

## Non-Standard Interactions in Solar Neutrinos

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SOUTH DAKOTA MINES

### Outline

- Non-standard neutrino interactions (NSI)
  - Solar neutrinos, oscillations
- Rates and measurement
- Potential experimental constraints



https://www.businessinsider.com/neutrinos-forged-in-the-heart-of-the-sun-2014-8

• Conclusions



Non-Standard Interactions in Solar Neutrinos

#### Non-standard neutrino interactions (NSI)

- One possible explanation
- Standard model (SM) neutrino interactions
  - Everything else: NSI



- Here focus on neutral-current models with heavy mediators  $(m_{\tau'}^2 \gg Q^2)$ 
  - $\mathcal{L}_{\alpha\beta}^{fP} = -2\sqrt{2}G_F \varepsilon_{\alpha\beta}^{fP} [\bar{\nu}_{\alpha}\gamma^{\mu}(1-\gamma^5)\nu_{\beta}] [\bar{f}\gamma_{\mu}Pf]$
  - Can change v flavor, not fermion (f)
  - Parameterized by ε's



#### <u>https://arxiv.org/pdf/1601.07179.pdf</u> A. Serenelli (2016) <u>https://pdg.lbl.gov/2020/reviews/rpp2020-rev-neutrino-mixing.pdf</u> PDG 2020

#### Solar neutrinos

## High flux of neutrinos < 20 MeV</li> produced in Sun

- Most < 1 MeV
- Produced as  $v_e$ 
  - ~0.5 change flavor



<sup>7</sup>Be

pep

## Solar neutrino oscillations

- Vacuum oscillations + matter effect in Sun
- v<sub>e</sub> survival probability
  - $P_{ee}^{3\nu} = \cos^4 \theta_{13} P_{ee}^{2\nu} + \sin^4 \theta_{13}$
  - $P_{ee}^{2\nu} = \frac{1}{2} [1 + \cos 2\theta \cos 2\theta_m]$





pp

0.9

0.8

0.7

0.6

#### O.G. Miranda, M.A. Tórtola, J.W.F. Valle (2006) arXiv:hep-ph/0406280 (extended with up-quark NSI and different NSI parameterization)

#### NSI in Sun

• 2-flavor model

•  $H = \begin{pmatrix} -\frac{\Delta m^2}{4E} \cos 2\theta + \sqrt{2}G_F N_e & \frac{\Delta m^2}{4E} \sin 2\theta \\ \frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix} + \sqrt{2}G_F (N_d + N_u) \begin{pmatrix} \varepsilon_D & \varepsilon_N \\ \varepsilon_N & -\varepsilon_D \end{pmatrix}$ 

• Measurement: 
$$P(v_e \rightarrow v_e) = \frac{1}{2} [1 + \cos 2\theta \cos 2\theta_m]$$

• 
$$\cos 2\theta_m = \frac{\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e + 2\varepsilon_D(N_d + N_u))}{[\Delta m^2]_{matter}}$$
  
•  $\left[\Delta m^2\right]_{matter}^2 = \left[\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e + 2\varepsilon_D(N_d + N_u))\right]^2$   
+  $\left[\Delta m^2 \sin 2\theta + 4\sqrt{2}\varepsilon_N EG_F(N_d + N_u)\right]^2$ 

Non-Standard Interactions in Solar Neutrinos

#### NSI in Sun

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• Measurement: 
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 NSI<sub>D</sub>  
•  $[\Delta m^2]_{matter}^2 = [\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e + 2\varepsilon_D(N_d + N_u))]^2 + [\Delta m^2 \sin 2\theta + 4\sqrt{2}\varepsilon_N EG_F(N_d + N_u)]^2$ NSI<sub>N</sub>

# NSI constraints (early 2020)

• All but plotted NSI coupling marginalized



#### **Excellent global fit and summary of NSI:**

P. Coloma, I. Esteban, M.C. Gonzalez-Garcia, M. Maltoni (2019) arXiv:1911.09109

07/07/2023

# NSI constraints (early 2020)

- All but plotted NSI coupling marginalized
- Looking primarily at  $\varepsilon^u_{ee}$  and  $\varepsilon^d_{ee}$ 
  - Most promising for non-zero NSI
  - What if we set  $\varepsilon_{ee}^{u} = \varepsilon_{ee}^{d}$ and rest to 0?

