

Probing Neutrinophilic Dark Matter: From Colliders to Supernovae

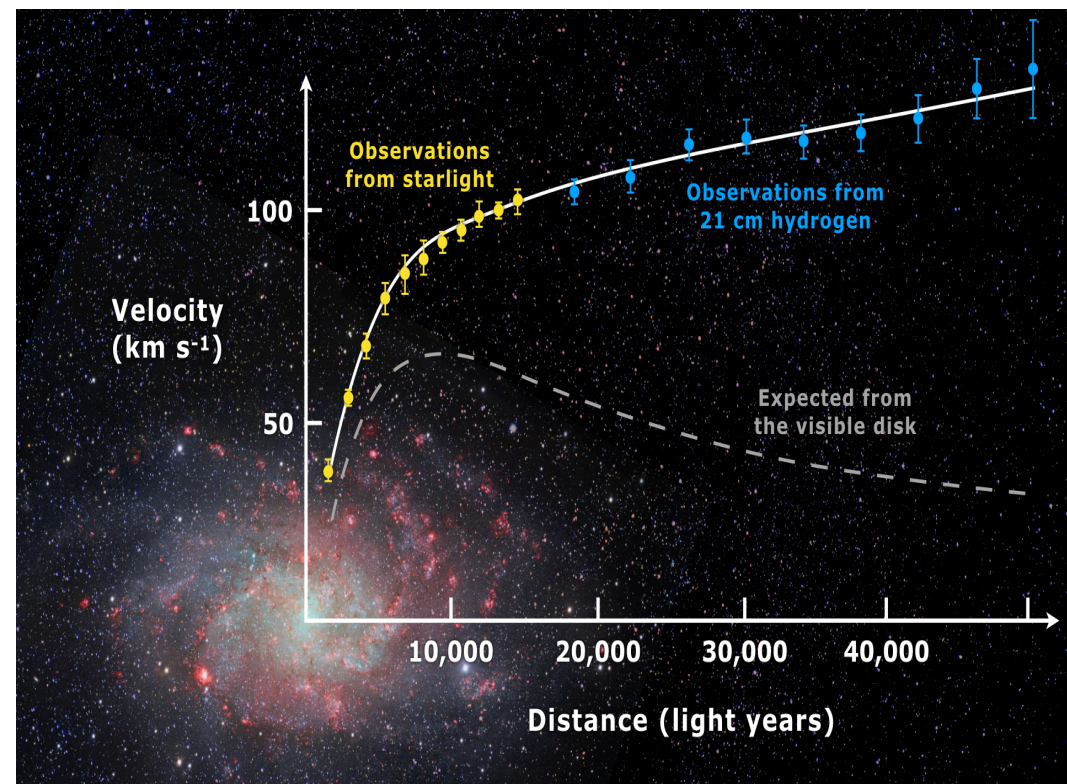
Douglas Tuckler

TRIUMF & Simon Fraser University

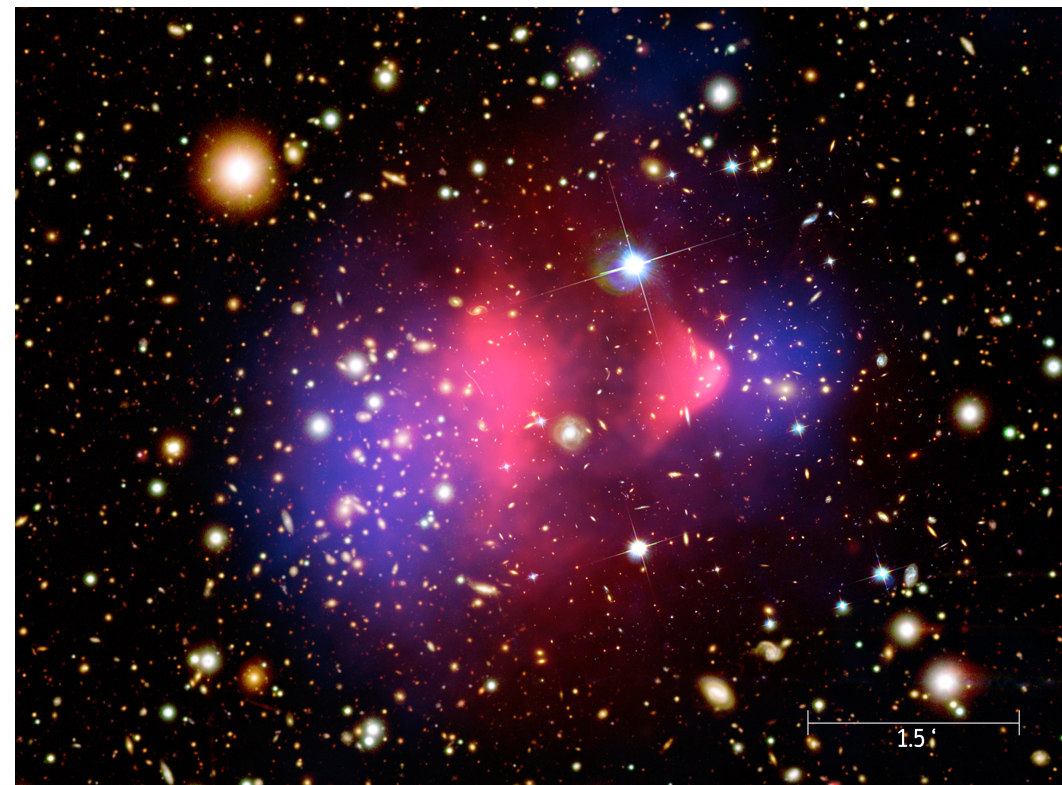
CETUP* 2024

June 21, 2024

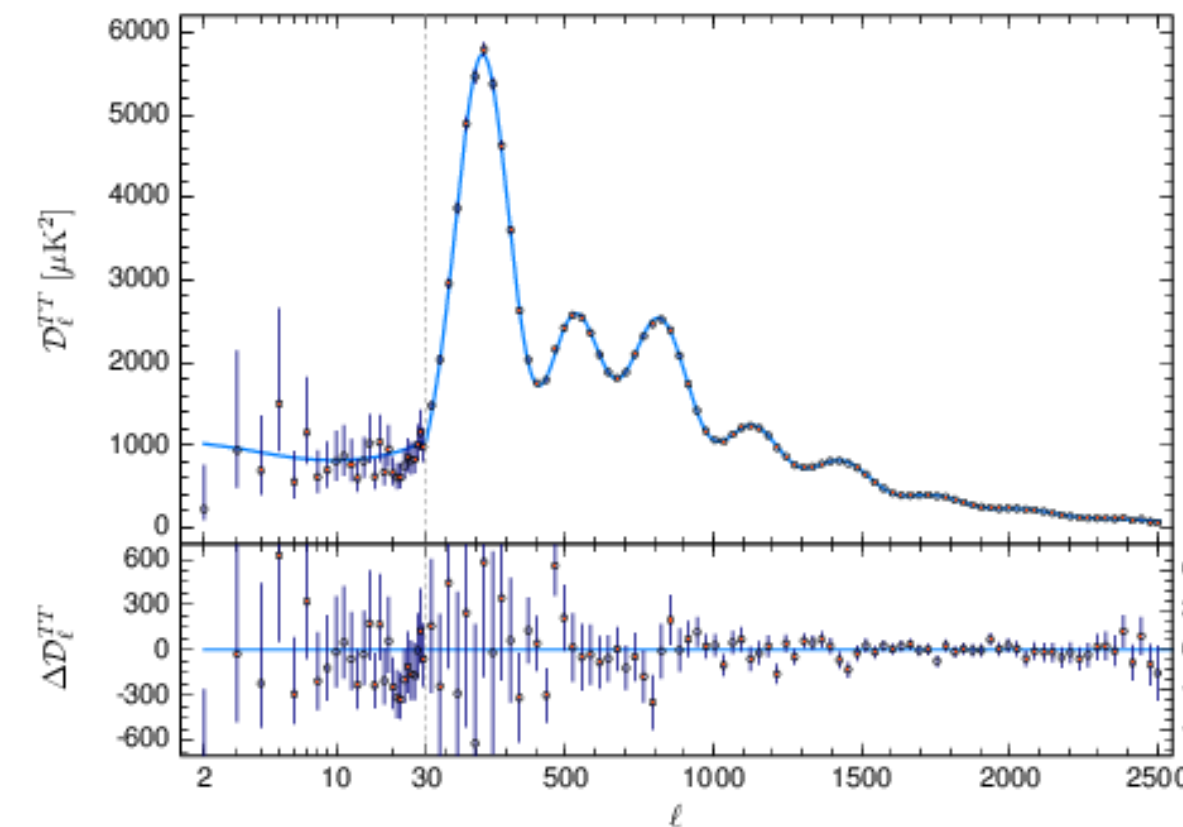
Evidence for Dark Matter



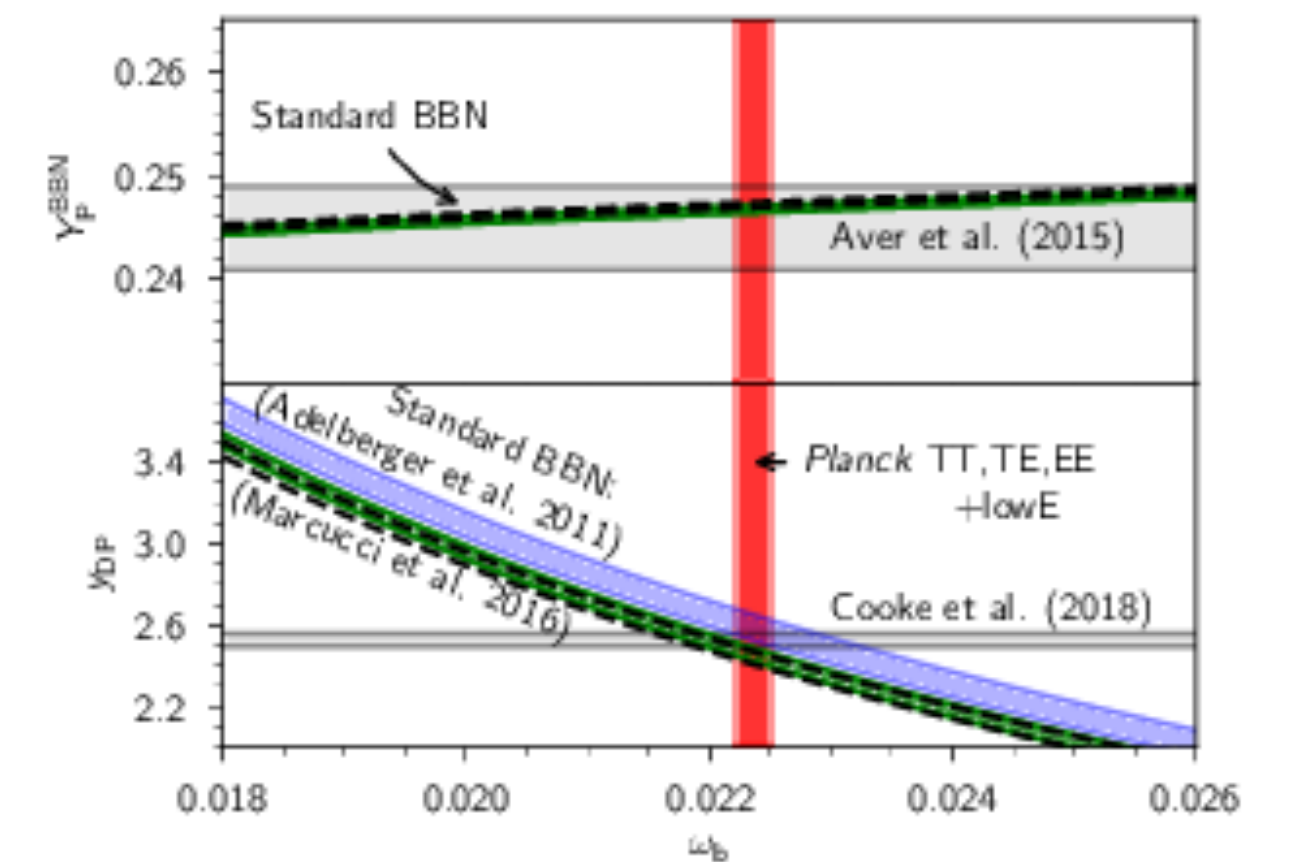
Galaxy rotation curves



Gravitational Lensing



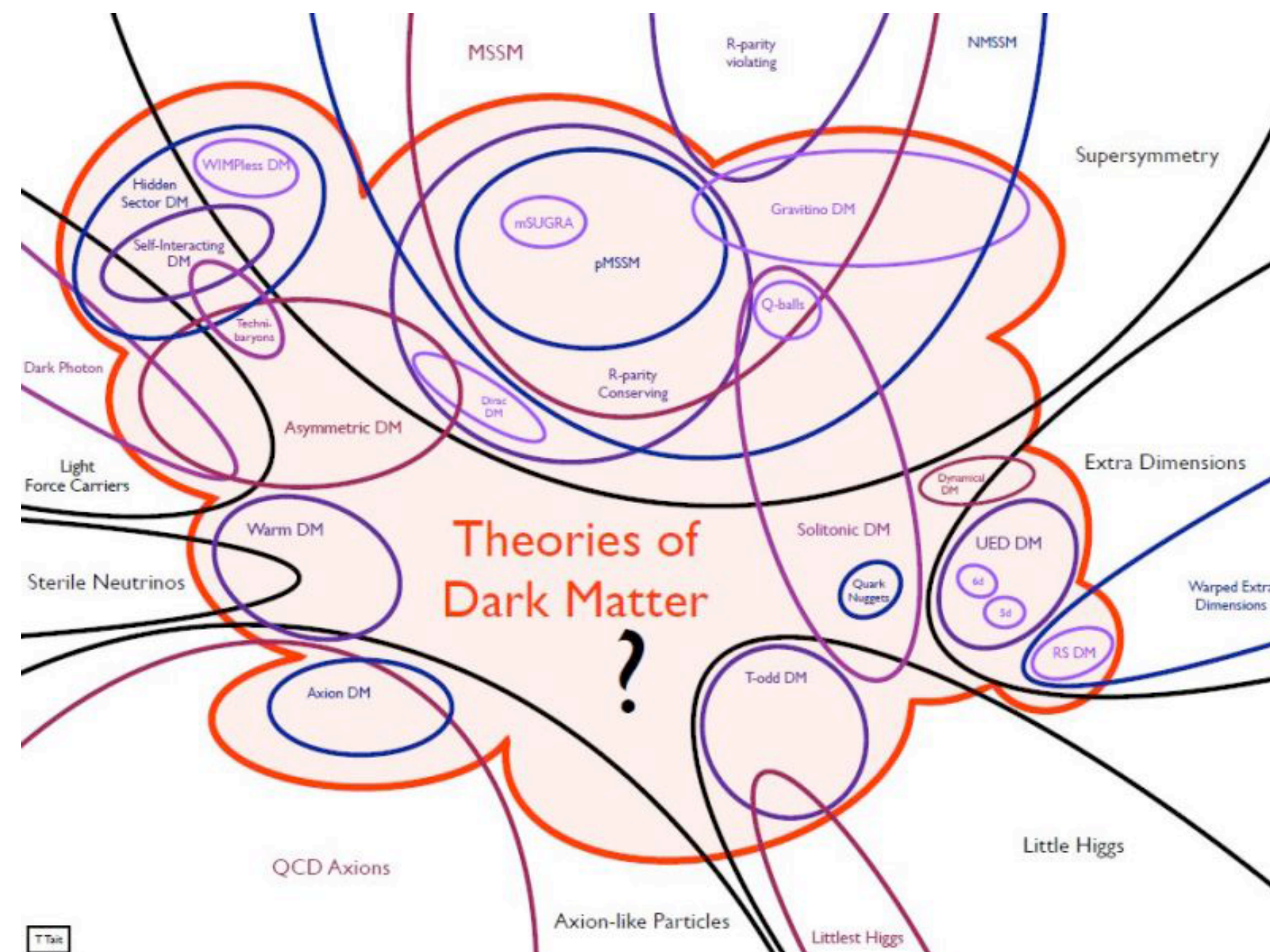
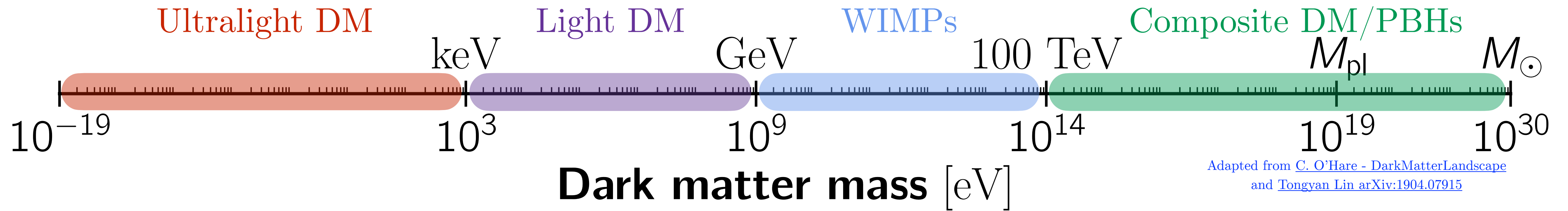
CMB



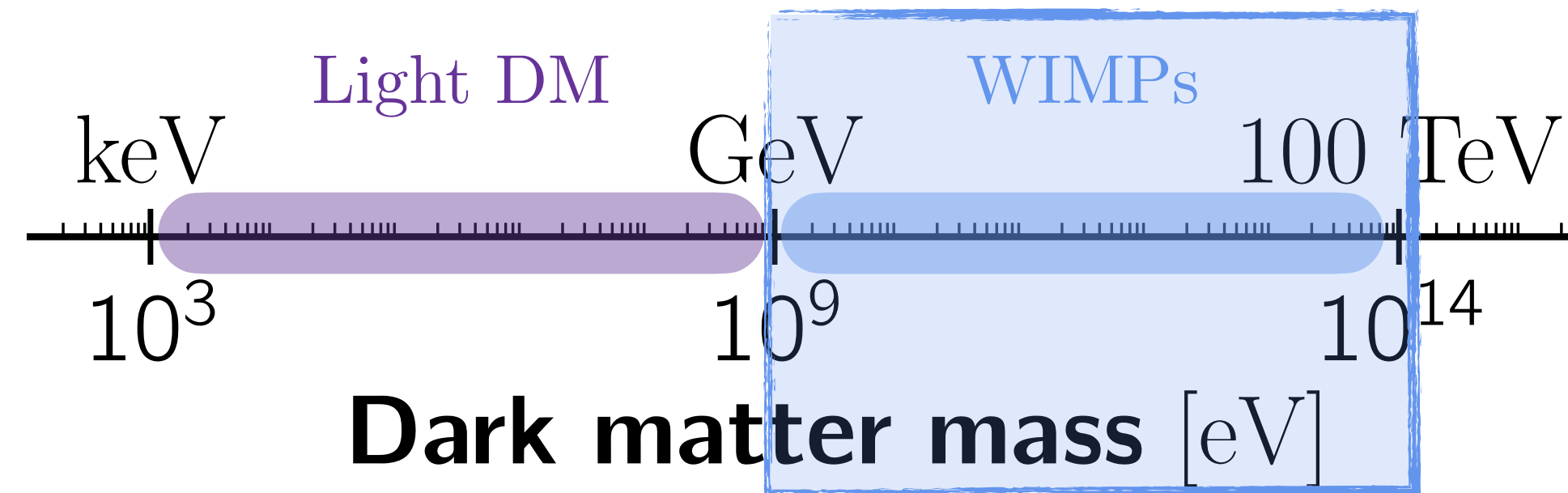
BBN

These observations tell us only about the *macroscopic* properties of DM. How can we probe the *microscopic* properties i.e. mass, non-gravitational interactions?

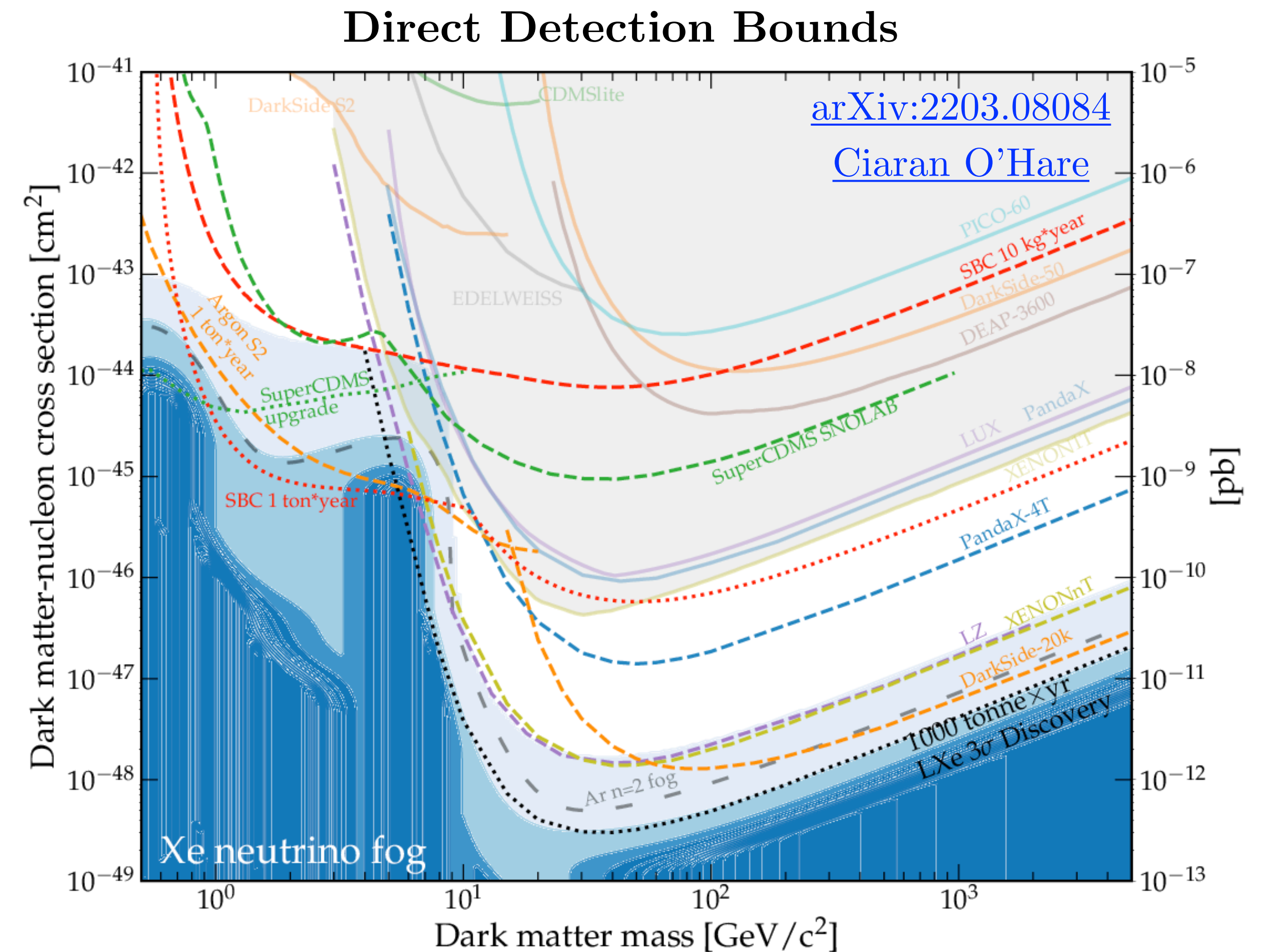
What even is DM?



