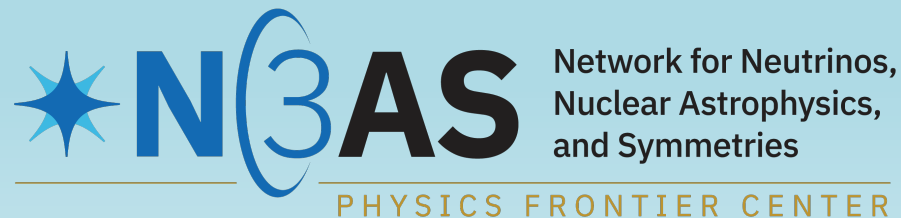
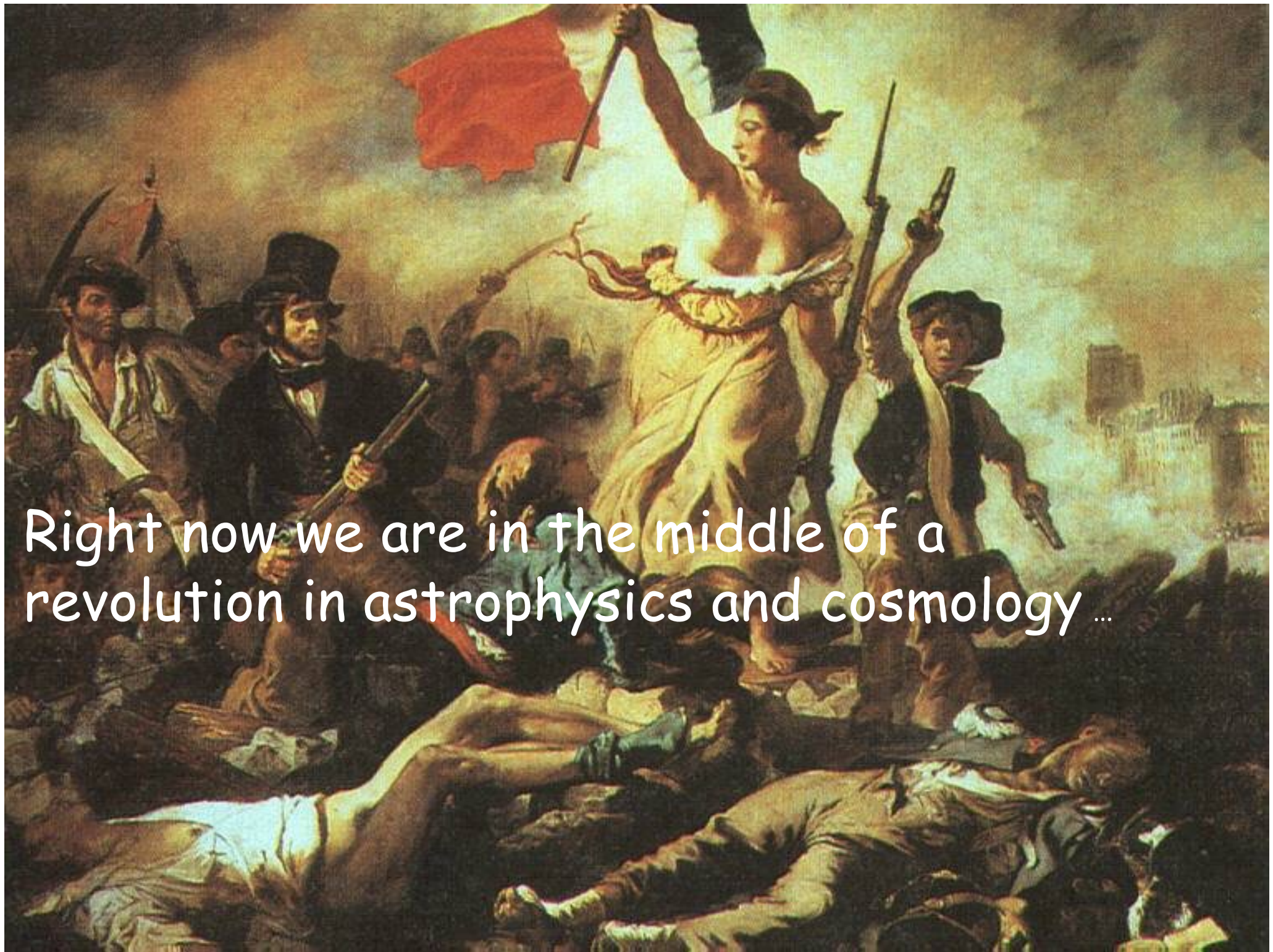