#### **Excellent global fit and summary of NSI:**

P. Coloma, I. Esteban, M.C. Gonzalez-Garcia, M. Maltoni (2019) arXiv:1911.09109

07/07/2023



#### Simplified model

O.G. Miranda, M.A. Tórtola, J.W.F. Valle (2006) arXiv:hep-ph/0406280 (extended with up-quark NSI and different NSI parameterization)

• 2-flavor model

$$SM \text{ vacuum oscillations + matter effect} \qquad NSI$$

$$H = \begin{pmatrix} -\frac{\Delta m^2}{4E} \cos 2\theta + \sqrt{2}G_F N_e & \frac{\Delta m^2}{4E} \sin 2\theta \\ \frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix} + \sqrt{2}G_F (N_d + N_u) \begin{pmatrix} \varepsilon_D & 0 \\ 0 & -\varepsilon_D \end{pmatrix}$$

$$Measurement: P(v_e \to v_e) = \frac{1}{2} [1 + \cos 2\theta \cos 2\theta_m] \qquad \varepsilon_D = \frac{\varepsilon_{ee}^u}{2} = \frac{\varepsilon_{ee}^d}{2}$$

$$\cos 2\theta_m = \frac{\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F (N_e + 2\varepsilon_D (N_d + N_u))}{[\Delta m^2]_{matter}} \qquad NSI_D \qquad NSI_N = 0$$

• 
$$\left[\Delta m^2\right]_{matter}^2 = \left[\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F\left(N_e + 2\varepsilon_D(N_d + N_u)\right)\right]^2 + \left[\Delta m^2 \sin 2\theta\right]^2$$

F. Capozzi, S.W. Li, G. Zhu, J.F. Beacom (2018) arXiv:1808.08232

### NSI-modified oscillations

- Correcting KamLAND result (**blue**) with  $\varepsilon_D = 0.0625$ (dot-dashed) produces same survival probability as solar fit (**green**)
- Could possibly explain neutrino anomaly





#### Back to full model

O.G. Miranda, M.A. Tórtola, J.W.F. Valle (2006) arXiv:hep-ph/0406280 (extended with up-quark NSI and different NSI parameterization)

• 2-flavor model

• 
$$H = \begin{pmatrix} -\frac{\Delta m^2}{4E} \cos 2\theta + \sqrt{2}G_F N_e & \frac{\Delta m^2}{4E} \sin 2\theta \\ \frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix} + \sqrt{2}G_F (N_d + N_u) \begin{pmatrix} \varepsilon_D & \varepsilon_N \\ \varepsilon_N & -\varepsilon_D \end{pmatrix}$$

• Measurement: 
$$P(v_e \rightarrow v_e) = \frac{1}{2} [1 + \cos 2\theta \cos 2\theta_m]$$

• 
$$\cos 2\theta_m = \frac{\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e + 2\varepsilon_D(N_d + N_u))}{[\Delta m^2]_{matter}}$$
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•  $[\Delta m^2]_{matter}^2 = [\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e + 2\varepsilon_D(N_d + N_u))]^2 + [\Delta m^2 \sin 2\theta + 4\sqrt{2}\varepsilon_N EG_F(N_d + N_u)]^2$  NSI<sub>N</sub>

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A. Friedland, M.L. Graesser, I.M. Shoemaker, L. Vecchi (2012) arXiv:1111.5331

#### How do NSI affect survival probability?

flavor-diagonal couplings

flavor off-diagonal couplings



#### Current constraints

- Thorough study in paper
  - Global fit of neutrino experiments (2020)
- Marginalized over oscillation parameters
- Green contours 90% and 3σ CL from atmospheric and LBL fit
- Other colors 1σ, 90%, 2σ, 99%, 3σ CL from solar+KamLAND fit
- What can we compare our result to?