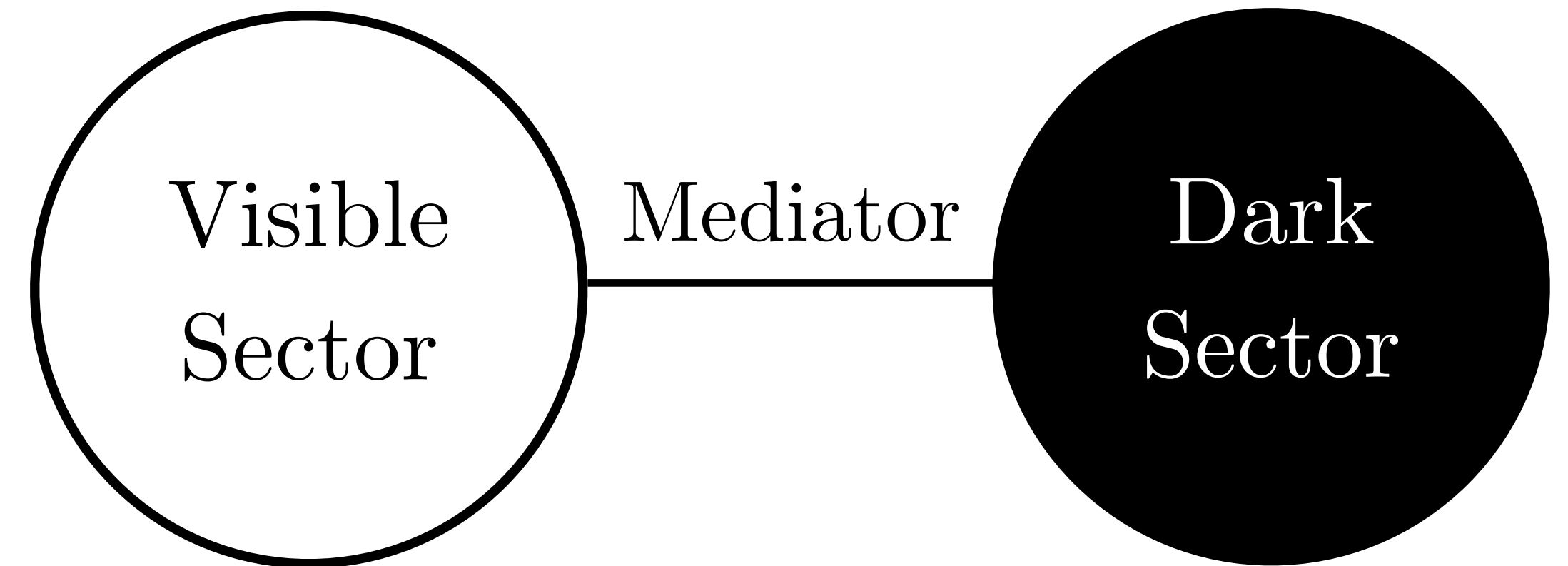
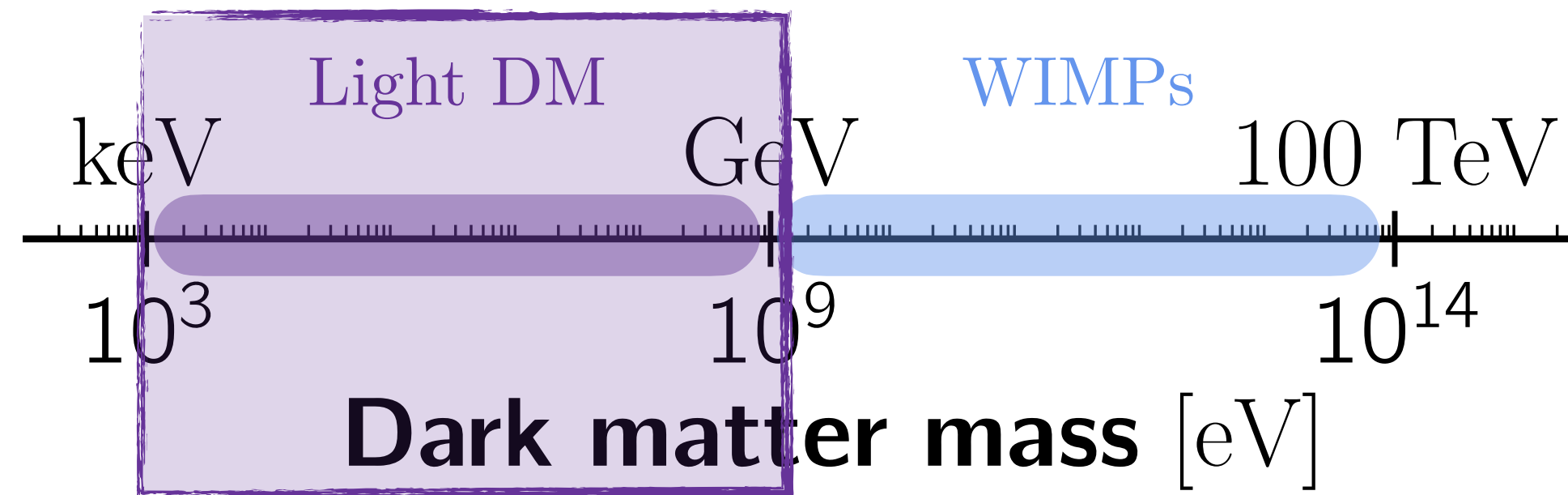
Weakly Interacting Massive Particles



- Traditional idea - *DM is a thermal relic*
- Direct detection bounds are becoming very constraining. *Push to smaller couplings. How to get beyond the neutrino floor?*
- *Alternative: go to lower masses where there are weaker bounds*

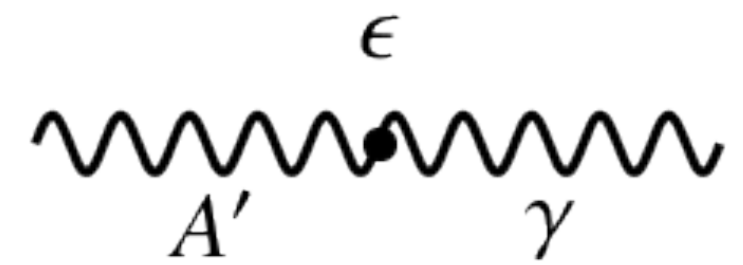


Light Dark Matter and Dark Sectors

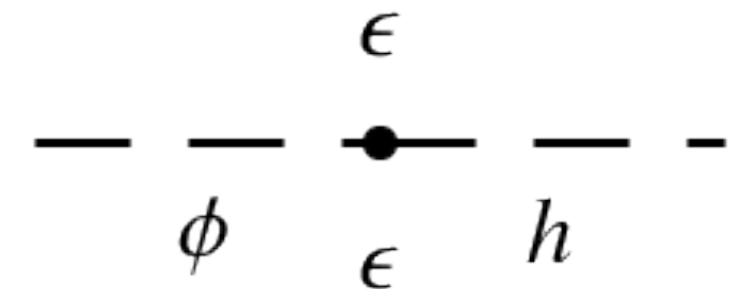


- Lee-Weinberg bound \rightarrow light thermal DM requires **new light mediators**
- Light mediators must be **SM singlets** \rightarrow **portal models**
- **Dark sectors** = DM + mediator + other SM singlet particles

1. Dark Photon: $\epsilon F^{\mu\nu} F'_{\mu\nu}$



2. Dark Higgs: $\epsilon |h|^2 |s|^2$



3. Heavy Neutrino: $\epsilon \ell h N$



Non-Thermal DM Candidates

Axions



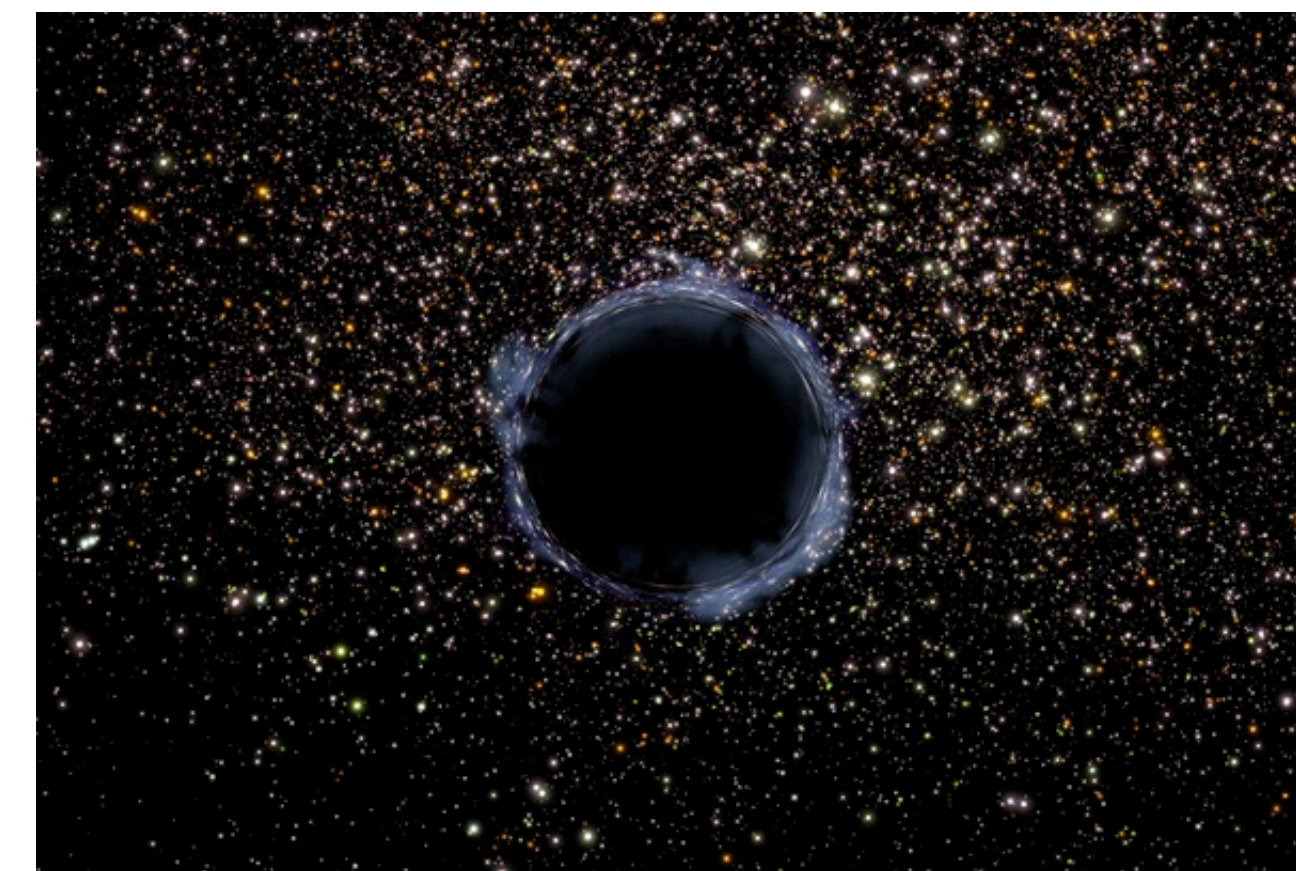
Composite/Heavy DM



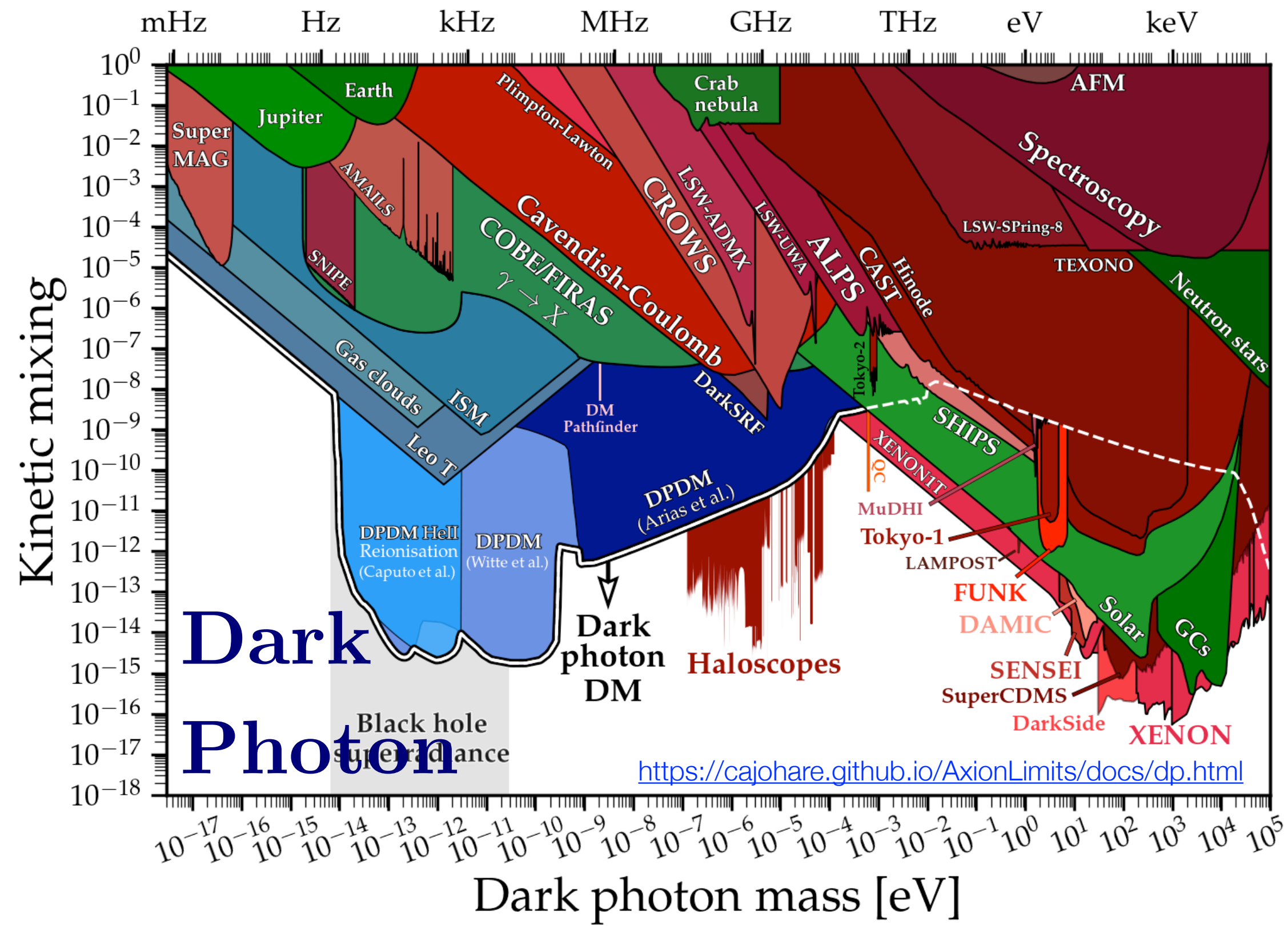
Ultra-light/wave DM



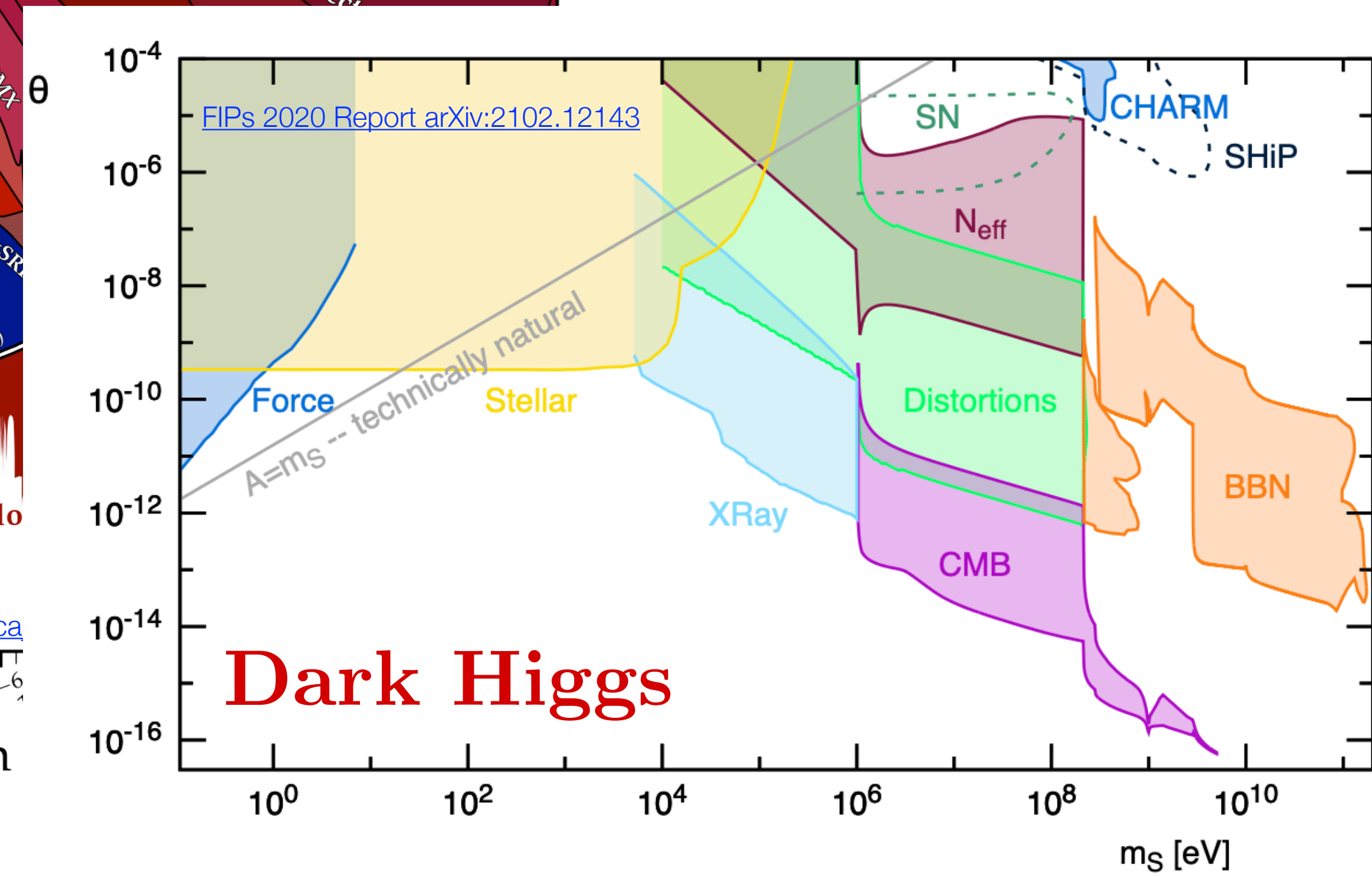
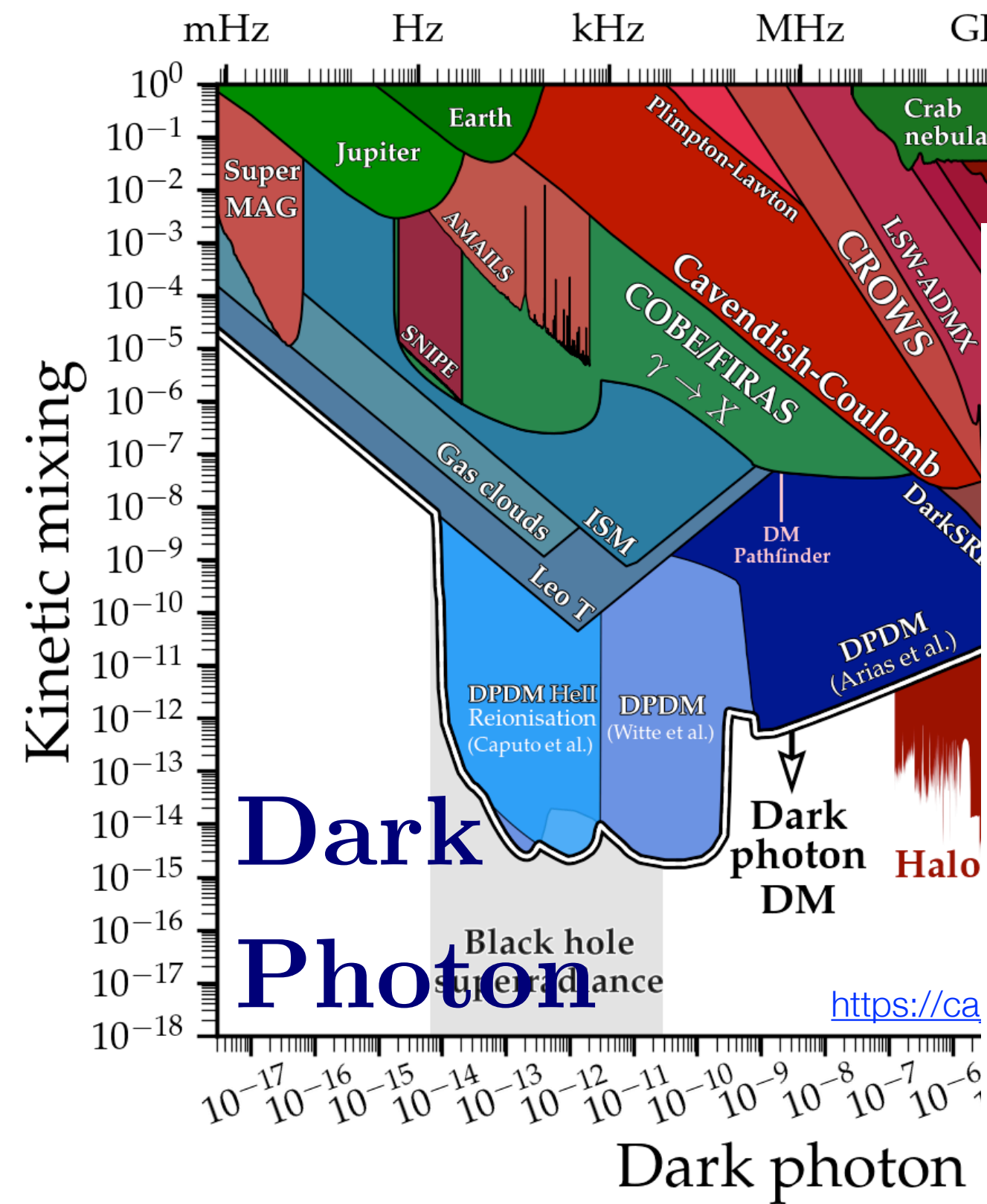
Primordial Black Holes