# An Overview of Physics at Underground Laboratories

A.B. Balantekin

Conference on Science at the Sanford  
Underground Research Facility 2022



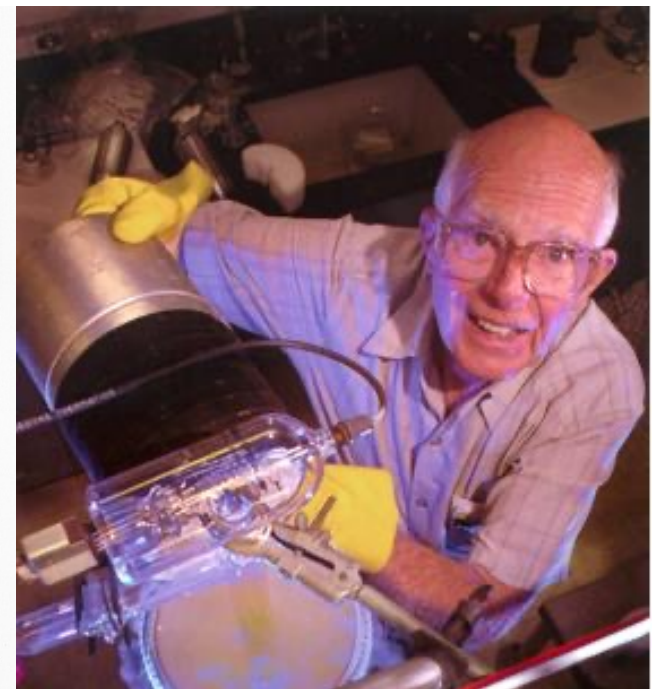
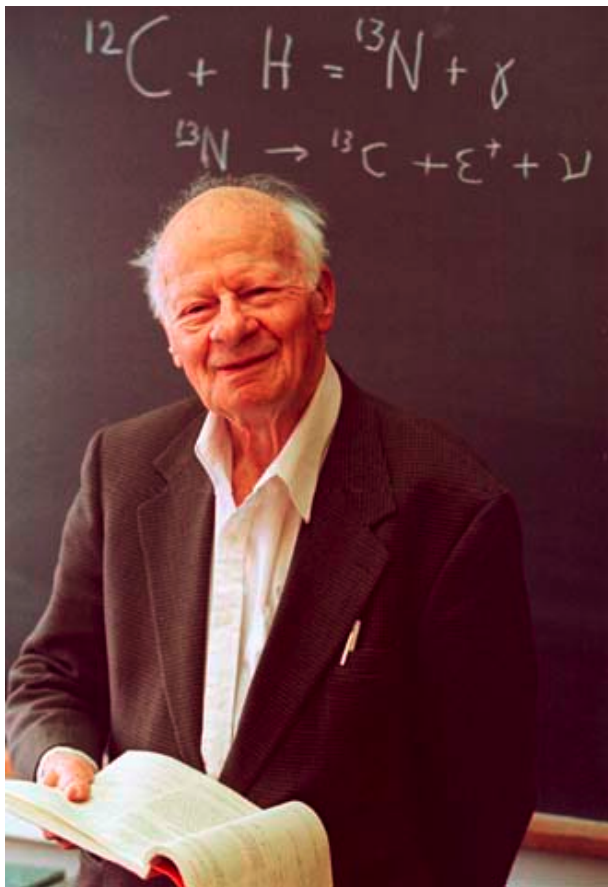


Right now we are in the middle of a revolution in astrophysics and cosmology ...