I. Esteban, M.C. Gonzalez-Garcia, M. Maltoni, I. Martinez-Soler, J. Salvado (2018) arXiv:1805.04530

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#### Current constraints

- Relevant limit corresponding to same NSI couplings for u and d quarks
  - Our assumptions



I. Esteban, M.C. Gonzalez-Garcia, M. Maltoni, I. Martinez-Soler, J. Salvado (2018) arXiv:1805.04530

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#### Fundamentals of Neutrino Physics and Astrophysics C. Giunti and C.W. Kim (2007)



## Elastic scattering (ES)



#### Rates

 Integrating previous plot 101 for exposure and number of electrons 101 101



## Statistical analysis

- Estimating sensitivity
  - Assume we observe SM prediction
  - What NSI we allow/exclude?
- Use Poisson negative log likelihood

• 
$$NLL = -2\log \mathcal{L} = 2\sum_{i=1}^{N} \left[ R_i^{SM} - R_i^{NSI} + R_i^{SM} \log \frac{R_i^{NSI}}{R_i^{SM}} \right]$$

- Allowed NSI to  $1\sigma,\,90\%$  CL,  $2\sigma,\,99\%$  CL,  $3\sigma$ 
  - *NLL* < 2.30, 4.61, 6.18, 9.21, 11.83
    - Critical values from 2-df  $\chi^2$  distribution



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#### Current-generation Xe DM experiments

- LZ, XENONnT
- Assume
  - 7 t of Xe
  - 3 years of running
  - No backgrounds, perfect energy resolution
- MeVs of deposited energy is high energy



D.S. Akerib et al. (2018), arXiv:1802.06039

#### Potential Xe constraint

- May exclude significant part of parameter space
  - If we detect every neutrino, have perfect energy resolution, etc.
- What is currently allowed?



#### Comparison with current constraints



I. Esteban, M.C. Gonzalez-Garcia, M. Maltoni, I. Martinez-Soler, J. Salvado (2018) arXiv:1805.04530



Global fit is better overall, but combining data may improve fit

B. Abi et al. (2020) arXiv:2002.02967

#### Future neutrino experiments

- Looked into possible constraint in 10s of years
  - Until then, solar neutrinos measured primarily by DM detectors
    - Maybe also SNO+
- Focus on Ar

K. Abe et al. (2018) arXiv:1805.04163

## HUGE DETECTORS





#### Low-background kt-scale Ar detector

- Potential module for Deep Underground Neutrino Experiment
- Assume
  - 3 kt of Ar
  - 1 year of data
    - Early result
  - No backgrounds, perfect energy resolution
  - Energy threshold of 1 MeV



T. Bezerra et al. (2023) arXiv:2301.11878

#### Potential future Ar constraint

- May exclude significant part of parameter space
  - If we detect every neutrino, have perfect energy resolution, etc.



#### Conclusions

- Expect many solar neutrinos to interact in current DM and future neutrino detectors
- Plan to use them to constrain NSI
  - May not significantly improve global fit at first, but will add independent measurement
  - Expect future neutrino experiments to constrain NSI much further
  - Could potentially explain solar neutrino anomaly

ν	Vµ
Ve	?

## Backup slides

A. Gando, et al. (2013) arXiv:1303.4667

#### KamLAND result

- Fit survival probability
  - Measured L/E over expected L/E for vs from reactors around KamLAND

$$\begin{split} P_{ee}^{3\nu} &= \cos^4 \theta_{13} \widetilde{P}_{ee}^{2\nu} + \sin^4 \theta_{13} \\ \widetilde{P}_{ee}^{2\nu} &= 1 - \sin^2 2\theta_{12M} \sin^2 \left(\frac{\Delta m_{21M}^2 L}{4E_{\nu}}\right) \\ \sin^2 2\theta_{12M} &= \frac{\sin^2 2\theta_{12}}{(\cos 2\theta_{12} - A/\Delta m_{21}^2)^2 + \sin^2 2\theta_{12}} \\ dm_{21M}^2 &= \Delta m_{21}^2 \sqrt{(\cos 2\theta_{12} - A/\Delta m_{21}^2)^2 + \sin^2 2\theta_{12}} \\ A &= \pm 2\sqrt{2}G_F \widetilde{N}_e E_{\nu} \qquad \widetilde{N}_e = N_e \cos^2 \theta_{13} \end{split}$$