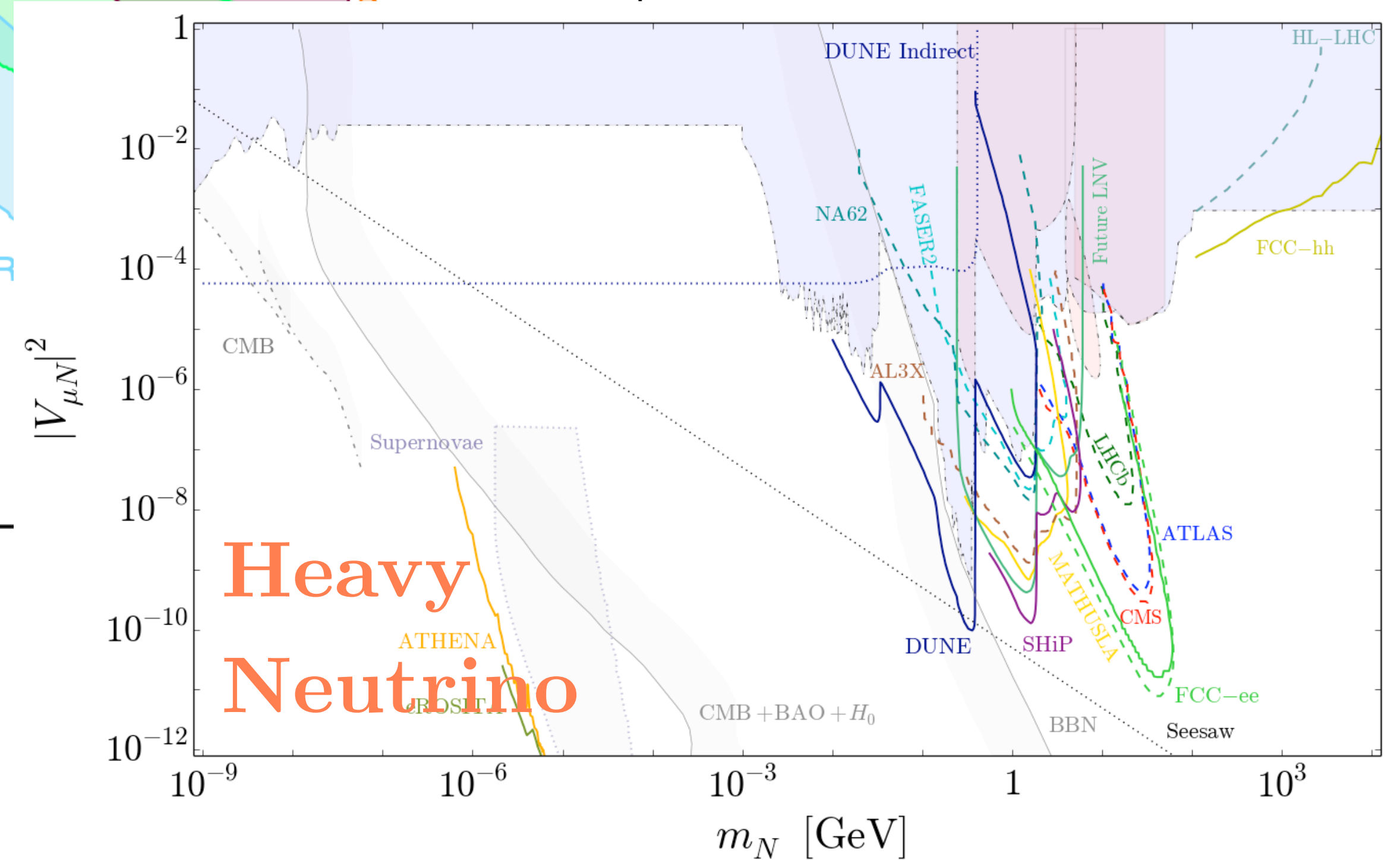
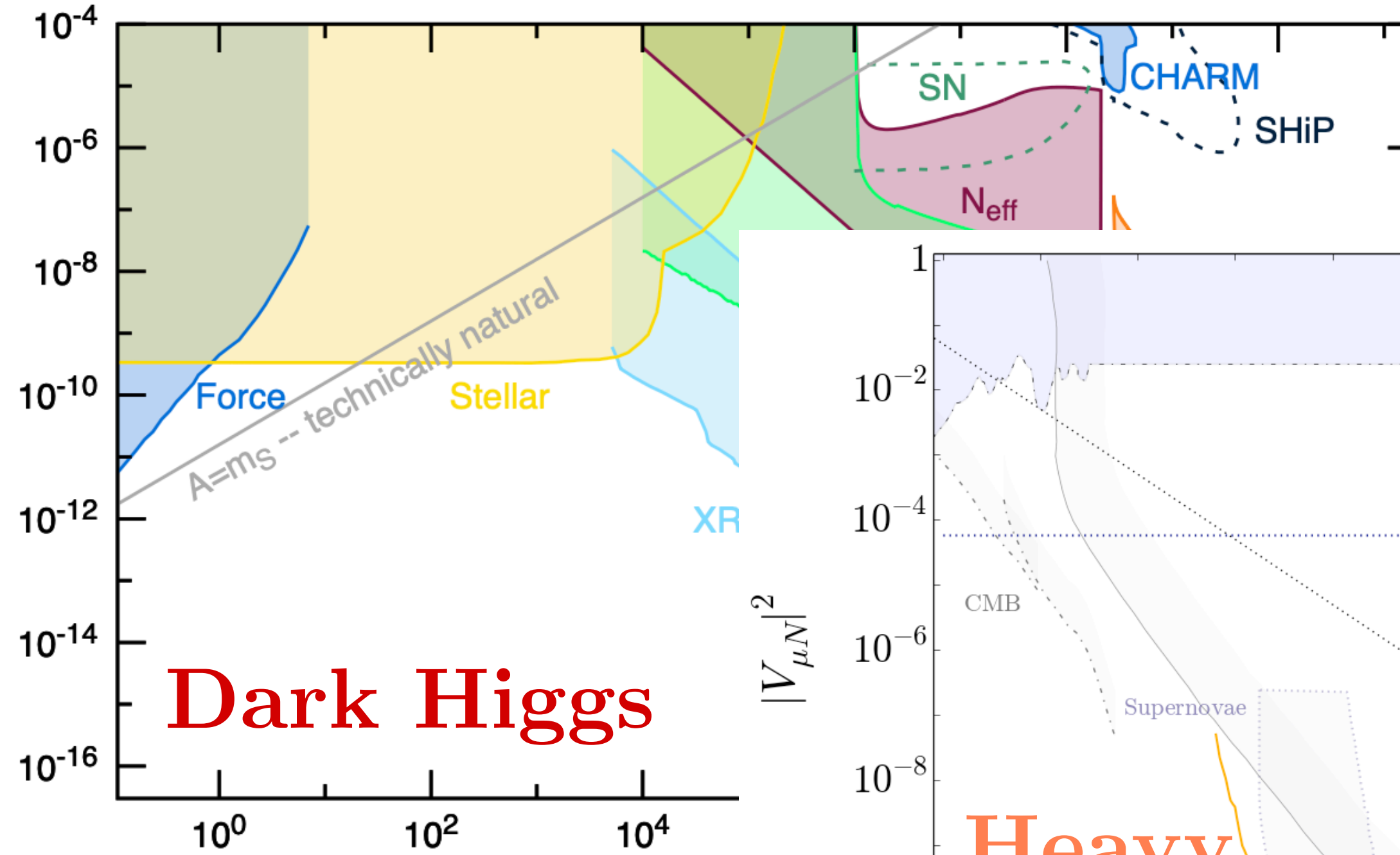
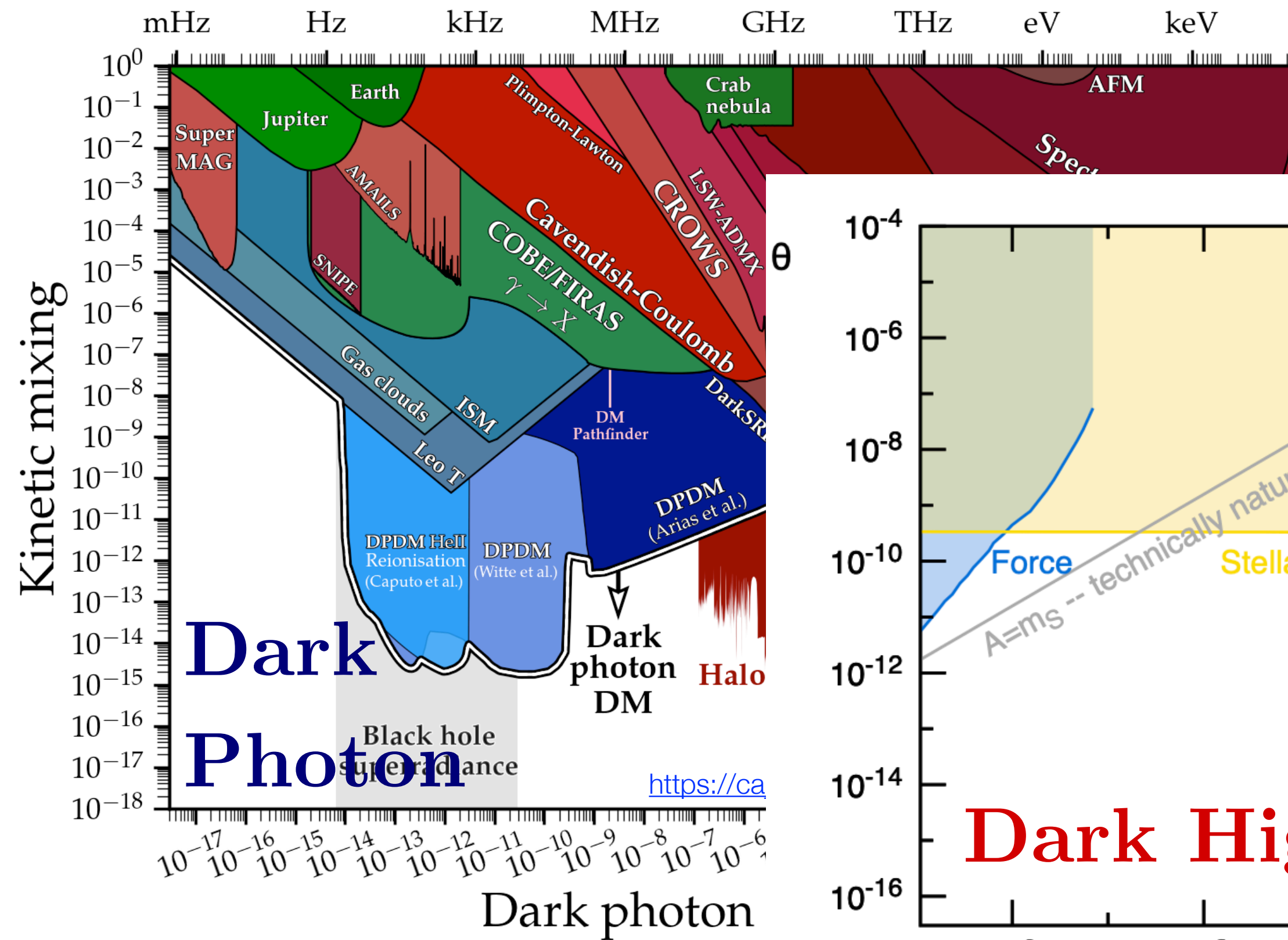
Experimental Results



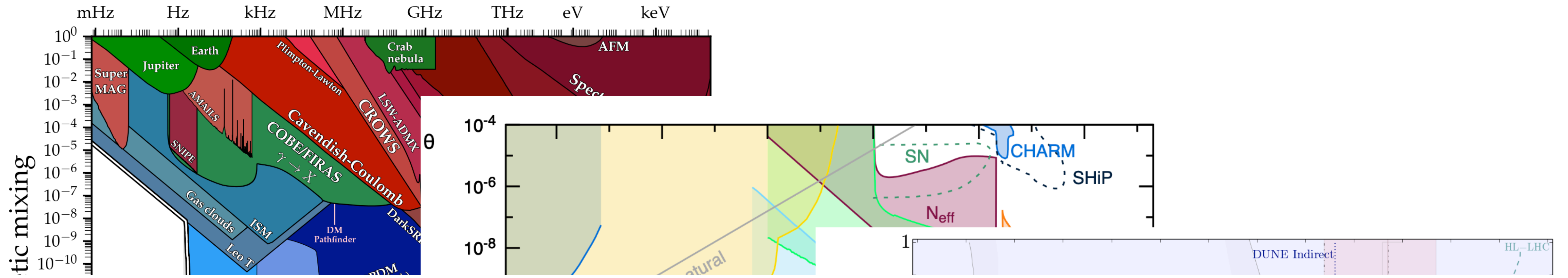
Experimental Results



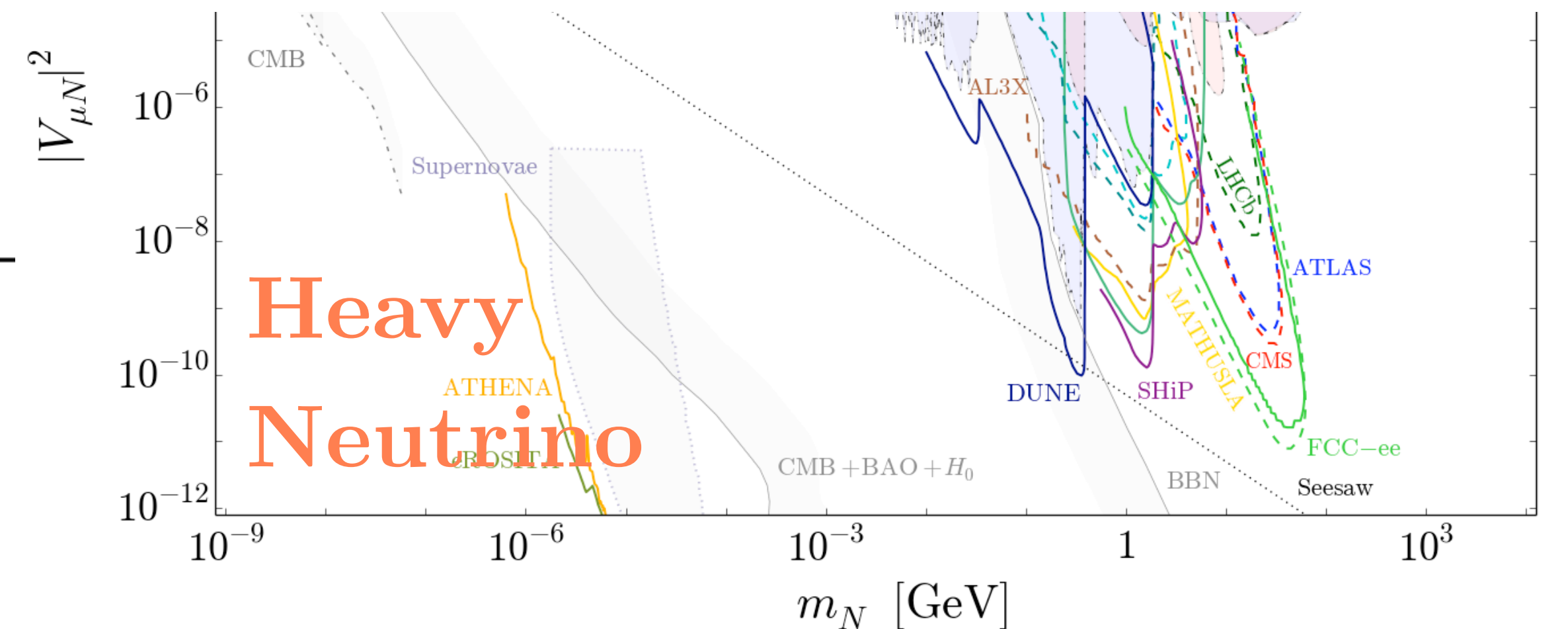
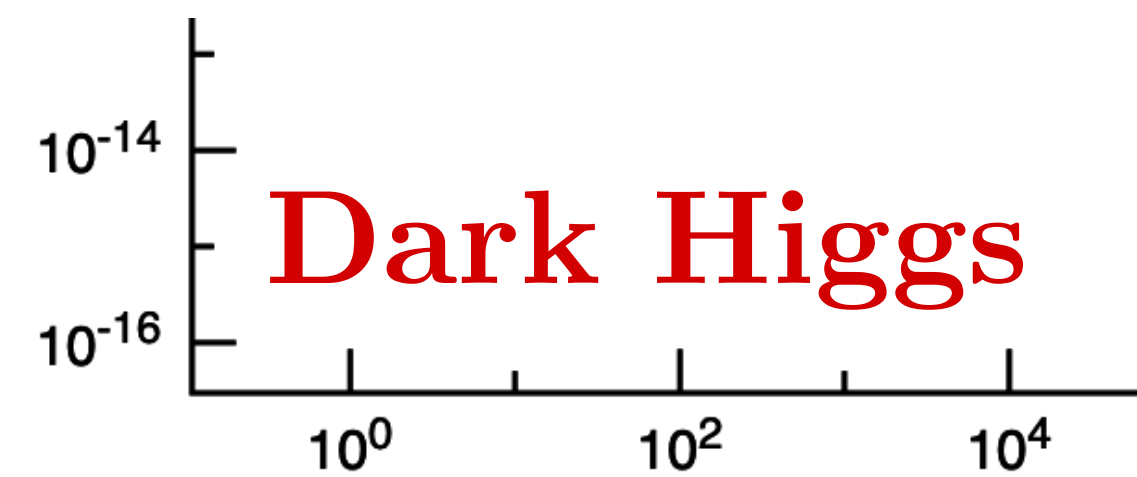
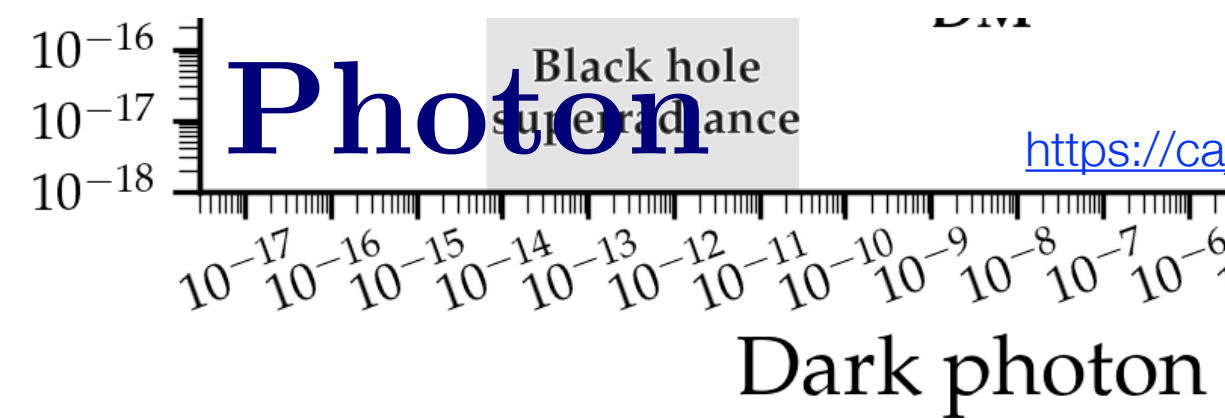
Experimental Results



Experimental Results

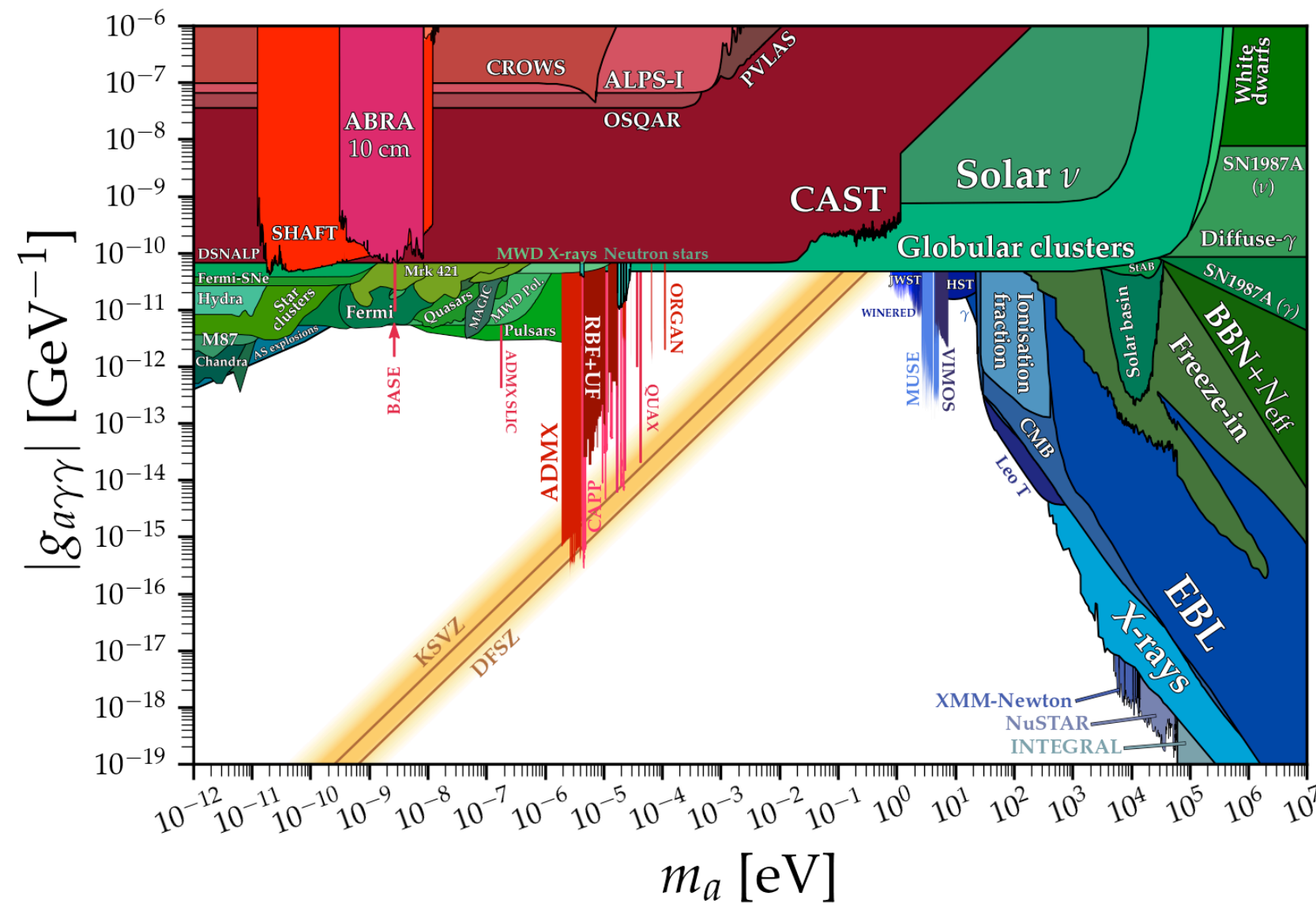


No DM/dark sector signal

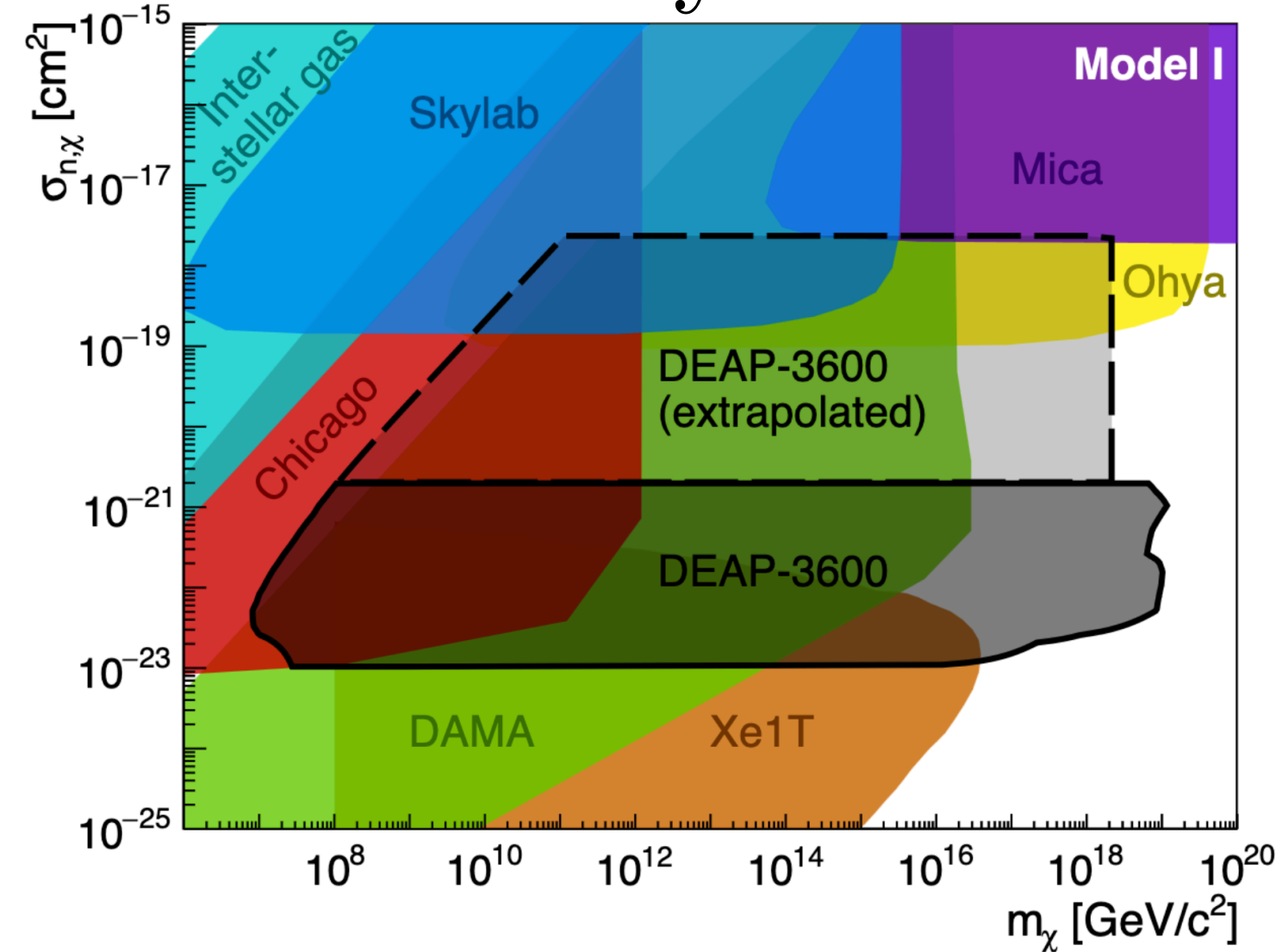


Experimental Results

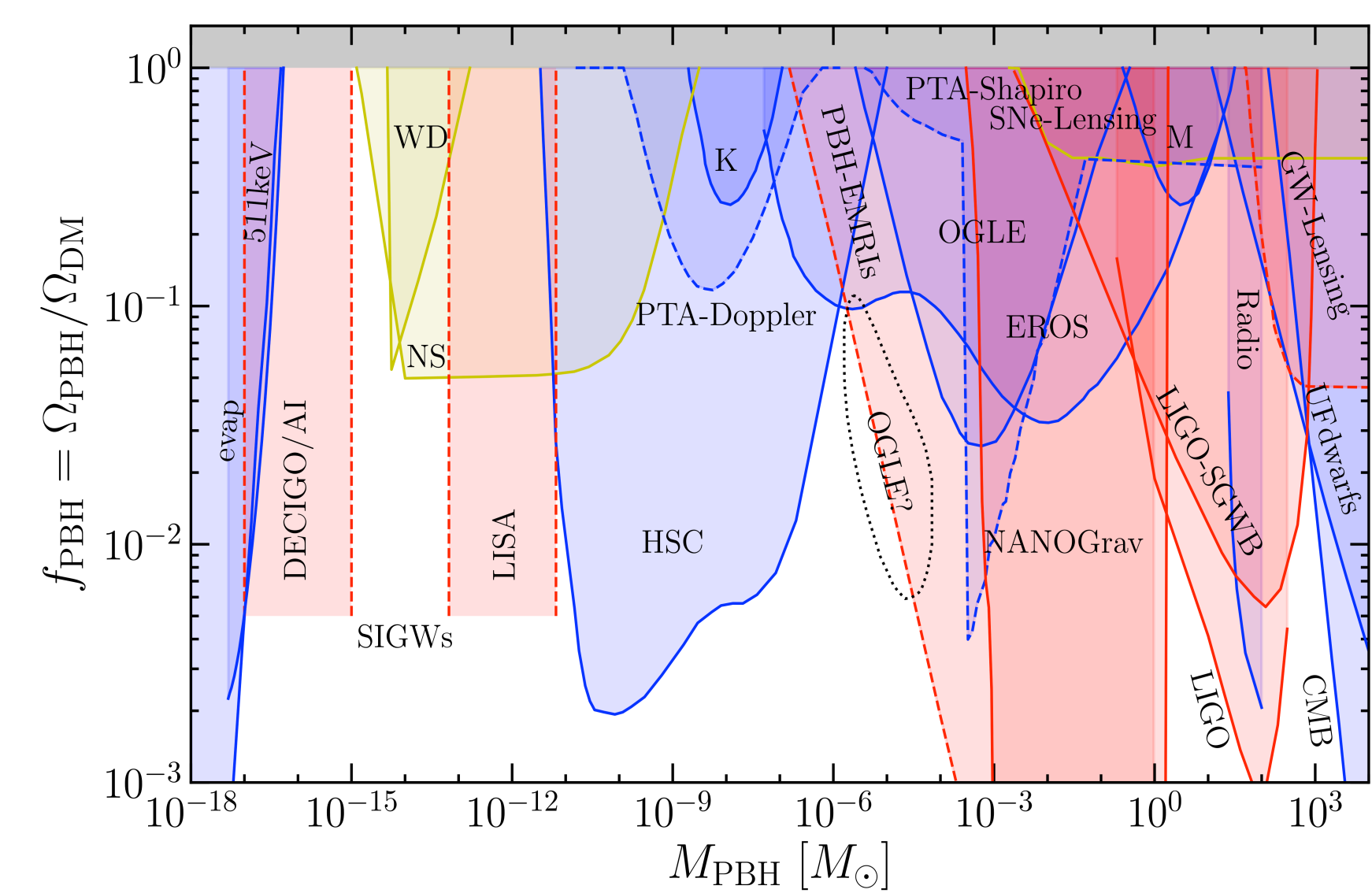
Axions



Heavy DM



PBHs



No DM/dark sector signal

Next steps?

- No dark matter signal has been observed. Where do we go from here?

Next steps?

- No dark matter signal has been observed. Where do we go from here?
 1. Maximally Optimistic option: We need to build all the experiments.

Next steps?

