<u>undertow</u>" which suppresses the abundances of unwanted relics and isocurvature contributions from the heavier φ_ℓ.





Summary

- <u>Stable, mixed-component cosmological eras</u> *i.e.* <u>stasis eras</u> are indeed a viable cosmological possibility – and one that can arise naturally in many extensions of the Standard Model.
- A <u>tower of scalar fields</u> which undergo a transition from overdamped to underdamped evolution can give rise to stasis.
- Stasis itself is an <u>attractor</u> in these systems, but several fundamental characteristics of the stasis epoch toward which the universe evolves depend on the initial conditions.
- In the presence of an additional background component with equationof-state parameter w_{BG} , the tower exhibits a <u>tracking behavior</u> in which its own equation-of-state parameter evolves toward w_{BG} .