 $\Delta$ 

## Solar results

- Many solar-neutrino experiments
  - Borexino, Super-K, SNO, ...
- Fit survival probability
  - $P_{ee}^{3\nu} = \cos^4 \theta_{13} P_{ee}^{2\nu} + \sin^4 \theta_{13}$
  - $P_{ee}^{2\nu} = \frac{1}{2} [1 + \cos 2\theta \cos 2\theta_m]$
  - $\cos 2\theta_m = \frac{\Delta m^2 \cos 2\theta 2\sqrt{2}EG_F N_e}{[\Delta m^2]_{matter}}$
  - $\left[\Delta m^2\right]_{matter}^2 = \left[\Delta m^2 \cos 2\theta 2\sqrt{2}EG_F N_e\right]^2 + \left[\Delta m^2 \sin 2\theta\right]^2$



https://indico.fnal.gov/event/43209/contributions/187863/attachments/129474/159089/nakajima\_Neutrino2020.pdf

Y. Nakajima, NEUTRINO2020

31

## Solar neutrino anomaly in 2020

- Started looking into NSI to explain solar neutrino anomaly
- Its significance decreased
   2019 → 2020
- Not game changer
  - still worthwhile
     NSI search with solar v's

SK+SNO fit disfavors the KamLAND best fit value at ~1.4 $\sigma$  (was ~2 $\sigma$ )



https://indico.fnal.gov/event/43209/contributions/187863/attachments/129474/159089/nakajima\_Neutrino2020.pdf

Y. Nakajima, NEUTRINO2020

#### New Super-K solar oscillation results

				$\Delta \chi^2$	8 <sup>3</sup> σ		· · · · · · · · · · · · · · · · · · ·			-		
					<sup>4</sup> 2σ 2			, , , , , , , , , , , , , , , , , , ,				
				eV <sup>2</sup>	_sin <sup>2</sup>	$(\Theta_{12})=0.316^{+0.034}_{-0.026}$ $(\Theta_{12})=0.306\pm0.014$	$\Delta m_{21}^2 = (7.54^{+0.19}_{-0.18})$ $\Delta m_{21}^2 = (6.11^{+1.21}_{-0.68})$	10 <sup>-5</sup> eV <sup>2</sup> sin <sup>2</sup> (Θ 10 <sup>-5</sup> eV <sup>2</sup>	9 <sub>13</sub> )=0.0219±0.0	)014		
		sin²(θ <sub>12</sub> )	Δm <sup>2</sup> <sub>21</sub> [10 <sup>-5</sup> eV <sup>2</sup> ]	10 <sup>-5</sup>	_sin <sup>2</sup>	(Θ <sub>12</sub> )=0.306 <sup>+0.013</sup> -0.012	$\Delta m_{21}^2 = (7.51^{+0.19}_{-0.18})$	10 <sup>-5</sup> eV <sup>2</sup>	relimi	nary		_
	KamLAND	$0.316^{+0.034}_{-0.026}$	$7.54^{+0.19}_{-0.18}$	.⊆_1	5	SK+SN		$\frown$	KamL			_
	SK+SNO	0.306±0.014	$6.11^{+1.21}_{-0.68}$	<sup>2</sup> <sup>2</sup>	_			$\frown$				-
A Fit _	Combined $(-3.6 \pm 1.6)$	$0.306^{+0.013}_{-0.012}$	$7.51^{+0.19}_{-0.18}$	∫ <b>1</b>		O e rech i						
$A_{DN}^{*u} =$ Best f	$= (-3.6 \pm 1.6)$	$(stat) \pm 0.6(syst))\%$ ar $\Delta m_{21}^2$ changed from	$\rightarrow A_{DN}^{1u} = (-2.1 \pm 10^{-5} \text{ eV}^2)$	(2019) to	5	Combin	show 1, 2,	5 σ cor	nfidence i	ntervals	a 2a	3σ
6.1 X <sup>-</sup>	10-5 eV2					0.1	0.2	0.3	0.4	0.5 sin <sup>2</sup> (θ)	24	68 $\Delta\chi^2$

#### Matter oscillations

- Survival probability depends on energy
- Measurements agree with theory
- Best oscillation fit from SNO+SK (black)
- Need better statistics and more measurements in transition region



#### KamLAND

- 1 kton of liquid scintillator
- Located in Kamioka, Japan
  - Detects neutrinos from reactors in Japan



#### arXiv:1808.04207





- Measured both  $v_{e}$  and total v fluxes

#### Super-Kamiokande







O.G. Miranda, M.A. Tórtola, J.W.F. Valle (2006) arXiv:hep-ph/0406280

#### Neutral-current NSI in Sun

2-flavor model

• Measurement: 
$$P(\nu_e \rightarrow \nu_e) = \frac{1}{2} [1 + \cos 2\theta \cos 2\theta_m]$$
  