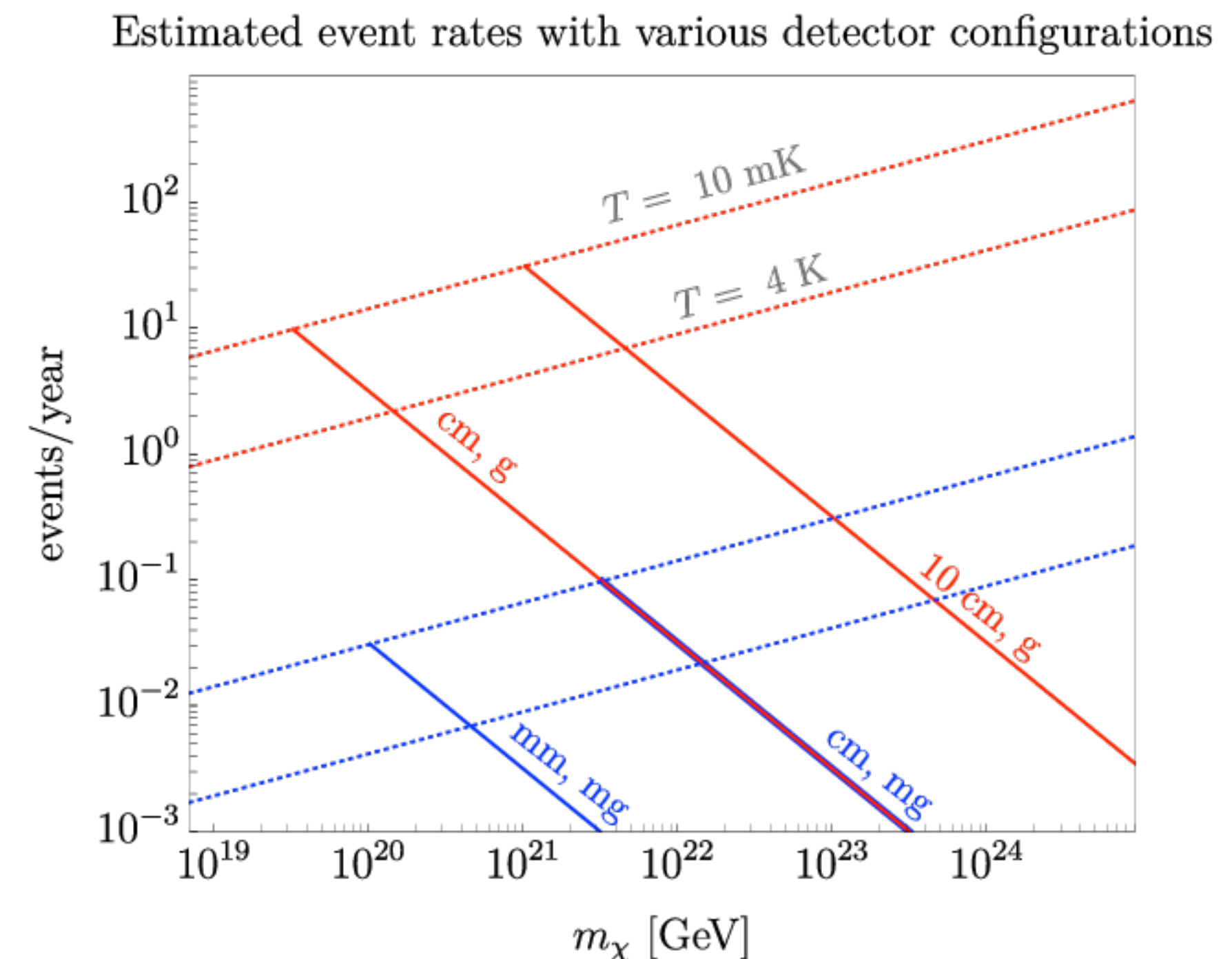
- No dark matter signal has been observed. Where do we go from here?
 1. Maximally Optimistic option: We need to build all the experiments.
 2. Maximally Pessimistic option: dark matter has **no non-gravitational interactions**.

Next steps?

- No dark matter signal has been observed. Where do we go from here?
 1. Maximally Optimistic option: We need to build all the experiments.
 2. Maximally Pessimistic option: dark matter has **no non-gravitational interactions**.

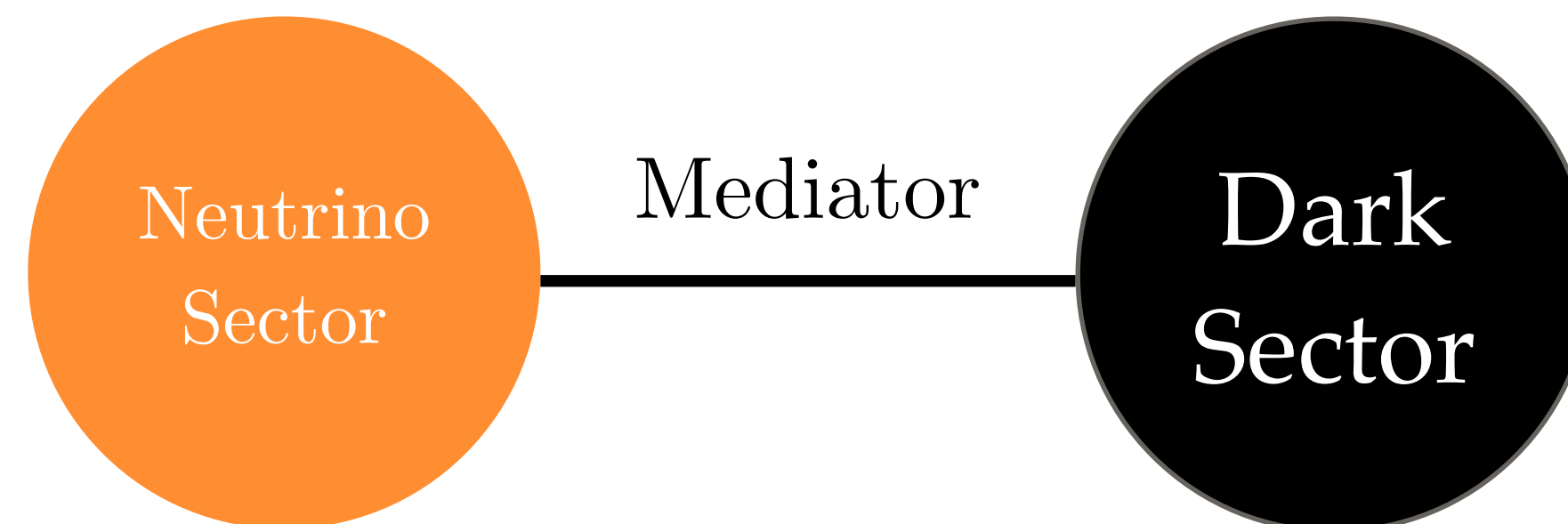
The Windchime Project: Gravitational Detection of Dark Matter in the Laboratory

Small window where this could work so
we better hope that DM has this mass!



Next steps?

- No dark matter signal has been observed. Where do we go from here?
 1. Maximally Optimistic option: We need to build all the experiments.
 2. Maximally Pessimistic option: dark matter has no non-gravitational interactions.
 3. Searches for DM assume that DM interacts with visible stuff (e.g. photons, electrons, protons). *What if DM is more elusive than we thought? What if DM only interacts with neutrinos?*

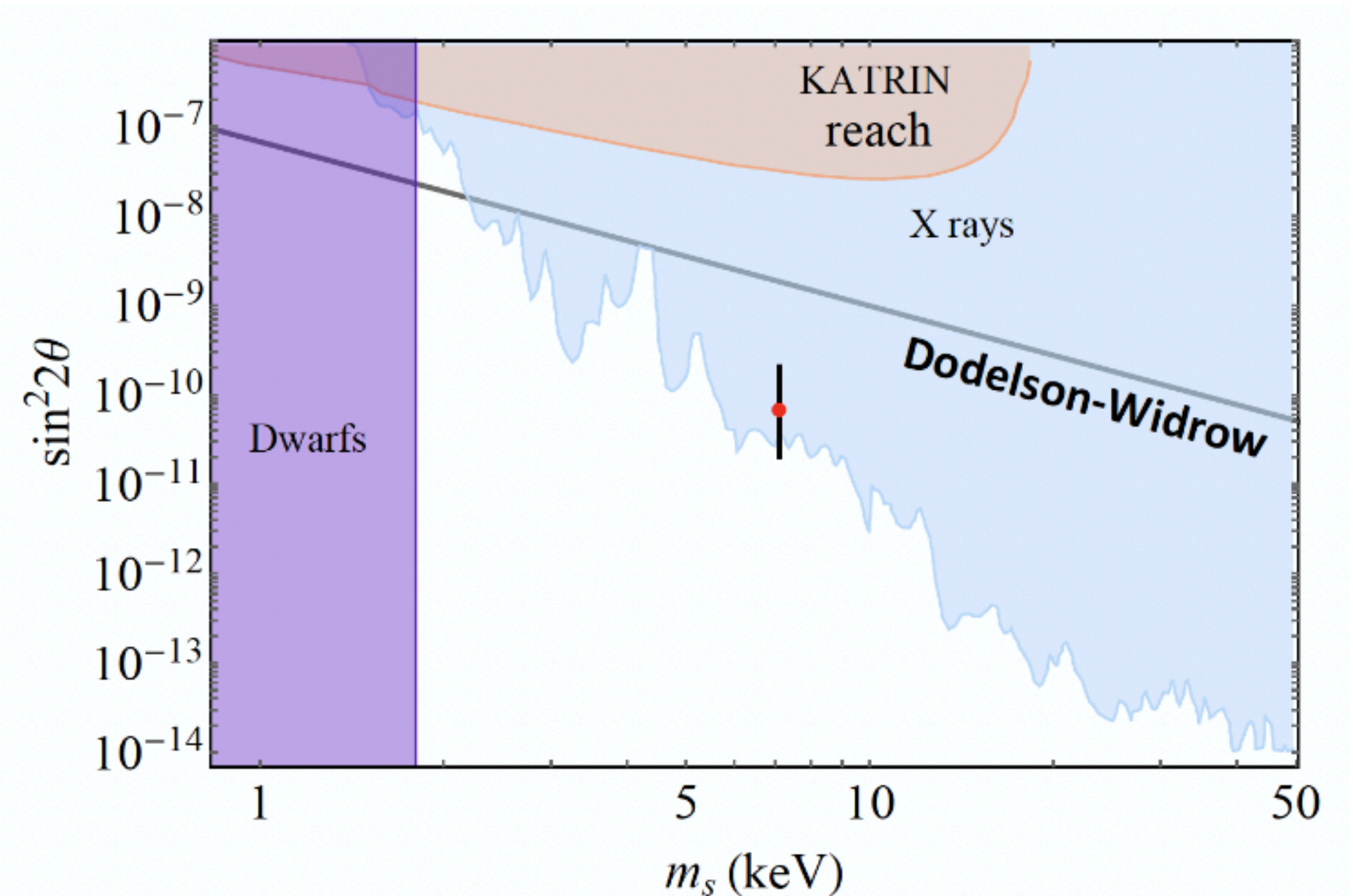


Sterile Neutrino Dark Matter

- keV-scale singlet fermion that mixes only with the SM neutrinos

$$\nu_4 = \nu_s \cos \theta + \nu_a \sin \theta$$

- Sterile neutrino produced via Dodelson-Widrow Mechanism
- Indirect detection via one-loop decay $\nu_s \rightarrow \nu_a \gamma$ with X-ray line at $E_\gamma = m_4/2$



*SνDM is almost completely excluded.
Can we save Dodelson-Widrow?*

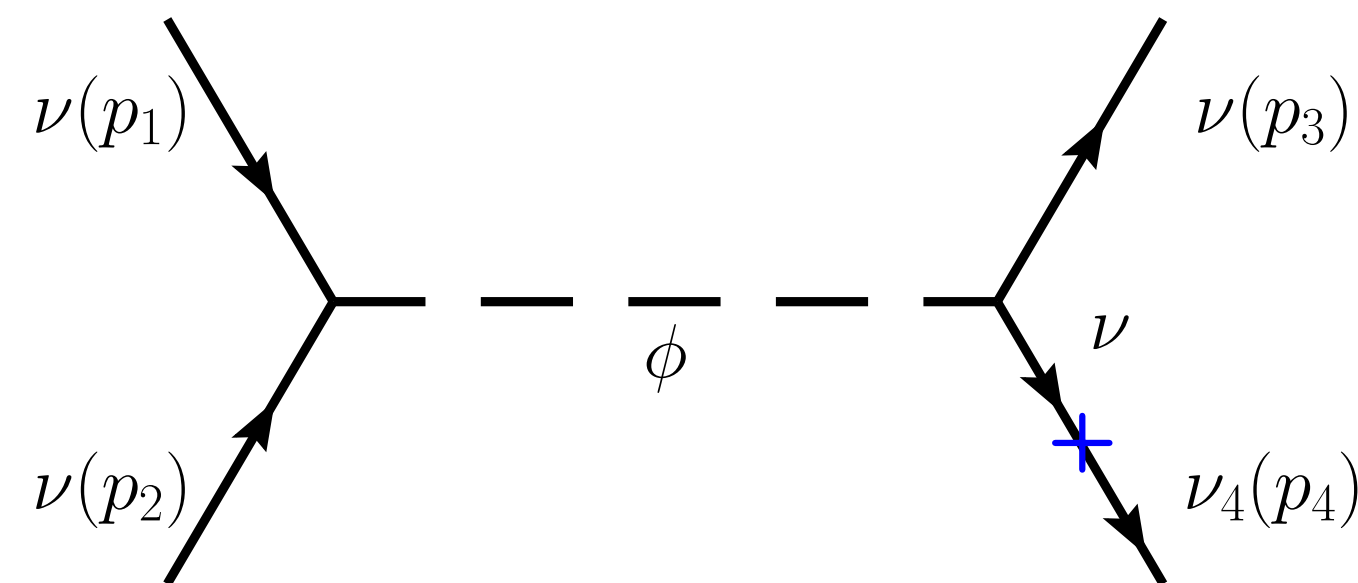
A Neutrinophilic Scalar Mediator

- Schematically, the sterile neutrino relic abundance is

$$\Omega \sim \Gamma \times \sin^2(2\theta)$$

- If $\Gamma = \Gamma_W$, then a large angle is required \rightarrow X-ray constraints.
- Smaller mixing angle by increasing the interaction rate? Yes! Introduce a scalar field ϕ of mass m_ϕ that mediates *new self interactions among SM neutrinos*.

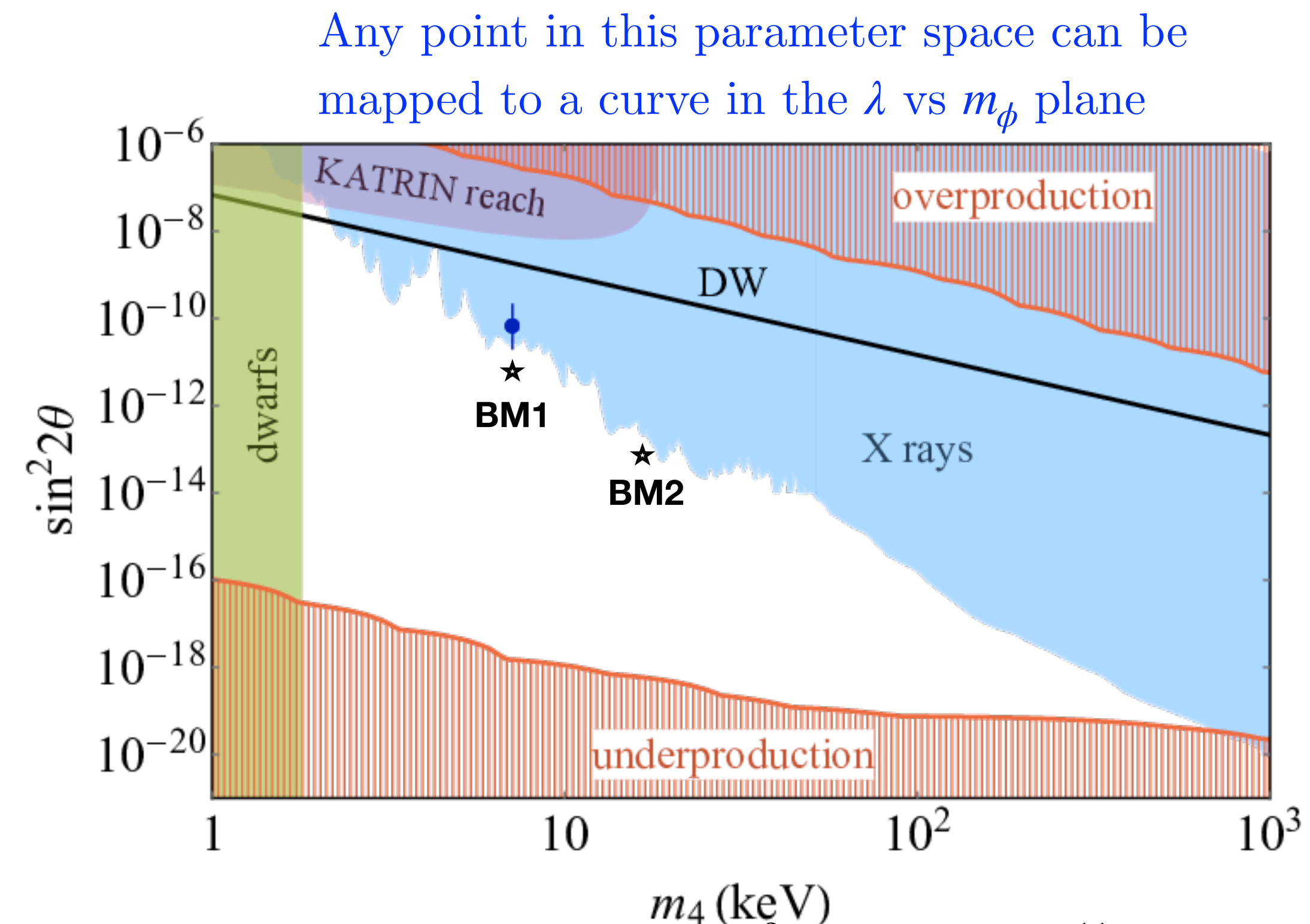
$$\mathcal{L} \supset \frac{1}{2} \lambda_{\alpha\beta} \nu_\alpha \nu_\beta \phi + \text{h.c.}$$



Larger rate than the weak interactions keeps SM neutrinos in contact for a longer period of time to build up the DM abundance!

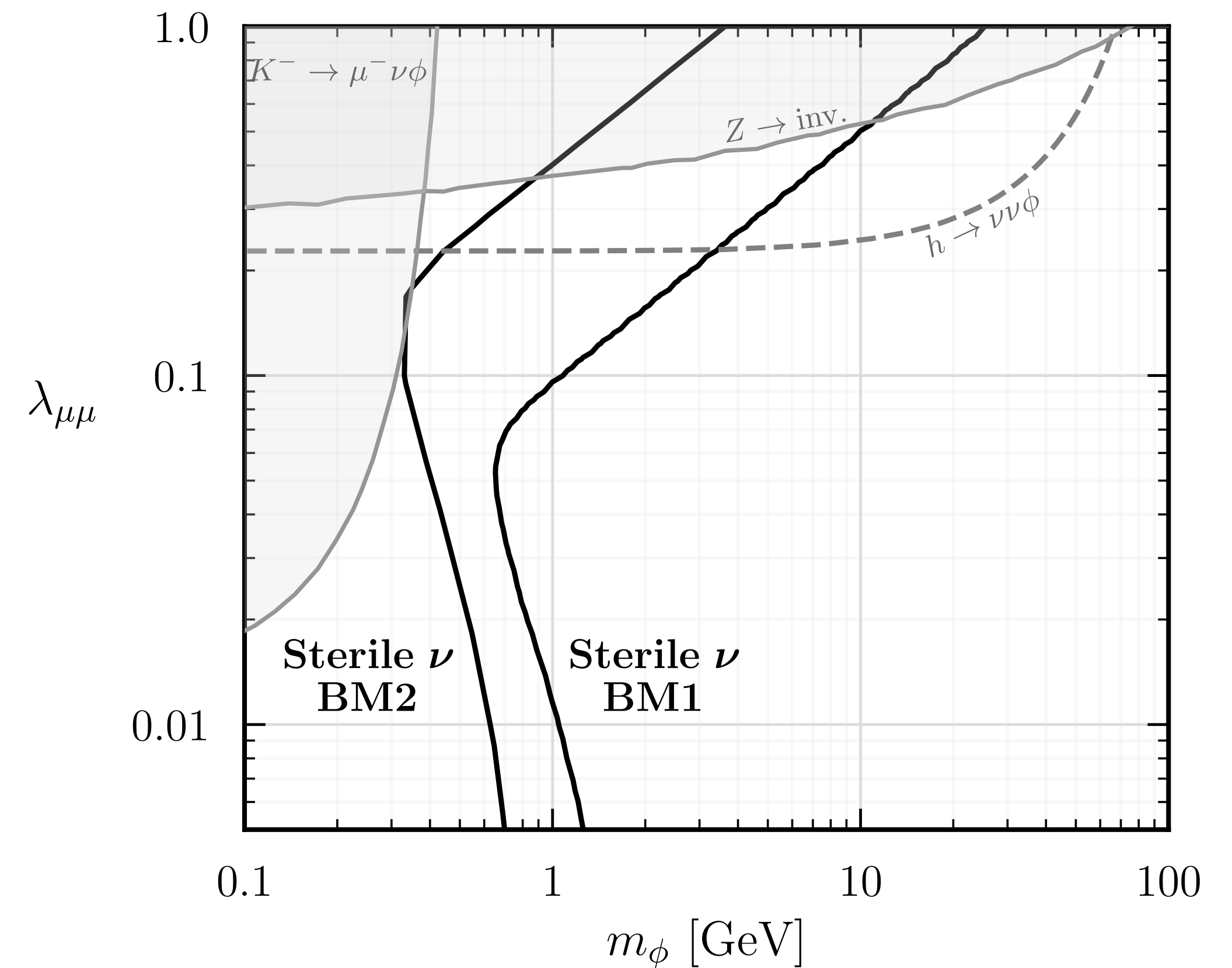
A Neutrinophilic Scalar Mediator

- New production mode \rightarrow don't have to live on DW line!



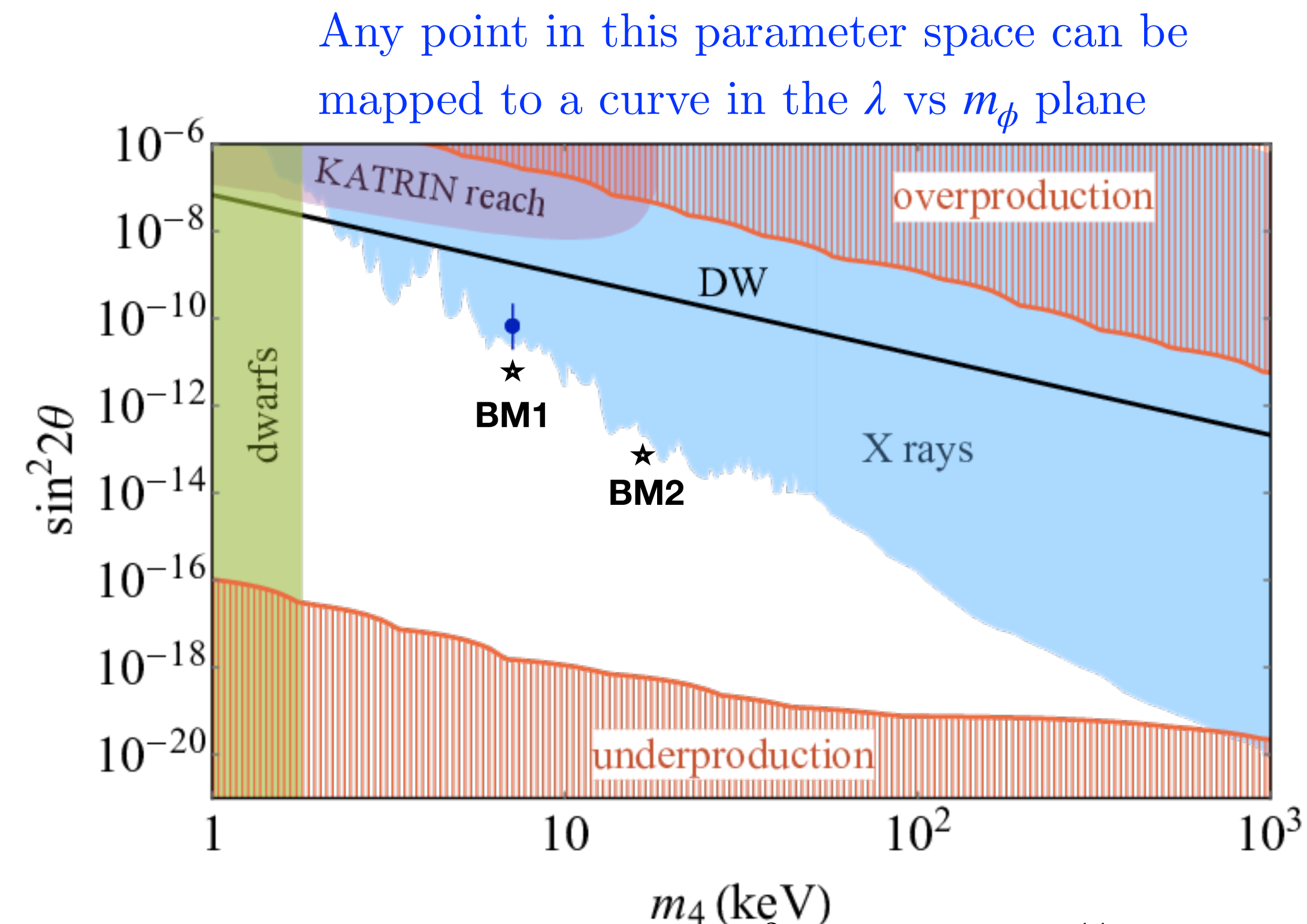
BM1 : $m_4 = 7\text{keV}$, $\sin^2(2\theta) = 7 \times 10^{-11}$