Right now we are in the middle of a revolution in astrophysics and cosmology ..

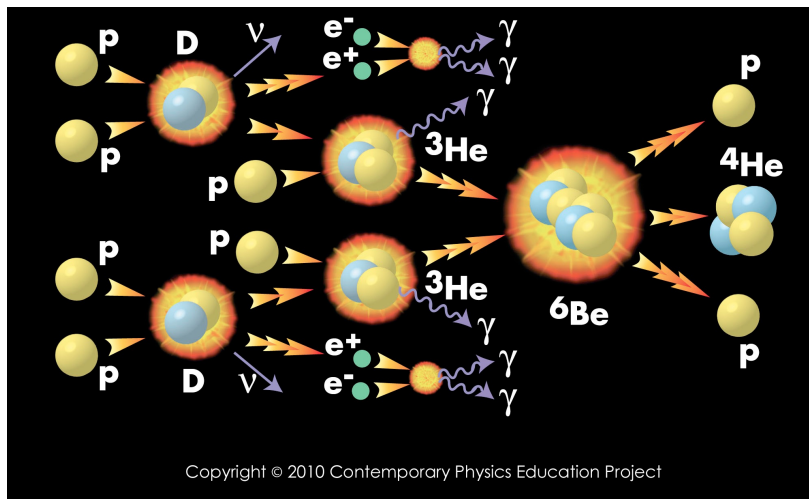
..which started not too far from here in an underground laboratory.



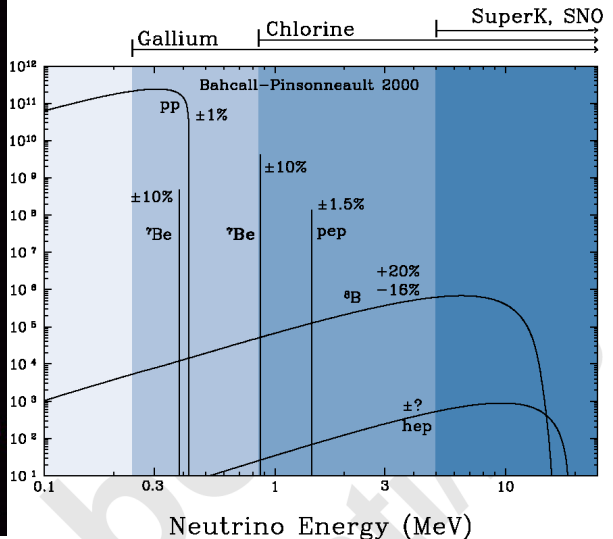
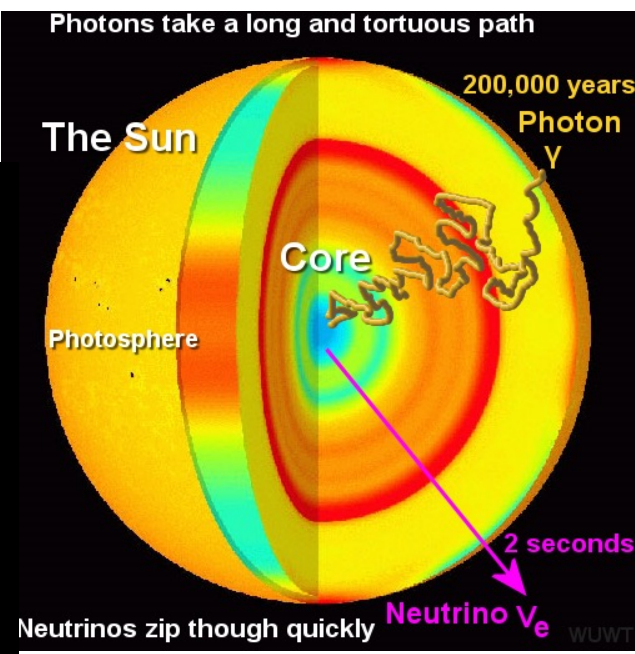
“...to see into the interior of a star and thus verify directly the hypothesis of nuclear energy generation..”