 $\varepsilon' = \sin^2 \theta_{23} \varepsilon_{\tau\tau}^{dV} - \varepsilon_{ee}^{d}$ 

• 
$$\cos 2\theta_m = \frac{\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e - \varepsilon' N_d)}{[\Delta m^2]_{matter}}$$
  
•  $[\Delta m^2]_{matter}^2 = [\Delta m^2 \cos 2\theta - 2\sqrt{2}EG_F(N_e - \varepsilon' N_d)]^2 + [\Delta m^2 \sin 2\theta + 4\sqrt{2}\varepsilon EG_F N_d]^2$ 

#### How well could DUNE do?

- SNO+SK detected ~80,000 v's
- Assume best-case 15,000 solar v's per 10 kt·year detected in DUNE
  - $E_{\nu}^{mean} = 8 \text{ MeV} (E_{\nu is}^{threshold} = 3 4 \text{ MeV})$
  - Ignore systematics
- Assume SNO+SK uncertainties for 80,000 solar v's in DUNE
  - Scale as  $\sqrt{\nu}$
  - Place at 8 MeV

#### How well could DUNE do?

• ~40 kt·years of DUNE could already validate SNO+SK

exposure	years	statistics (relative to SNO+SK)	uncertainty
10 kt·years	1 (1 module)	0.19	2.3 σ <sub>SNO+SK</sub>
40 kt∙years	1 (4 modules)	0.75	1.2 $\sigma_{SNO+SK}$
160 kt∙years	4 (4 modules)	3	0.58 $\sigma_{SNO+SK}$
400 kt∙years	10 (4 modules)	7.5	0.37 $\sigma_{SNO+SK}$
1,600 kt·years	40 (4 modules)	30	0.18 σ <sub>SNO+SK</sub>

## Statistical analysis for DUNE

- Estimating sensitivity
  - Assume we observe SM prediction
  - What NSI we allow/exclude?
- Use this negative log likelihood

• 
$$NLL = -2 \log \mathcal{L} = \sum_{i=1}^{N} \frac{\left(P_i^{SM} - P_i^{NSI}\right)^2}{\sigma_i^2}$$

- Allowed NSI to  $1\sigma,\,90\%$  CL,  $2\sigma,\,99\%$  CL,  $3\sigma$ 
  - *NLL* < 2.30, 4.61, 6.18, 9.21, 11.83
    - Critical values from 2-df  $\chi^2$  distribution

#### Differential ES cross section

 Taken from Fundamentals of Neutrino Physics and Astrophysics by C. Giunti and C.W. Kim

$$\frac{\mathrm{d}\sigma}{\mathrm{d}T_e}(E_\nu, T_e) = \frac{\sigma_0}{m_e} \left[ g_1^2 + g_2^2 \left( 1 - \frac{T_e}{E_\nu} \right)^2 - g_1 g_2 \frac{m_e T_e}{E_\nu^2} \right] \qquad \qquad T_e^{\max}(E_\nu) = \frac{2 E_\nu^2}{m_e + 2 E_\nu}$$

$$\sigma_0 = \frac{2 \, G_{\rm F}^2 \, m_e^2}{\pi} \simeq 88.06 \times 10^{-46} \, {\rm cm}^2 \qquad \qquad m_e = 0.511 \, {\rm MeV}$$
$$g_1^{(\nu_e)} = g_2^{(\bar{\nu}_e)} = 1 + \frac{g_V^l + g_A^l}{2} = 1 + g_L^l = \frac{1}{2} + \sin^2 \vartheta_{\rm W} \simeq 0.73$$

$$g_2^{(\nu_e)} = g_1^{(\bar{\nu}_e)} = \frac{g_V^l - g_A^l}{2} = g_R^l = \sin^2 \vartheta_W \simeq 0.23$$

#### Elastic scattering of solar neutrinos in DUNE

- Below  $E_{\nu_e} = 7$  MeV will have more ES interactions in Ar
- Close to threshold, but...