BM2 : $m_4 = 21\text{keV}$, $\sin^2(2\theta) = 1.4 \times 10^{-13}$



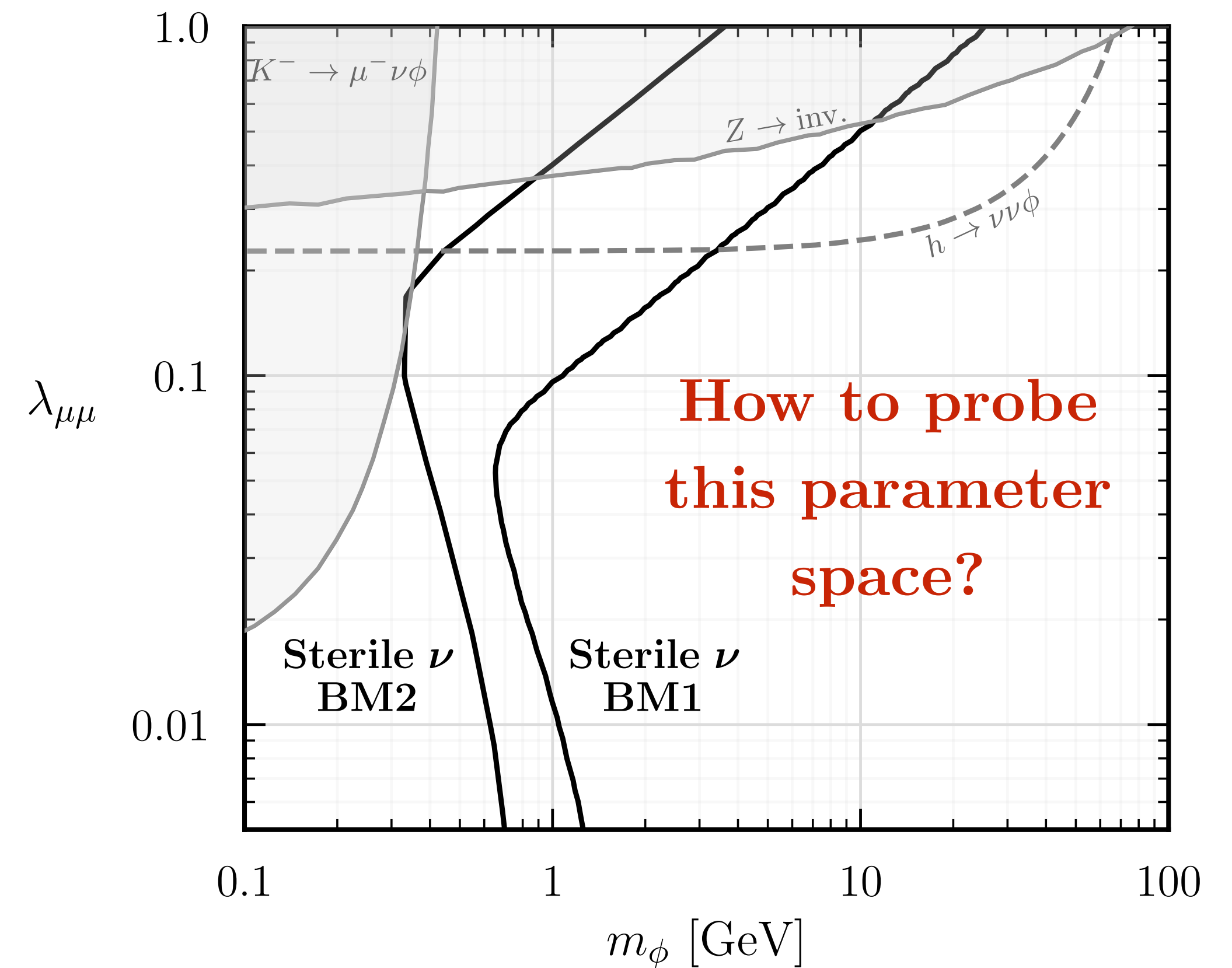
A Neutrinophilic Scalar Mediator

- New production mode \rightarrow don't have to live on DW line!



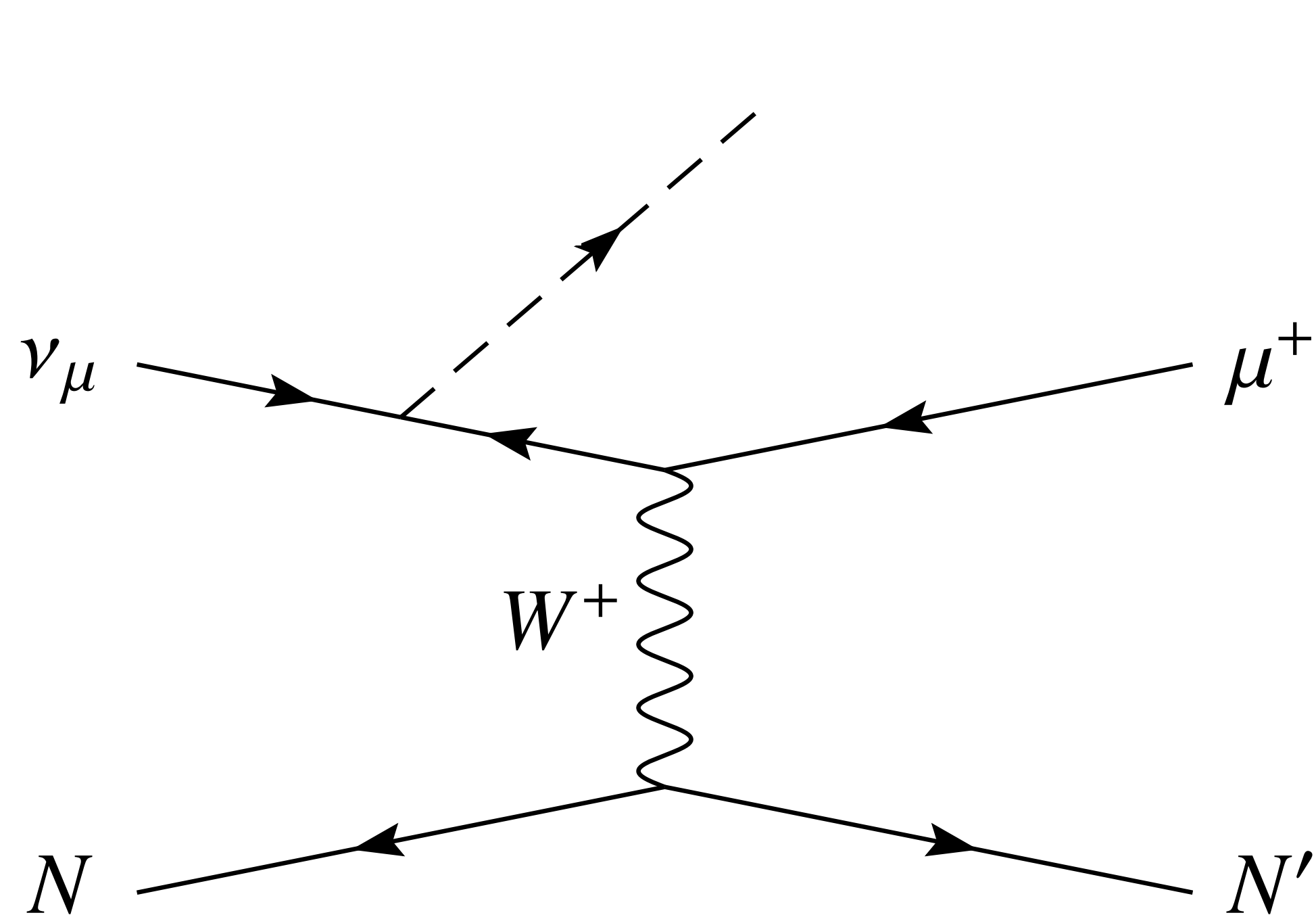
BM1 : $m_4 = 7\text{keV}$, $\sin^2(2\theta) = 7 \times 10^{-11}$

BM2 : $m_4 = 21\text{keV}$, $\sin^2(2\theta) = 1.4 \times 10^{-13}$



The Mono-neutrino Signature

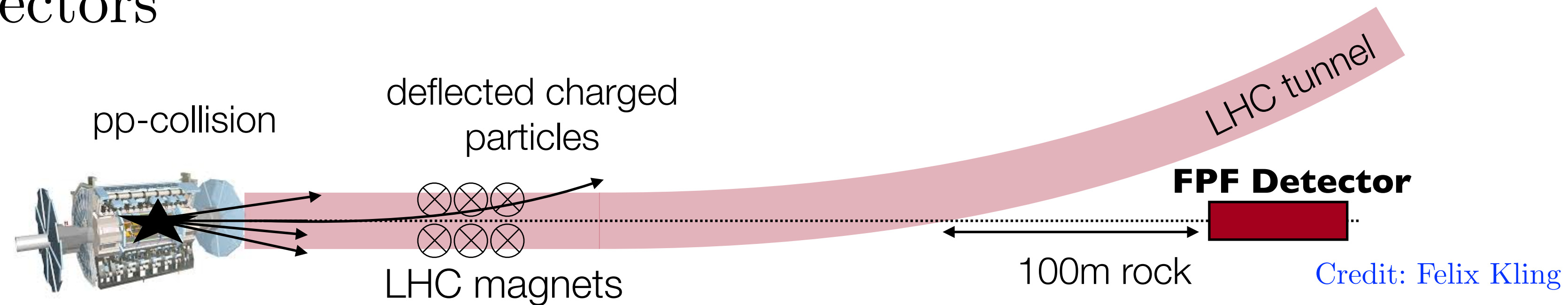
- Unique signature due to the neutrinophilic nature of the mediator: incoming neutrino radiates a scalar particle and then converts to a muon via CC interactions. [K. J. Kelly and Y. Zhang arXiv:1901.01259](#)



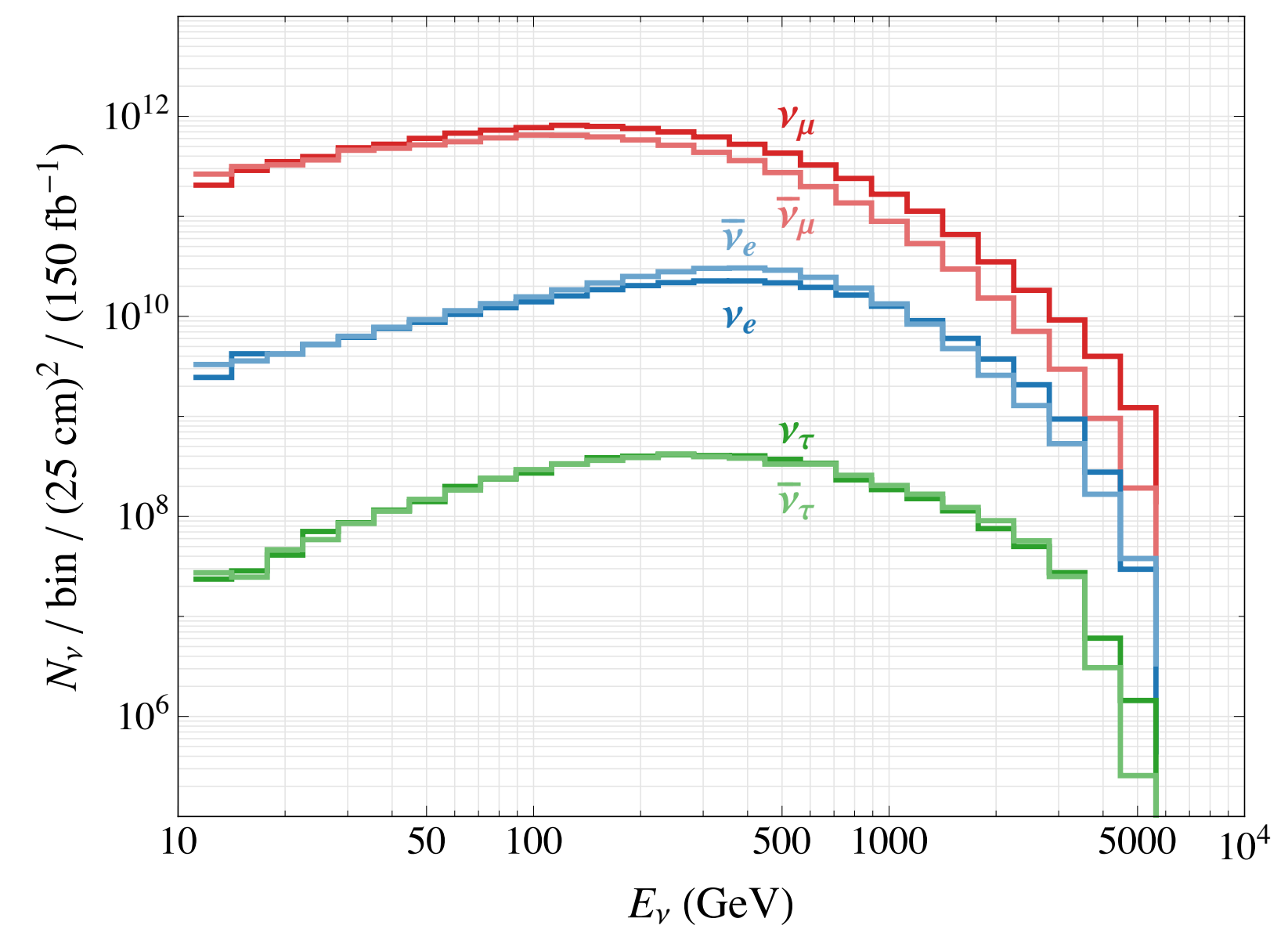
- Observable: **Missing transverse momentum** carried away by ϕ
- Similar in spirit to mono-X searches at the LHC, missing transverse momentum technique @ LDMX/DarkLight
- **High energy/intensity neutrino environments are excellent to probe this signature!**

LHC Forward Physics Facility

- A proposal to explore SM and BSM physics in the far forward region of LHC detectors



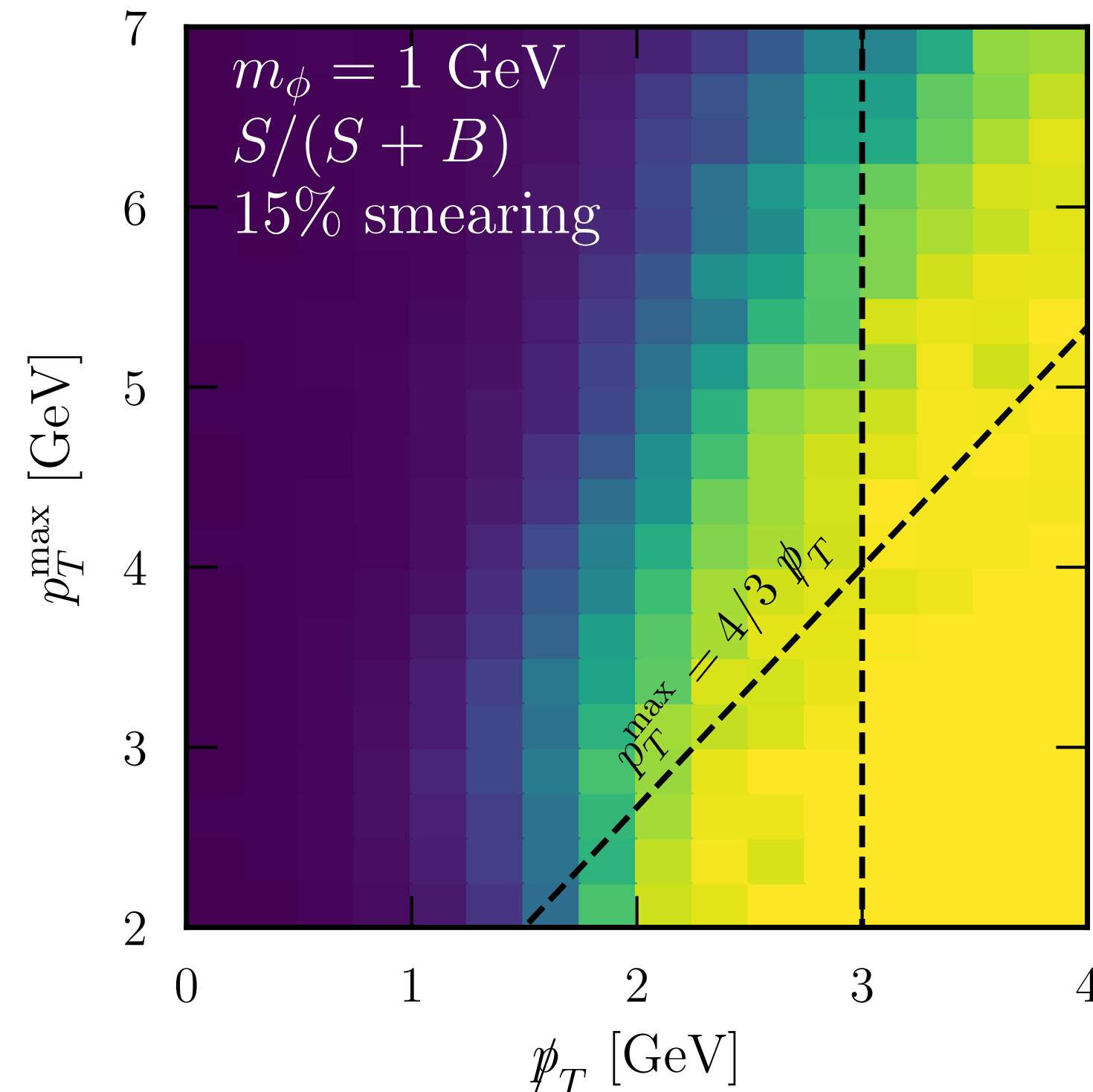
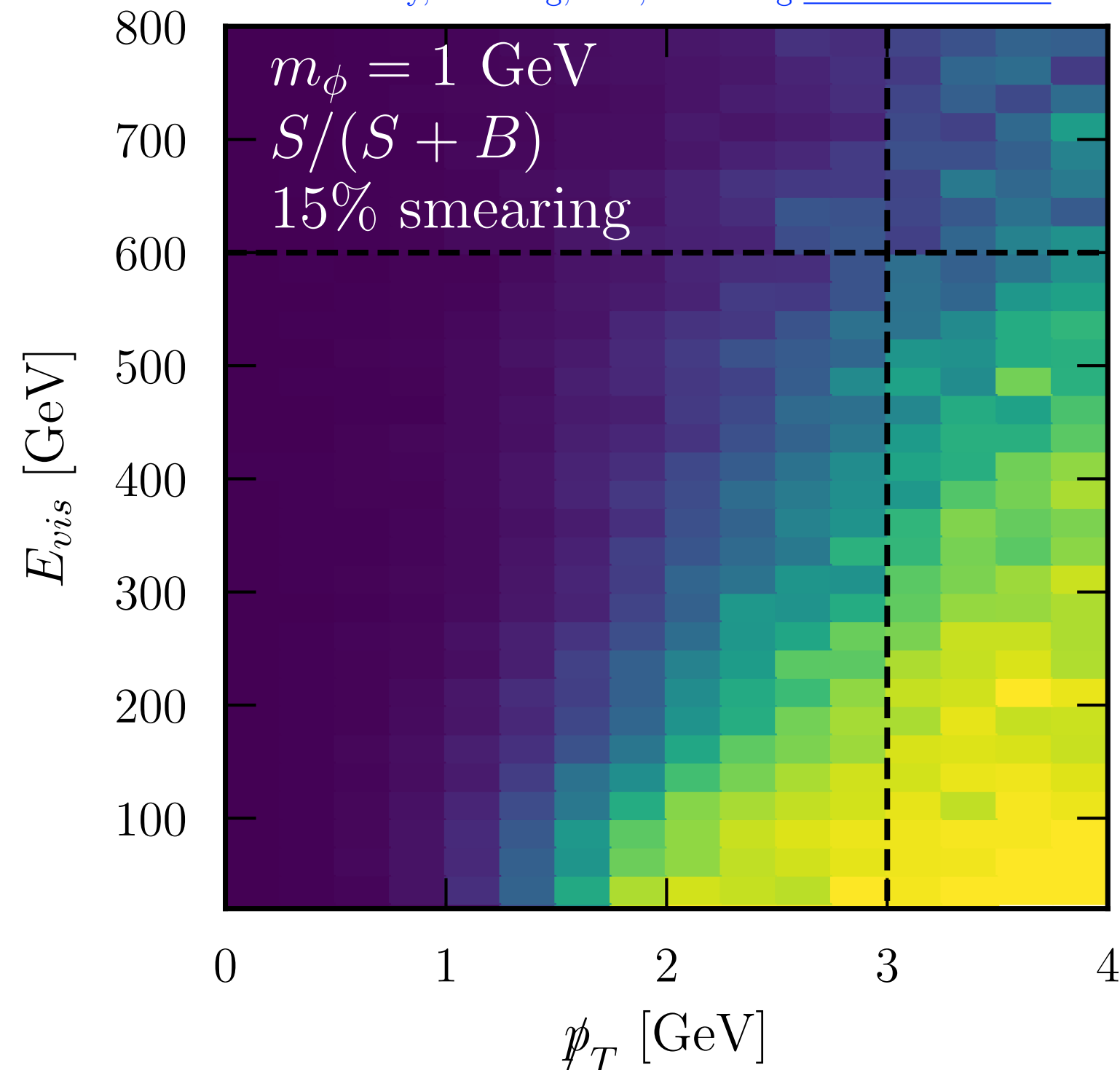
- Flux of high energy neutrinos can be used to probe our model!
- Advantages of LHC neutrinos:
 - High energy neutrinos can probe higher scalar masses
 - Neutrino scattering is DIS \rightarrow smaller uncertainties



Analysis Strategy

- Relevant observables:
 - Missing transverse momentum \cancel{p}_T
 - Total energy of all visible final states E_{vis}
 - Highest transverse momentum of visible final state objects p_T^{max}

K.J. Kelly, F. Kling, DT, Y. Zhang [arXiv:2111.05868](https://arxiv.org/abs/2111.05868)

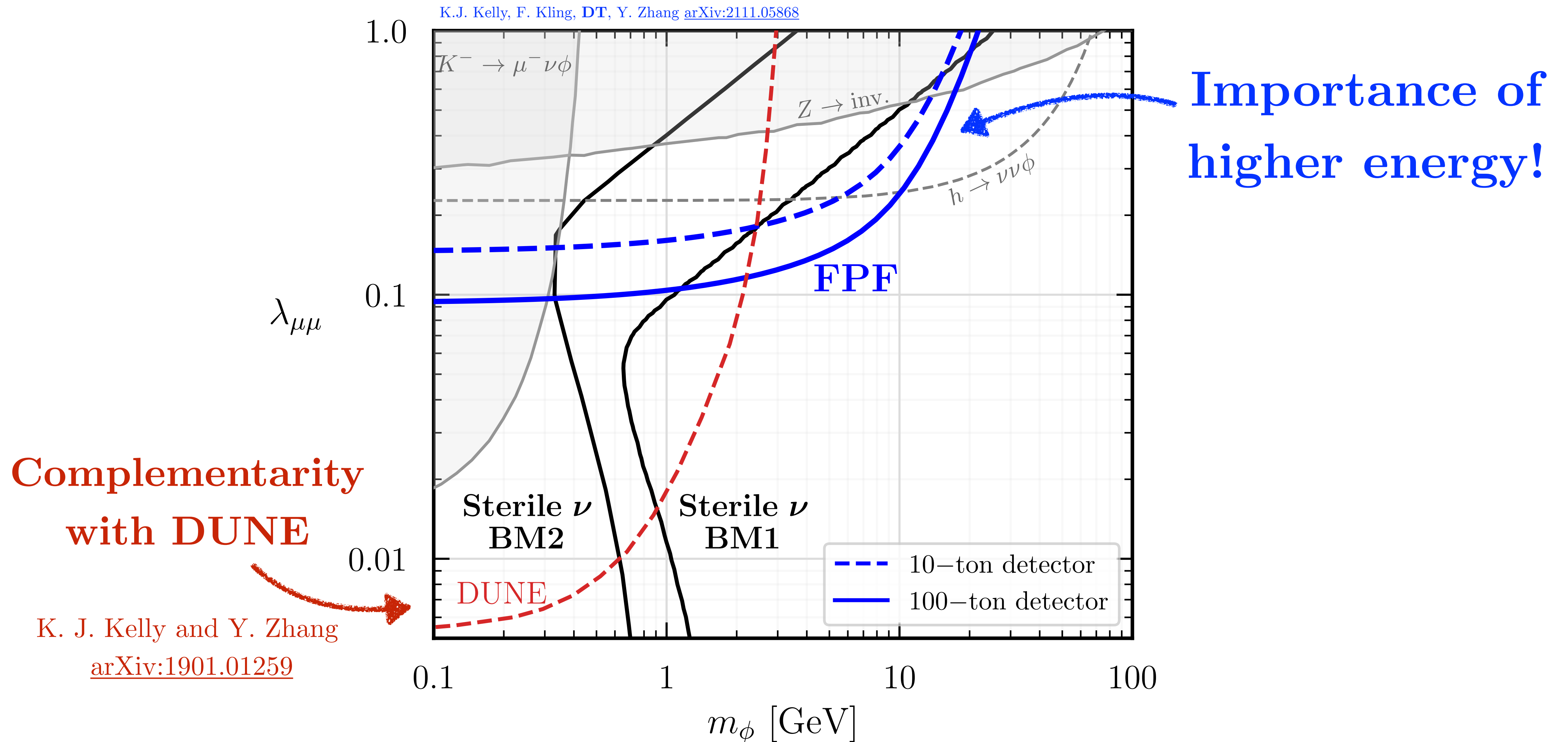


Cut and Count

	$\nu_\mu + \bar{\nu}_\mu \text{ CC}$	$m_\phi = 1 \text{ GeV}$
$E_{vis.} < 600 \text{ GeV}$	61%	76%
$\cancel{p}_T > 3 \text{ GeV}$	0.2%	26%
$p_T^{max} < \frac{4}{3} \cancel{p}_T$	10^{-5}	15%

Significant reduction in bkg. *from missing transverse momentum cut!*

FPF Reach: Sterile Neutrinos

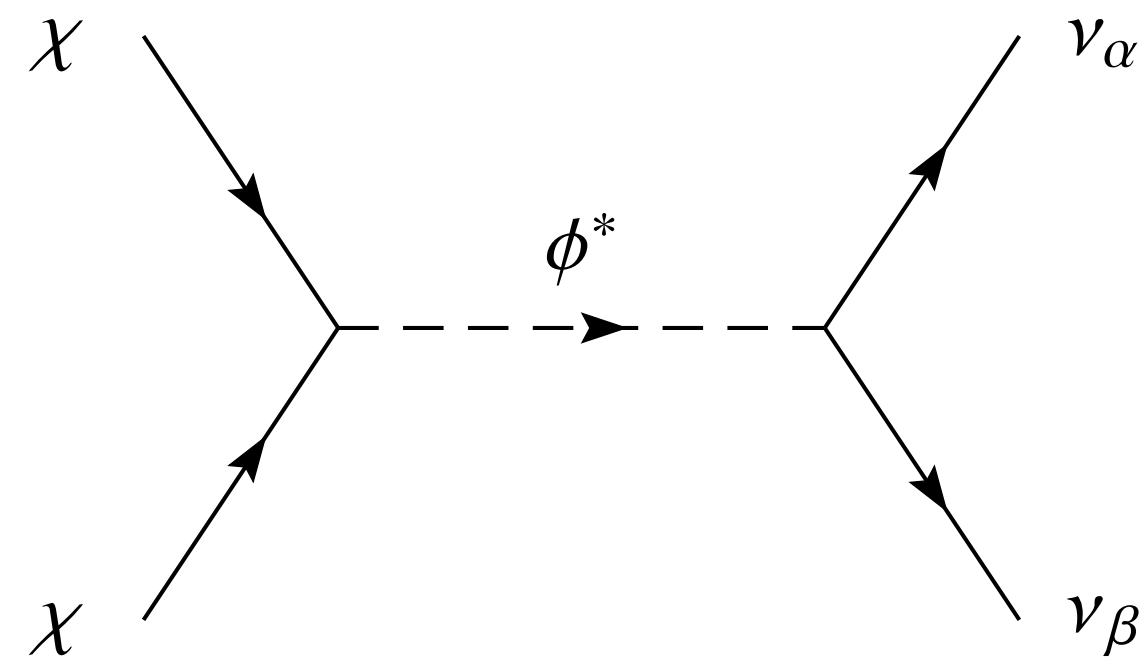


FPF Reach: Thermal DM

- The neutrinophilic scalar ϕ can also be a mediator to thermal DM

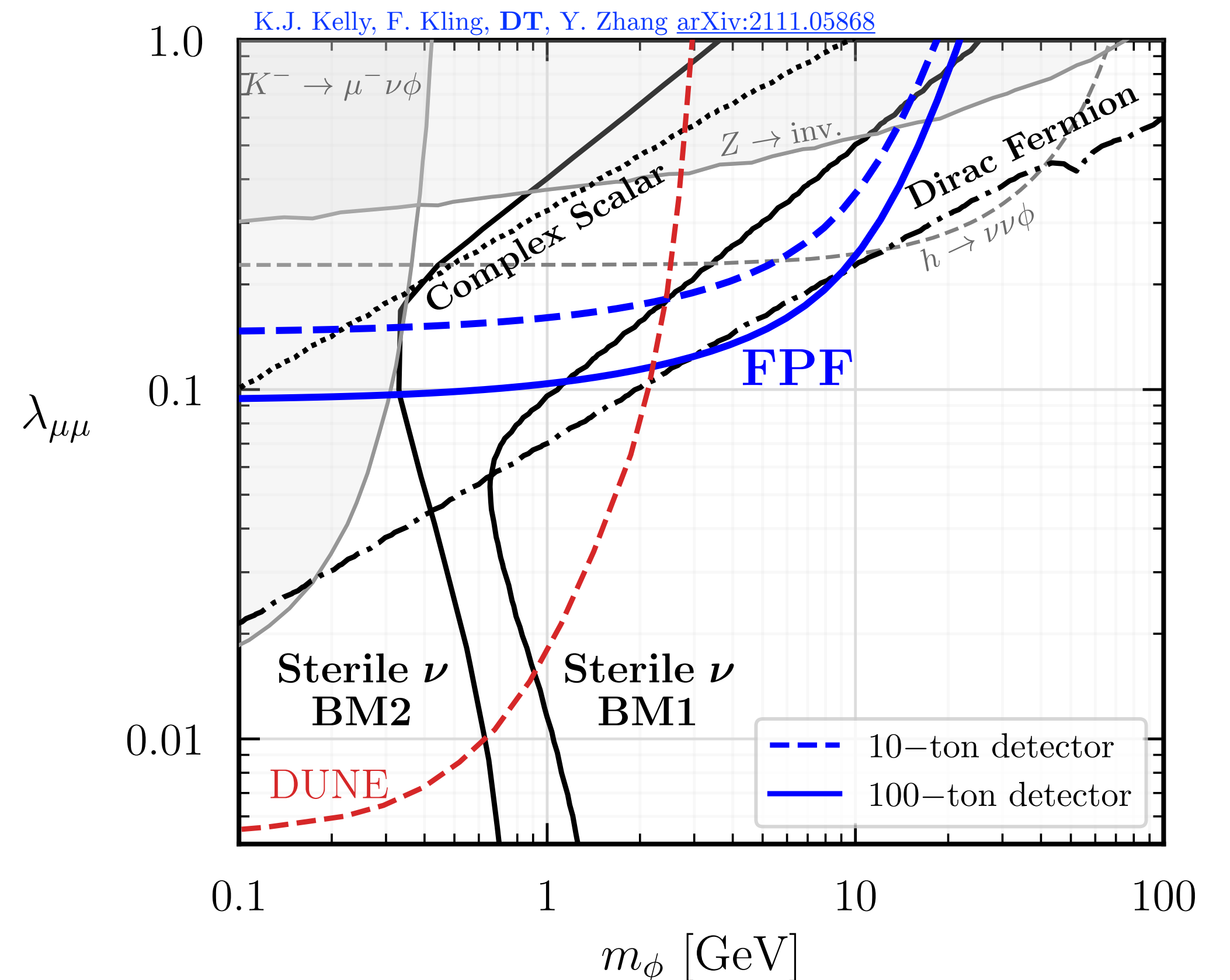
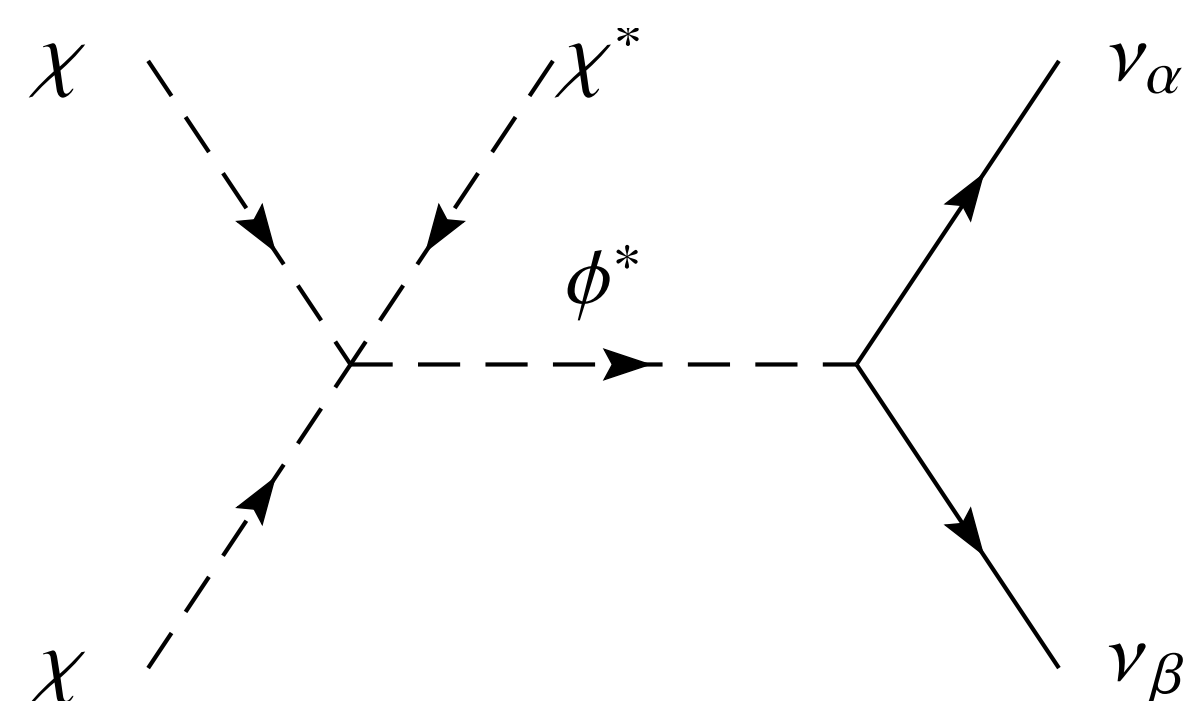
Fermion DM

$$\mathcal{L} = \frac{1}{2} y_\chi \bar{\chi} \chi \phi + \text{h.c.}$$

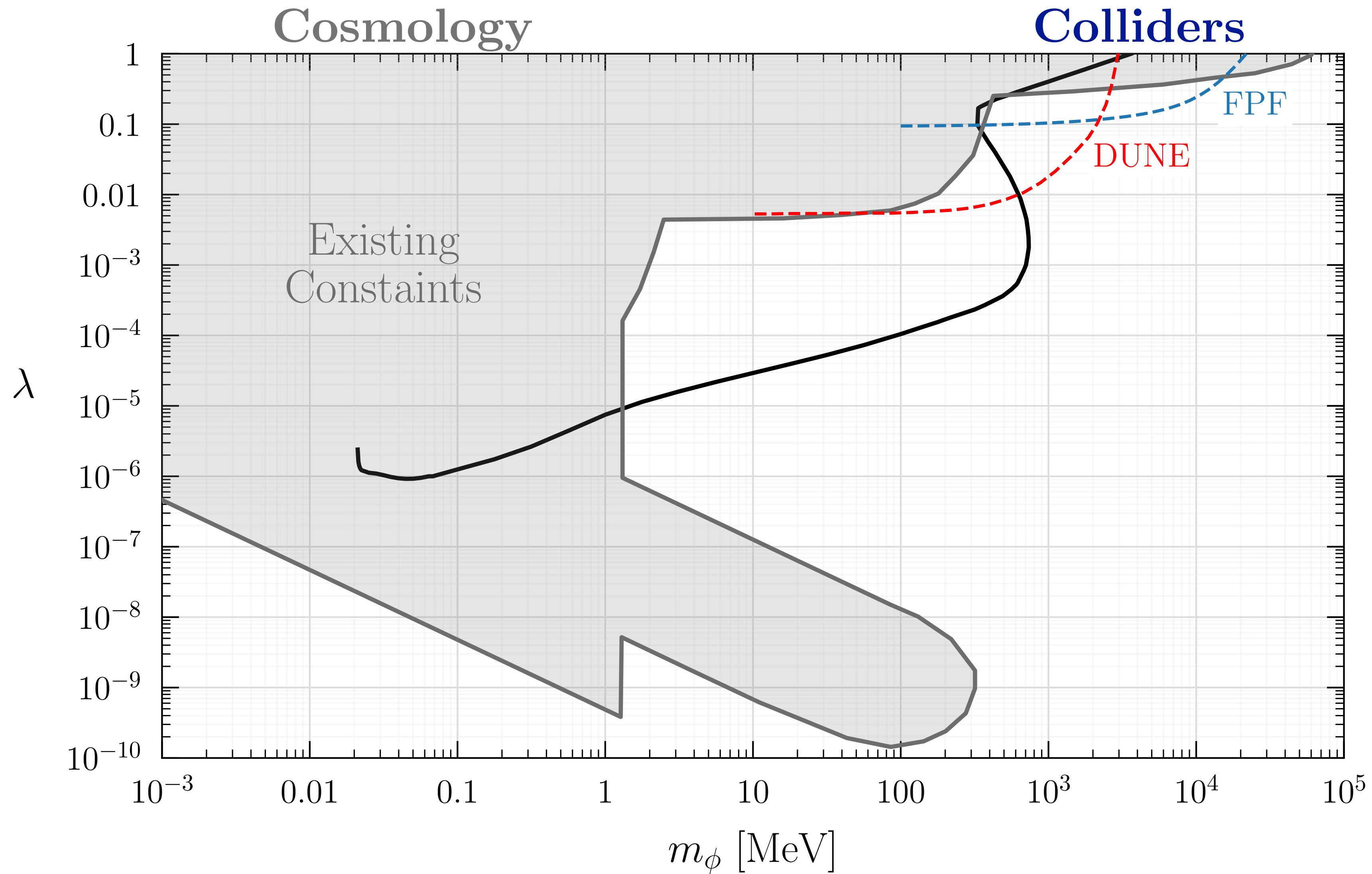


Scalar DM

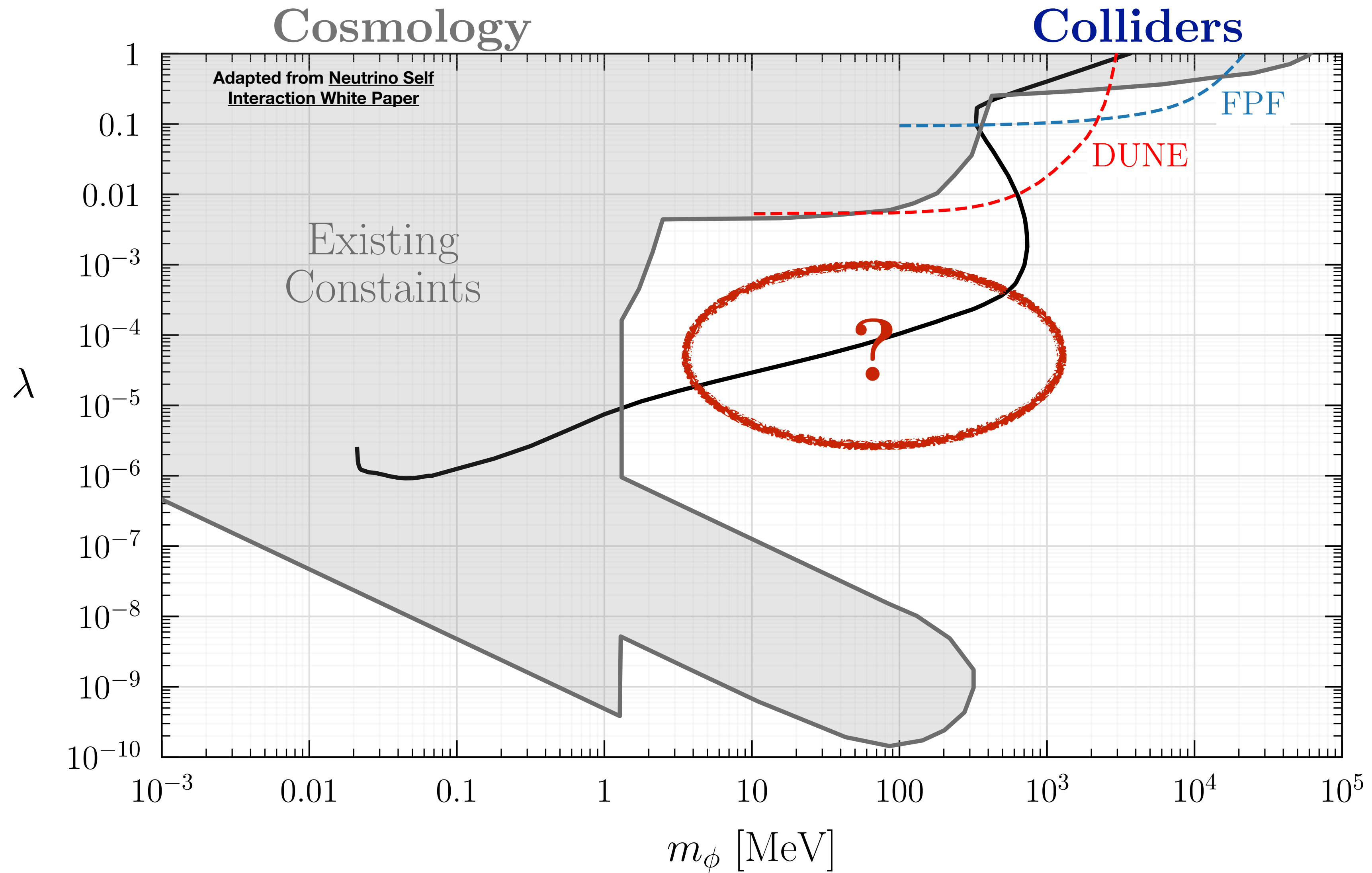
$$\mathcal{L} = \frac{1}{6} y_\chi \chi^3 \phi + \text{h.c.}$$



Big Picture

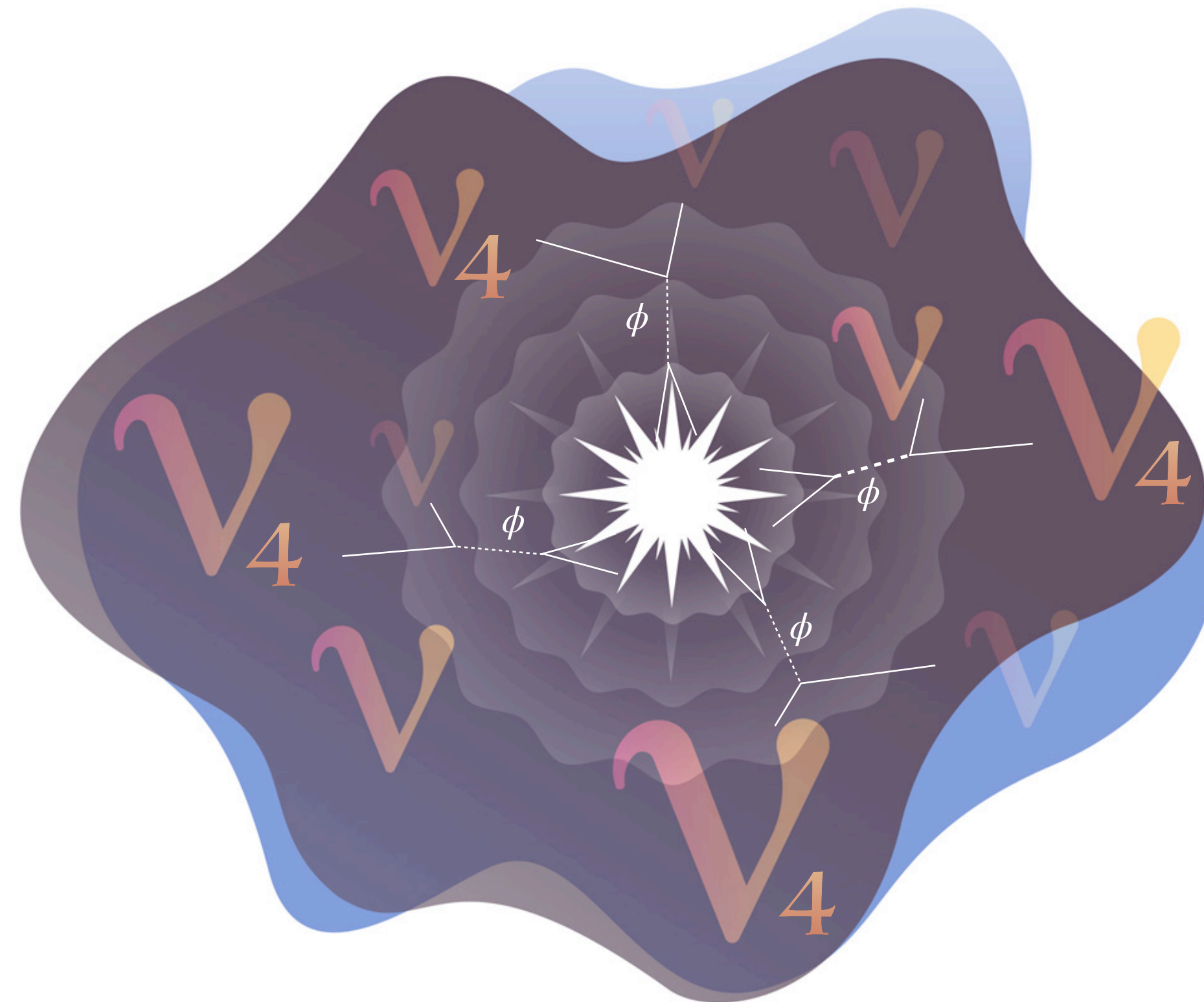


Big Picture



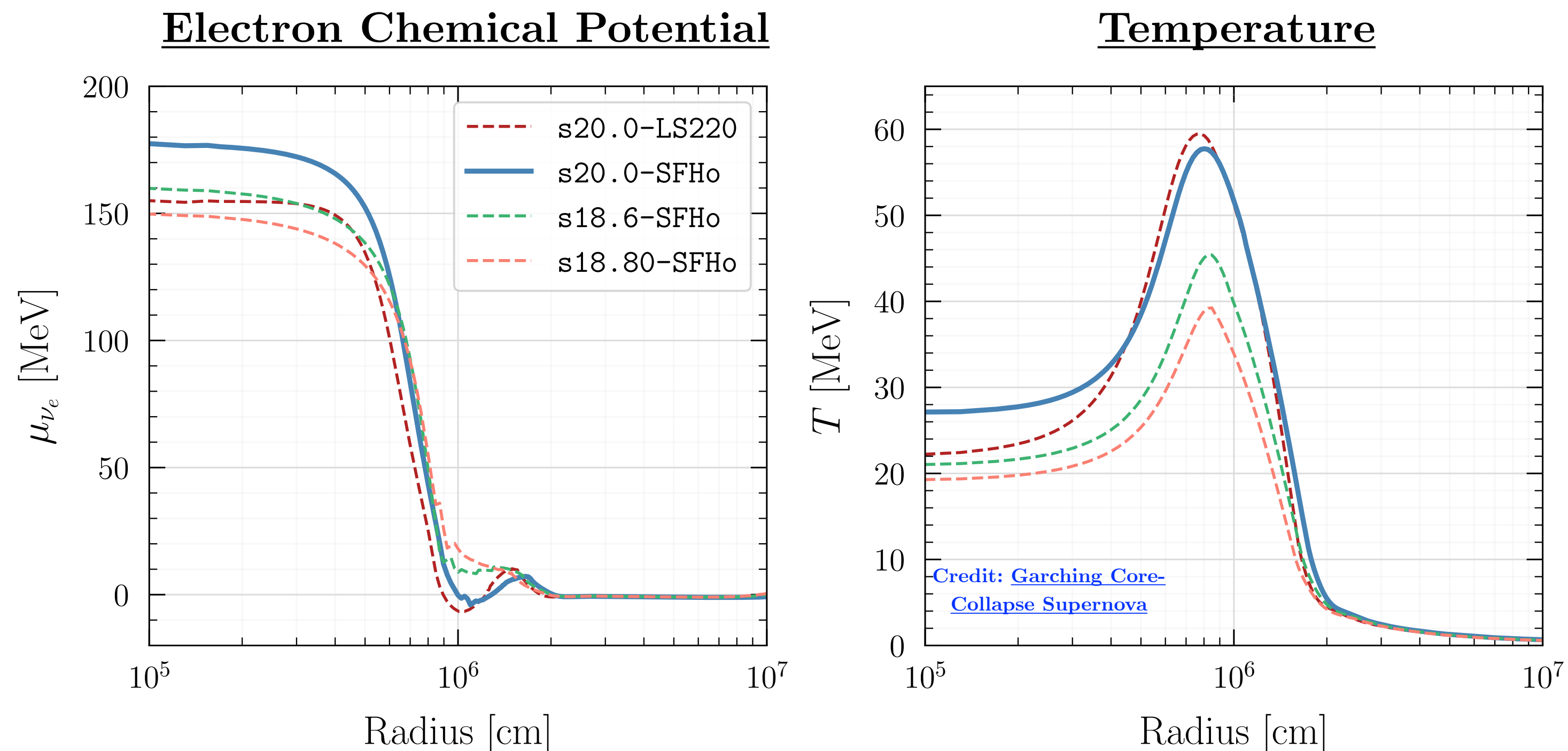
Supernovae

- Another neutrino dense environment!
- Same process that generates $S\nu$ DM relic abundance in early universe produces $S\nu$ DM in the supernova → **excessive supernova cooling!**



Cooling Rate Calculation: A Sketch

- **Step 1: Get supernova profile** $\mu_{\nu}(r), T(r), \rho(r), Y_e(r)$



- $\mu_{\nu_e}/T > 1 \rightarrow$ Fermi-Dirac Distributions are not exponentially suppressed! **Enhanced cooling rate**
 $\mu \neq 0 \rightarrow$ probe smaller couplings!
- $T_{SN} \sim 60$ MeV \rightarrow can probe m_ϕ of 1 MeV up to few 100s of MeV. Exactly where we are missing probes!