CC and 18·ES cross sections cross at ~7 MeV



#### 07/07/2023



#### 07/07/2023

I. Esteban, M.C. Gonzalez-Garcia, M. Maltoni, I. Martinez-Soler, J. Salvado (2018) arXiv:1805.04530

## Parameter definitions

- Matter Hamiltonian  $H_{\text{mat}}^{\text{eff}} = \sqrt{2}G_F N_e(x) \begin{bmatrix} c_{13}^2 & 0\\ 0 & 0 \end{bmatrix} + \begin{bmatrix} \xi^p + Y_n(x)\xi^n \end{bmatrix} \begin{pmatrix} -\varepsilon_D^\eta & \varepsilon_N^\eta\\ \varepsilon_N^{\eta*} & \varepsilon_D^\eta \end{pmatrix} \end{bmatrix}$
- Diagonal and non-diagonal NSI

$$\begin{aligned} \varepsilon_D^{\eta} &= c_{13} s_{13} \operatorname{Re} \left( s_{23} \varepsilon_{e\mu}^{\eta} + c_{23} \varepsilon_{e\tau}^{\eta} \right) - \left( 1 + s_{13}^2 \right) c_{23} s_{23} \operatorname{Re} \left( \varepsilon_{\mu\tau}^{\eta} \right) \\ &- \frac{c_{13}^2}{2} \left( \varepsilon_{ee}^{\eta} - \varepsilon_{\mu\mu}^{\eta} \right) + \frac{s_{23}^2 - s_{13}^2 c_{23}^2}{2} \left( \varepsilon_{\tau\tau}^{\eta} - \varepsilon_{\mu\mu}^{\eta} \right) \\ \varepsilon_N^{\eta} &= c_{13} \left( c_{23} \varepsilon_{e\mu}^{\eta} - s_{23} \varepsilon_{e\tau}^{\eta} \right) + s_{13} \left[ s_{23}^2 \varepsilon_{\mu\tau}^{\eta} - c_{23}^2 \varepsilon_{\mu\tau}^{\eta*} + c_{23} s_{23} \left( \varepsilon_{\tau\tau}^{\eta} - \varepsilon_{\mu\mu}^{\eta} \right) \right] \end{aligned}$$

 Assumed usual NSI couplings can be factorized into neutrino and charged-fermion parts

$$\varepsilon_{\alpha\beta}^{f} \equiv \varepsilon_{\alpha\beta}^{f,L} + \varepsilon_{\alpha\beta}^{f,R} = \varepsilon_{\alpha\beta}^{\eta} \xi^{f} \qquad \xi^{p} = \sqrt{5} \cos \eta$$
$$\xi^{n} = \sqrt{5} \sin \eta$$

#### Neutrinos in LZ

- <sup>8</sup>B neutrinos via CEvNS
  - CEvNS (nuclear) recoils look exactly like WIMP recoils ( $m_{\chi} \approx 6 \text{ GeV}$ )
  - Expect few events
  - Have not been observed yet
    - important measurement
- Many other neutrinos will interact in LZ
  - What can we learn from them?

B solar vs

S1c [phd]

4.5

 $^{8}B + hep$ 

10

.og10(S2c [phd])

48

#### What can LZ do?

- Not seeing reactor neutrinos
  - Not many reactors around
  - Detector not big enough
  - Cannot improve KamLAND measurement
- Expect to see many solar neutrinos
  - More matter effect from high solar density
  - Interested to pursue this



### Deep Underground Neutrino Experiment

- 40 kt of liquid argon
  - Staged
- First couple years assume 40 kt-years
- Solar neutrinos via CC
- Possible module with low-threshold (~100 keV) 05.0 <sup>G</sup> 1 kt of liquid argon
  - Solar neutrinos via ES
- E. Church, C. Jackson, R. Saldanha (2020) arXiv:2005.04824





#### Not strong early constraint due to high threshold

2.5

5.0

7.5 10.0 12.5 15.0 17.5 20.0

neutrino energy (MeV)

0.35

0.25

දි 0.20 Ę 0.15

0.10 0.05 0.00 ∔ 0.0

#### Hyper-Kamiokande

- 217 kt of water
  - 2<sup>nd</sup> module may be added
- Assume 4 years of statistics
- Solar neutrinos via ES
  - Recent Super-K analysis with 3.5-MeV threshold

Very strong constraint due to high statistics (no systematics or backgrounds here)