Cooling Rate Calculation: A Sketch

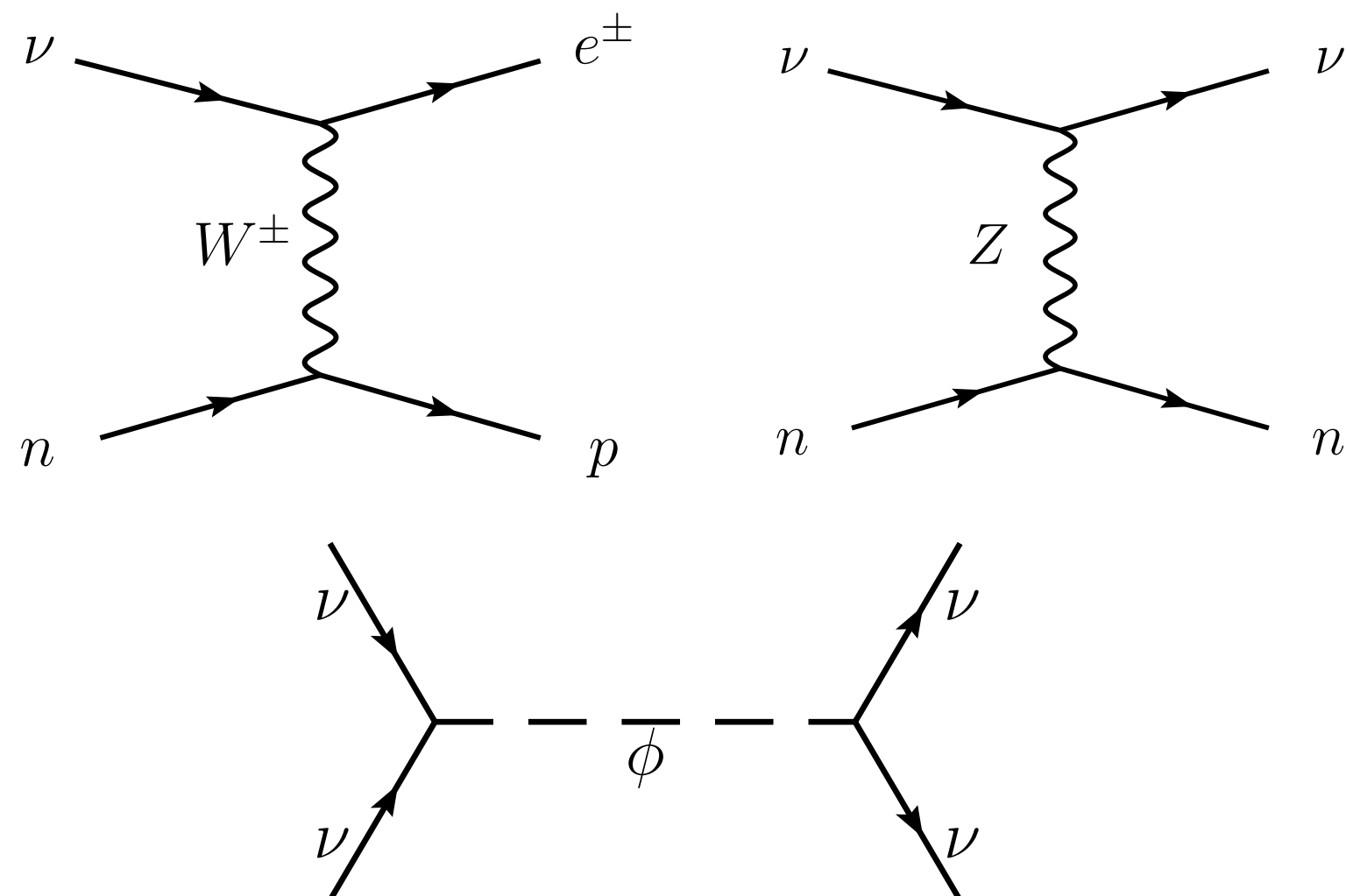
- **Step 2: Calculate active-sterile neutrino mixing in matter**

$$\sin^2(2\theta_{eff}) = \frac{\Delta^2 \sin^2(2\theta)}{\Delta^2 \sin^2(2\theta) + \Gamma^2 + (\Delta \cos(2\theta) - V)^2}$$

← Vacuum
Mixing Angle

Interaction Rate

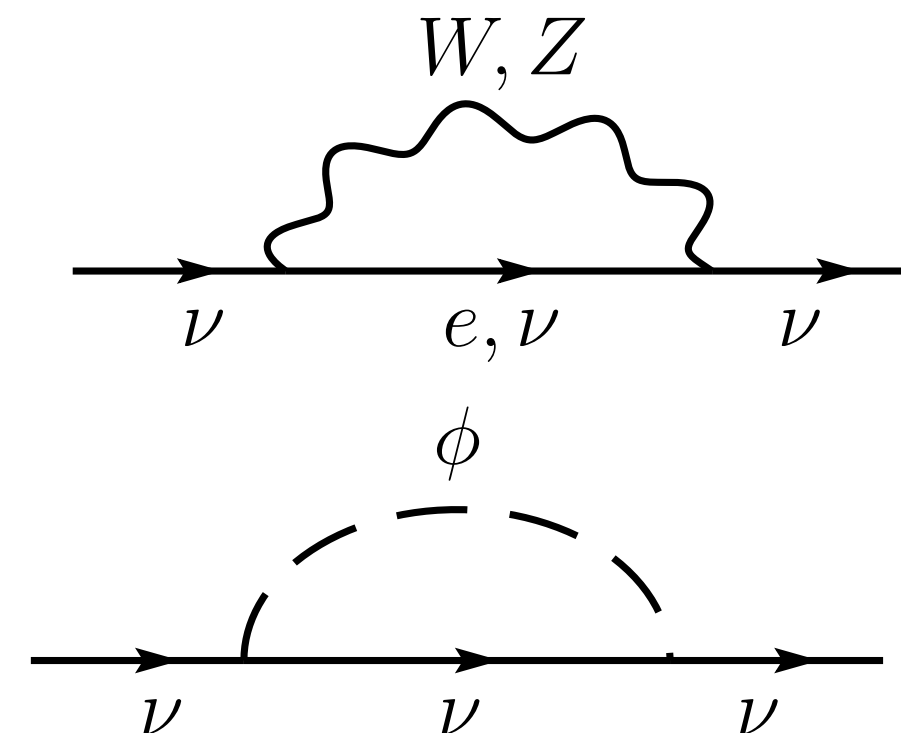
$$\Gamma = \Gamma_{weak} + \Gamma_{\phi}$$



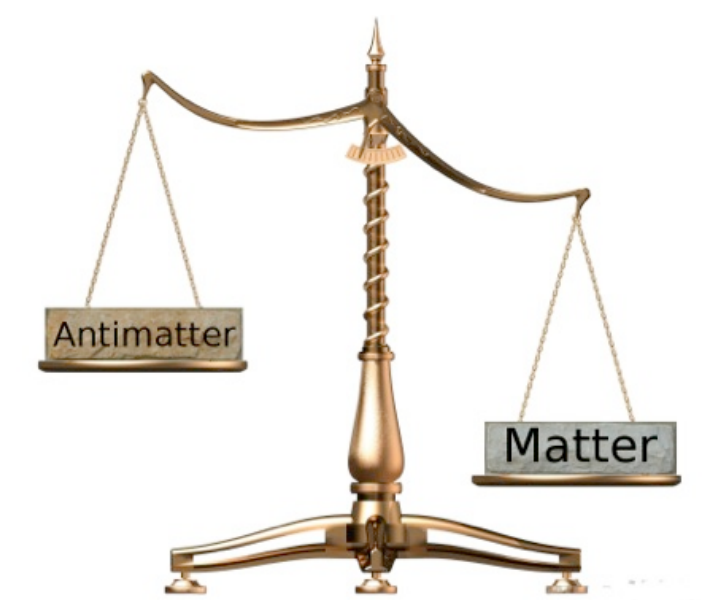
Effective
Potential

$$V = V_{weak} + V_{\phi}$$

Thermal potential



Matter asymmetries

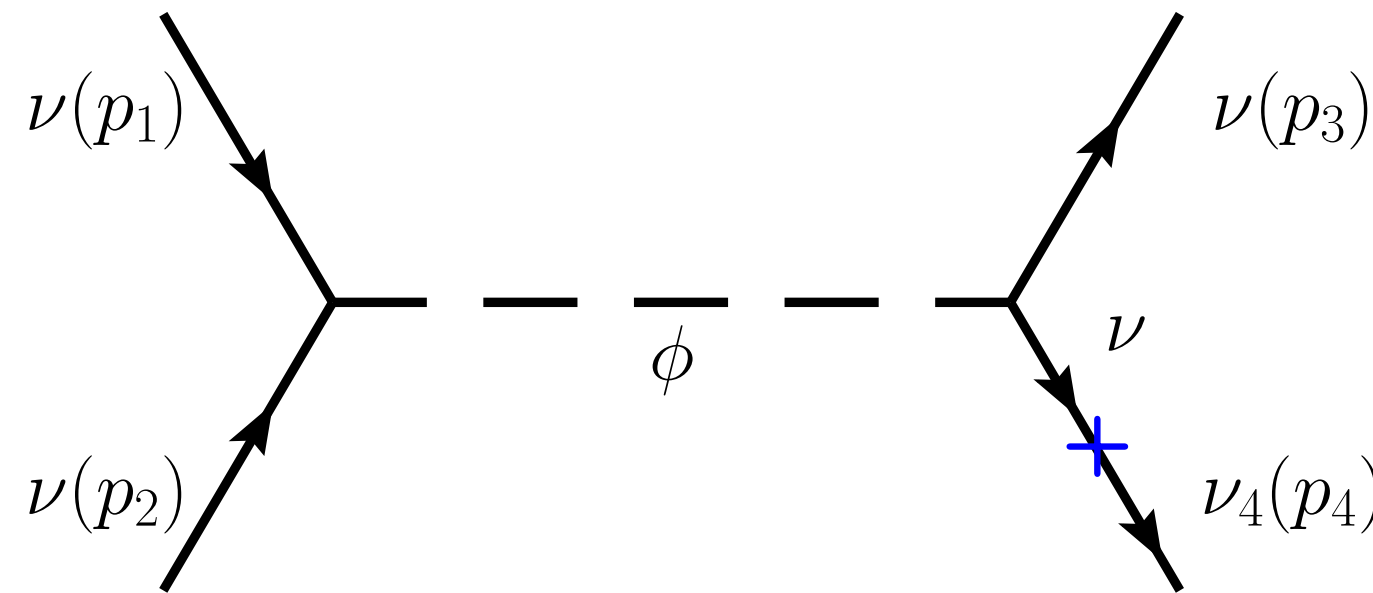


Cooling Rate Calculation: A Sketch

- **Step 3: Optical depth, or ν_4 energy loss due to scattering**

$$\tau = \int_r^\infty dr \sin^2(2\theta_{eff}) \Gamma(E, r) \quad \begin{array}{l} \text{Interaction Rate} \\ \Gamma = \Gamma_{weak} + \Gamma_\phi \end{array}$$

- **Step 4: Sterile neutrino production matrix element**



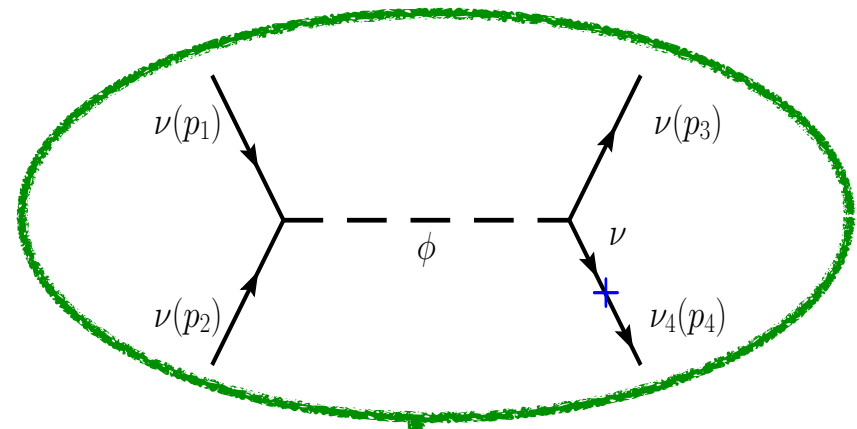
$$|\mathcal{M}|^2 = 32\pi^2 \lambda^2 m_\phi^2 \delta(s - m_\phi^2) \sin^2 \theta_{eff}(r, E_4)$$

- **Step 4.5: Profit**

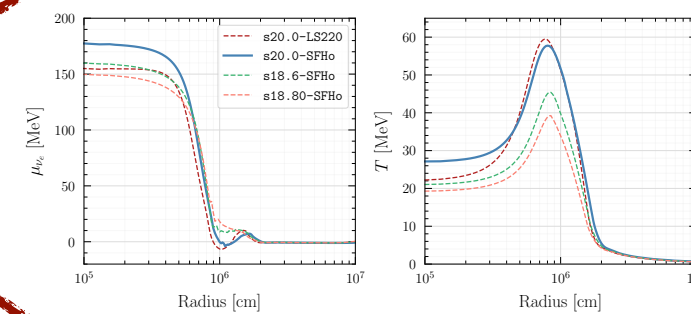
Cooling Rate Calculation: A Sketch

- Step 5: Put everything together to calculate the luminosity

SνDM production



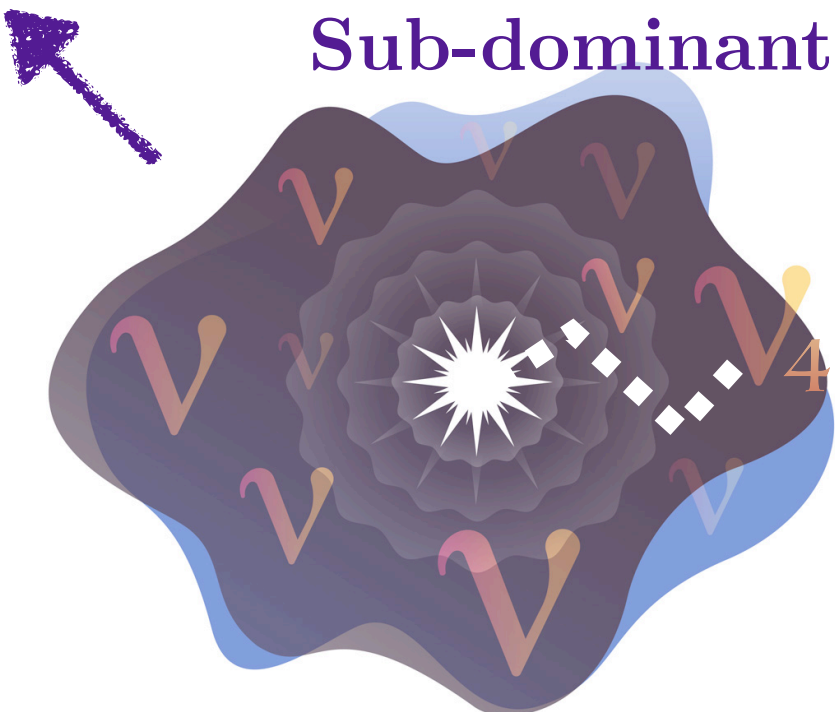
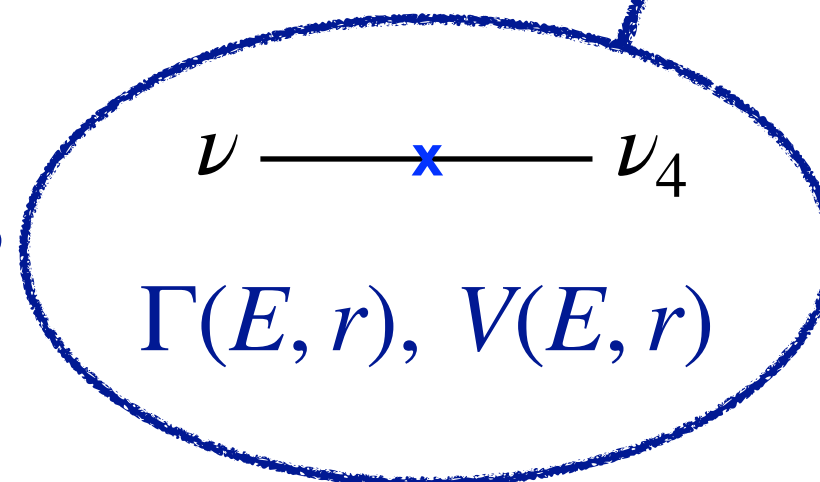
SN Profile



$$L = \frac{\lambda^2 m_\phi^2}{4\pi^2} \int_0^{4R_c} r^2 dr \int_0^\infty dE_1 f(E_1, r) \int_{m_\phi^2/(4E_1)}^\infty dE_2 f(E_2, r) \frac{1}{\sqrt{(E_1 + E_2)^2 - m_\phi^2}}$$

$$\times \int_{\frac{1}{2}(E_1 + E_2 - \sqrt{(E_1 + E_2)^2 - m_\phi^2})}^{\frac{1}{2}(E_1 + E_2 + \sqrt{(E_1 + E_2)^2 - m_\phi^2})} dE_4 \sin^2 \theta_{\text{eff}}(r, E_4) E_4 e^{-\tau(E_4, r)}$$

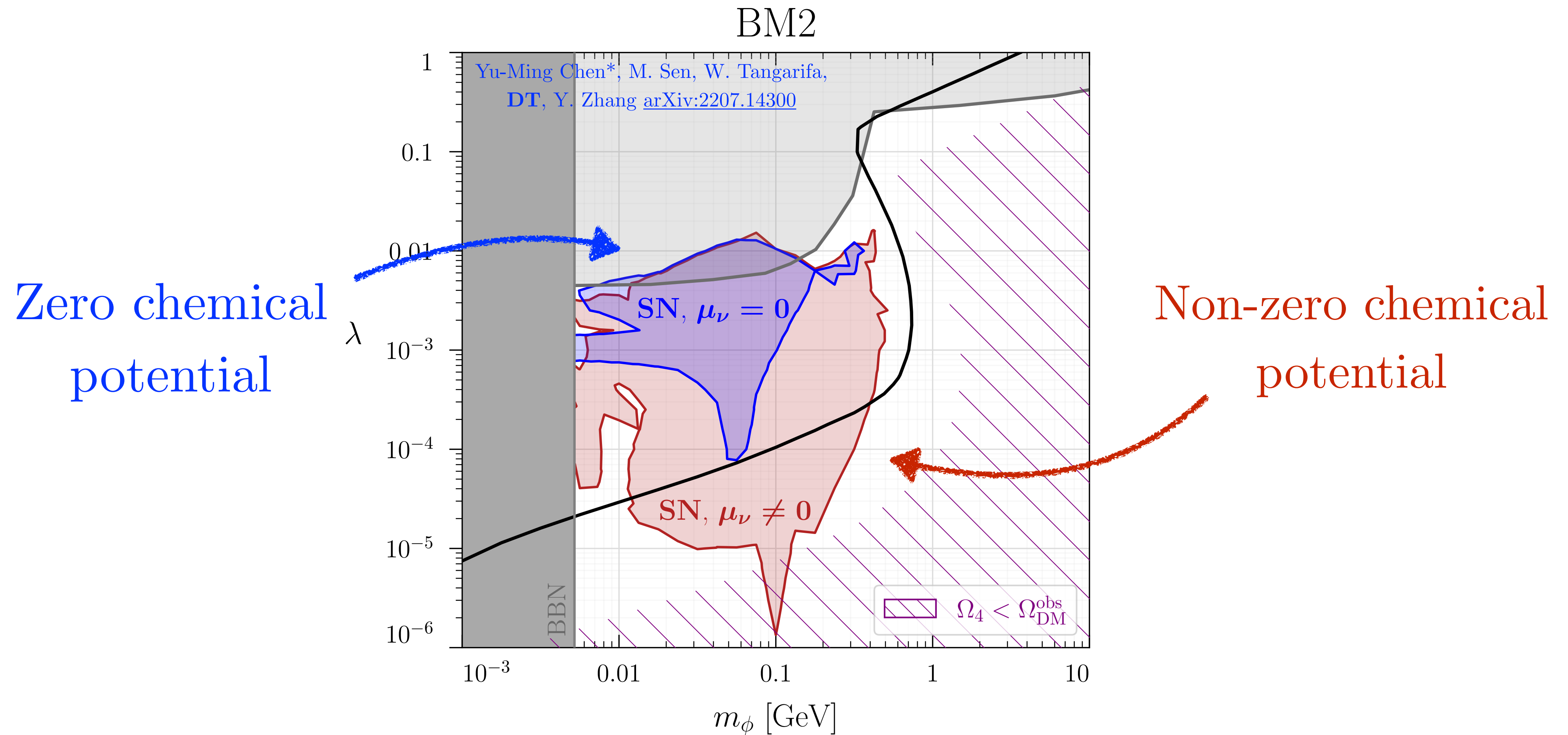
Matter effects



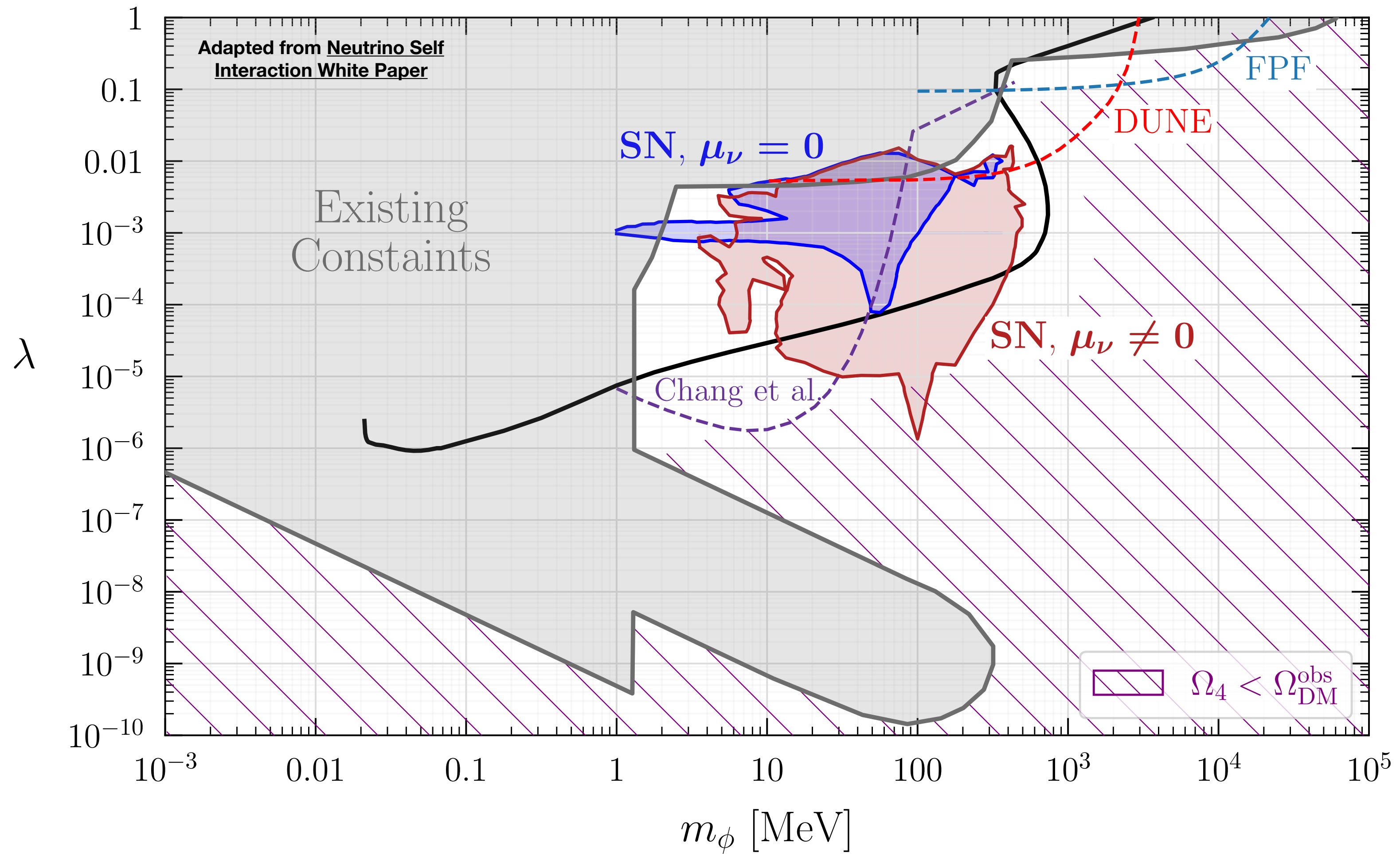
Re-absorption.
Sub-dominant effect

Supernova Cooling Bounds

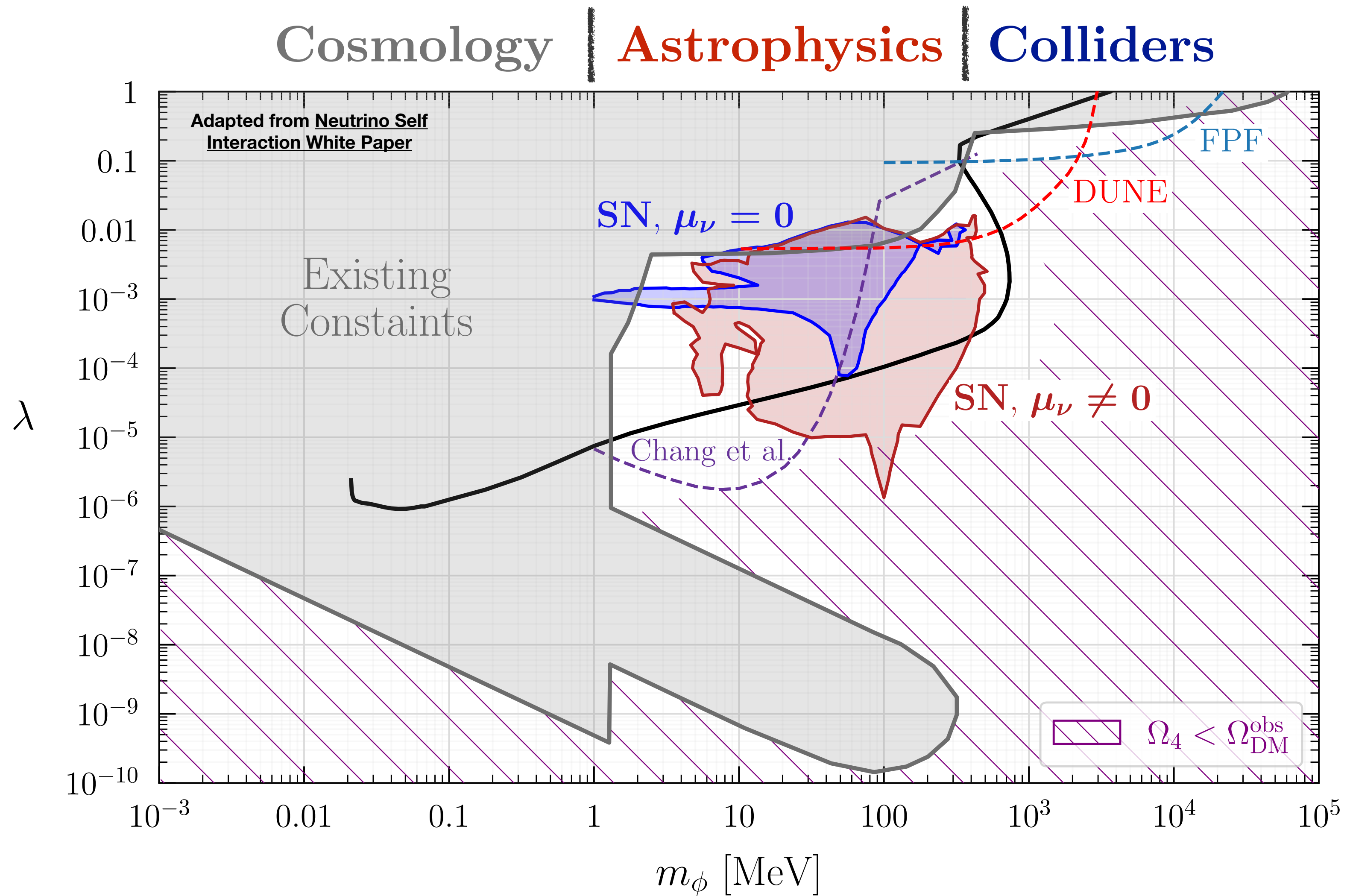
- Observations of SN1987 bound the emission luminosity to be $L \lesssim 3 \times 10^{52}$ ergs/s



Big Picture



Big Picture



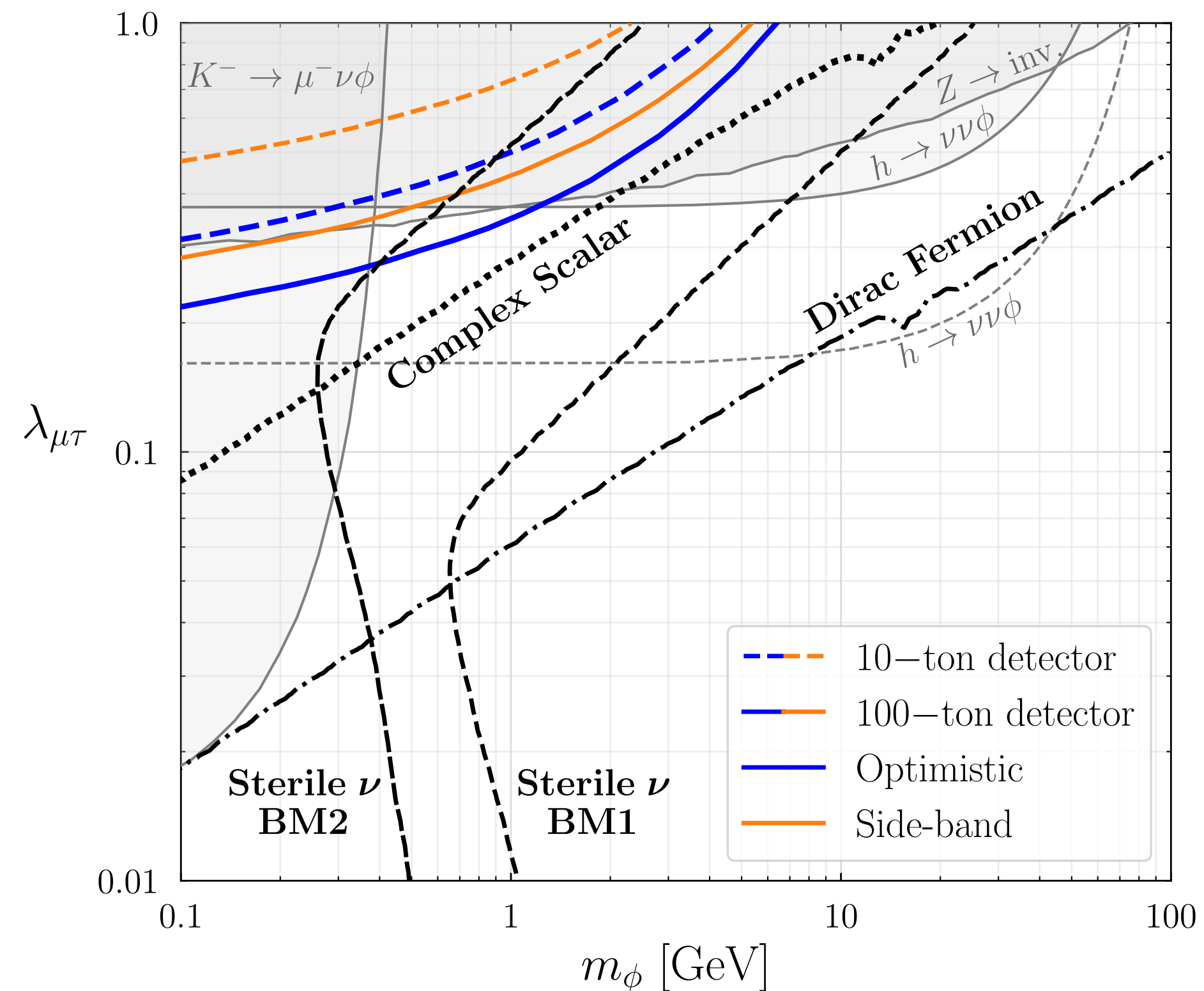
Great complementarity between different probes of neutrinophilic DM!

Thanks!
Questions?

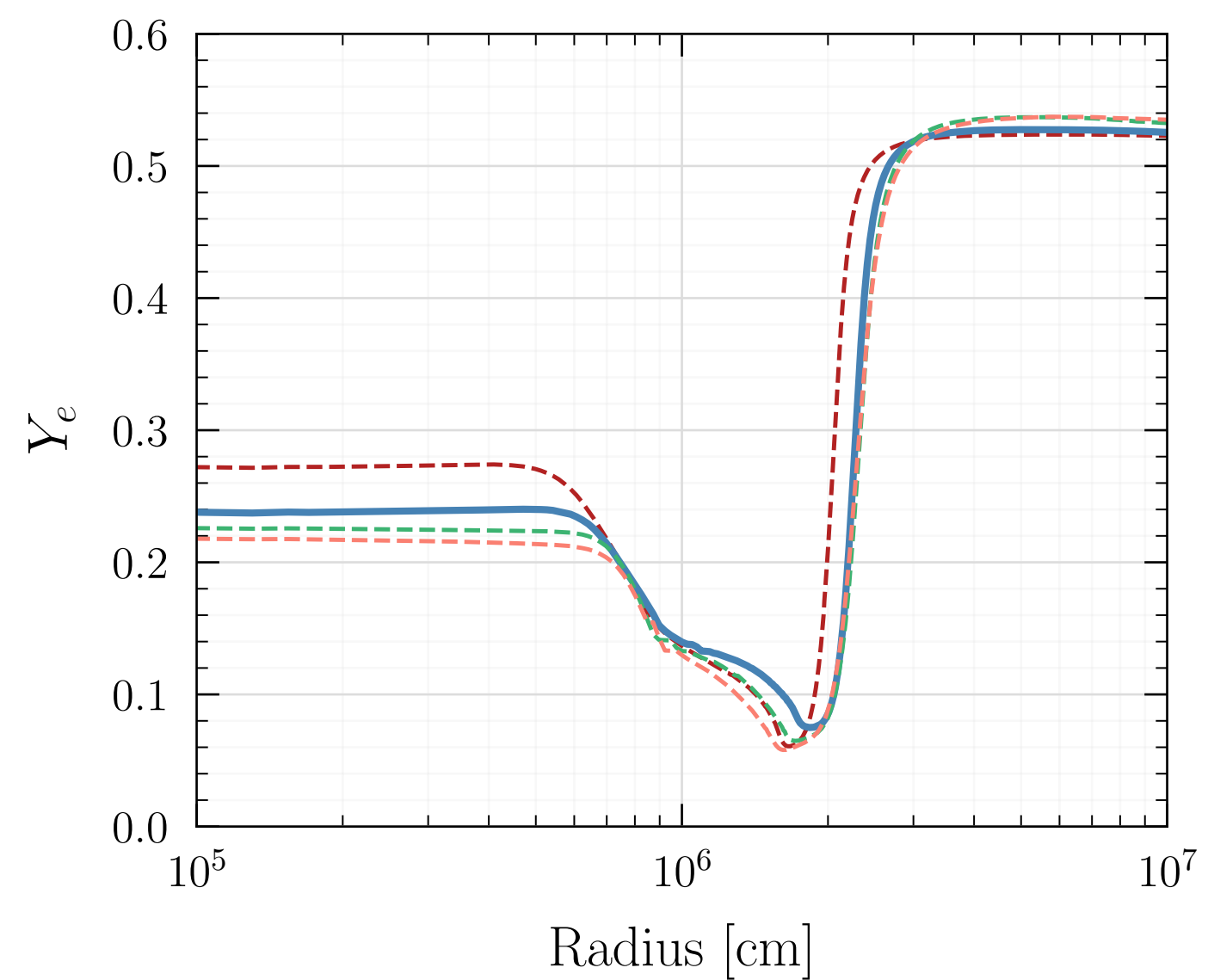
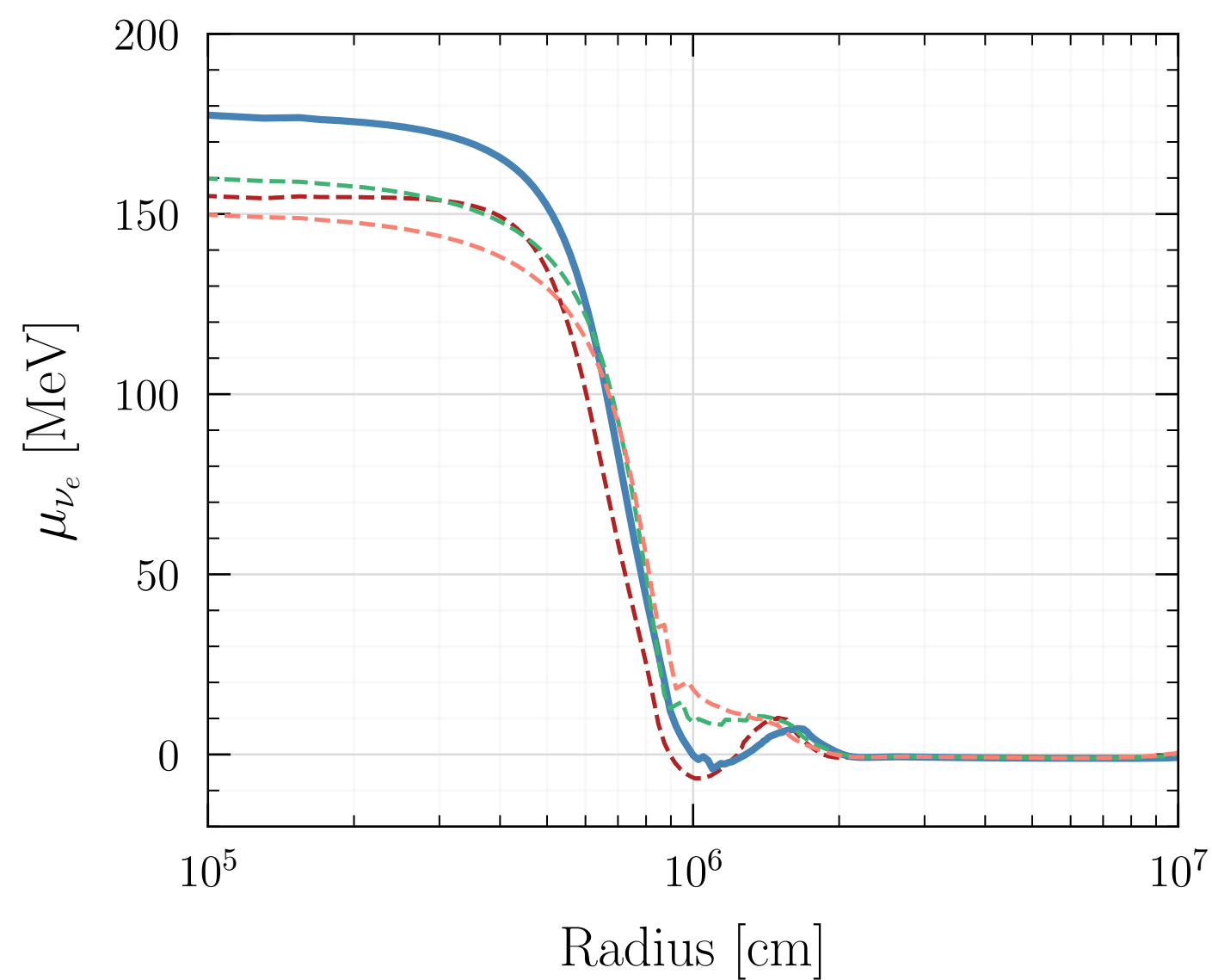
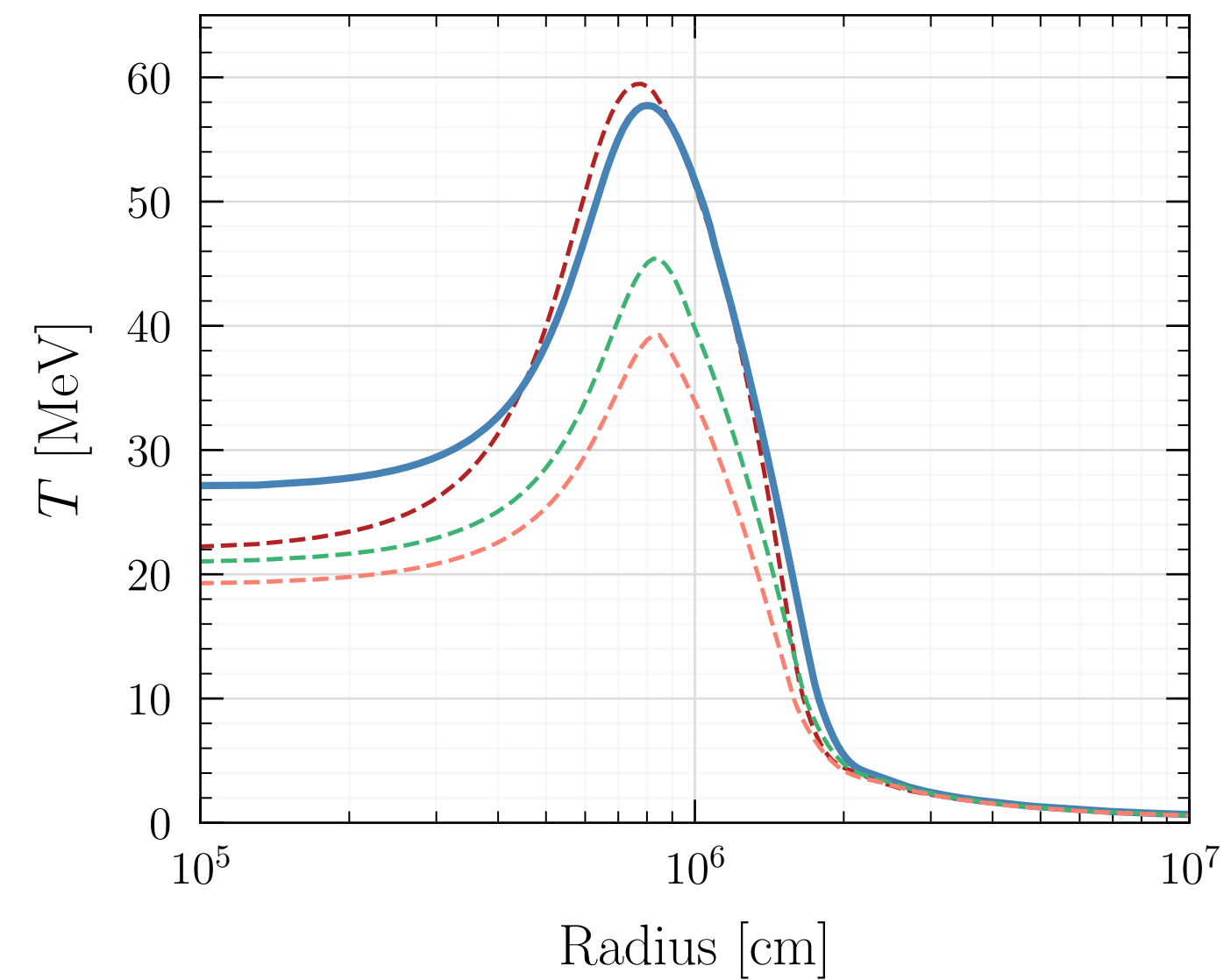
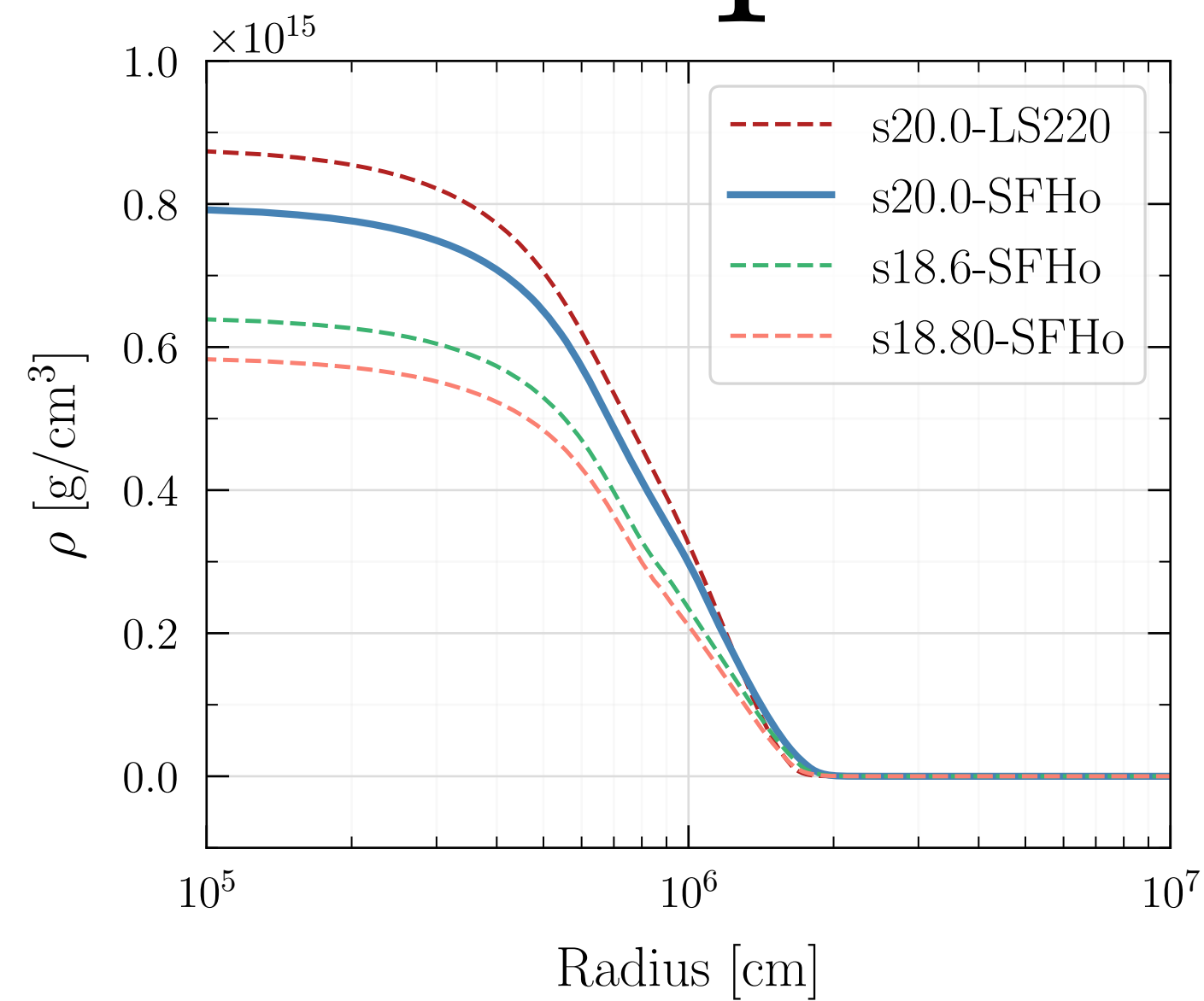
Back up

FPF Reach: Final State Tau Leptons

- For $\lambda_{\mu\tau} \neq 0$, the signal is a tau + \cancel{p}_T coming from a muon-neutrino beam.
- Only $\mathcal{O}(100)$ tau neutrinos are expected to interact with the detector. The signal will result in an excess of tau events compared to the SM.
- Simple analysis: count the number of signal events with a tau in the final state



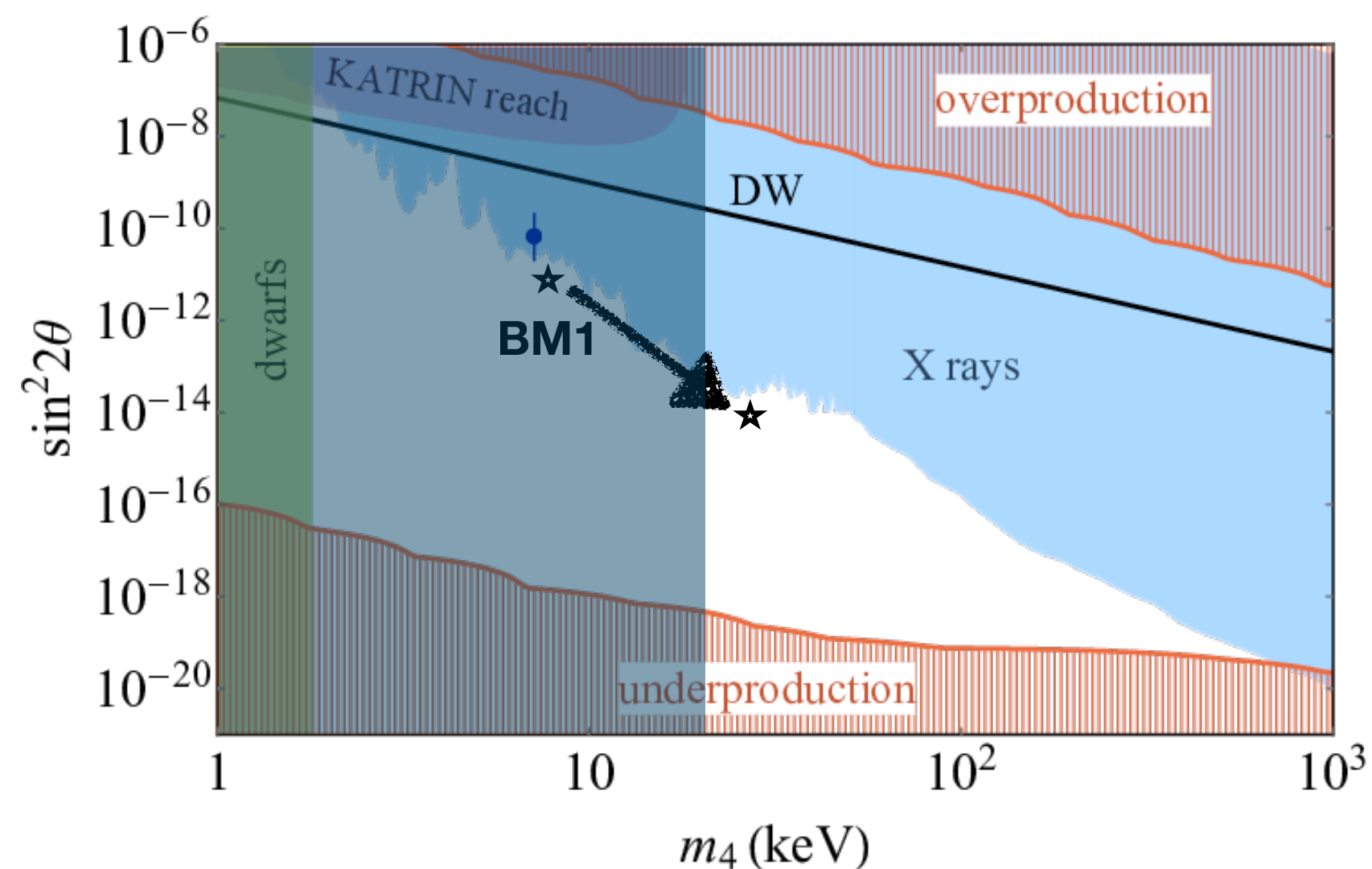
Supernova Profile



Constraints from MW Dwarf Galaxies

- Spoiler alert: There is a lower limit on sterile neutrino dark matter mass in the presence of a neutrinophilic scalar mediator!

$$\Omega \sim \Gamma \times \sin^2(2\theta)$$



Smaller mixing \rightarrow larger λ .
Run into existing constraints.
See [arXiv:2301.08299](https://arxiv.org/abs/2301.08299)