Bahcall and Davis, 1964

# Solar Neutrinos

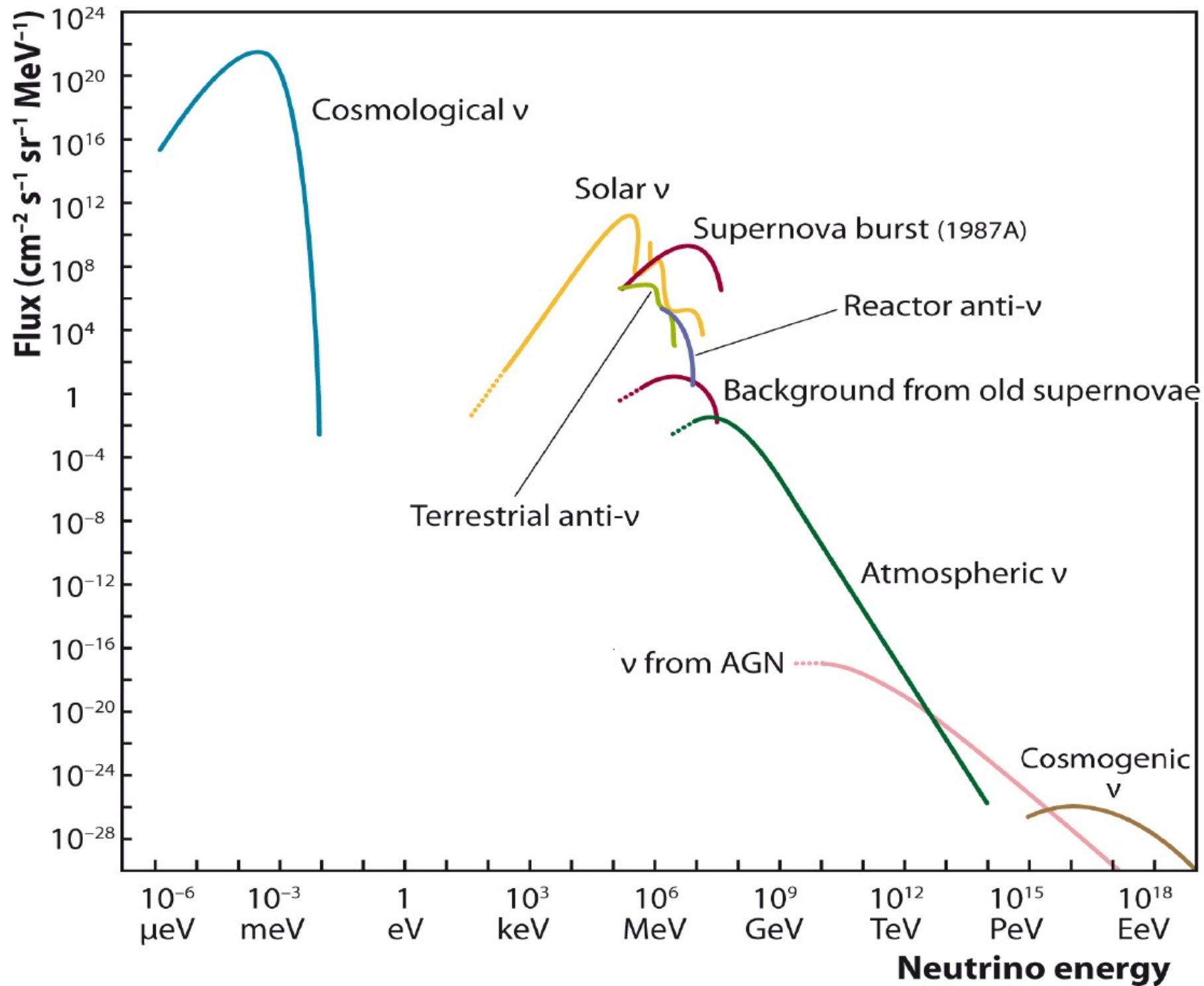


Copyright © 2010 Contemporary Physics Education Project



# Scientific reach of the underground laboratories

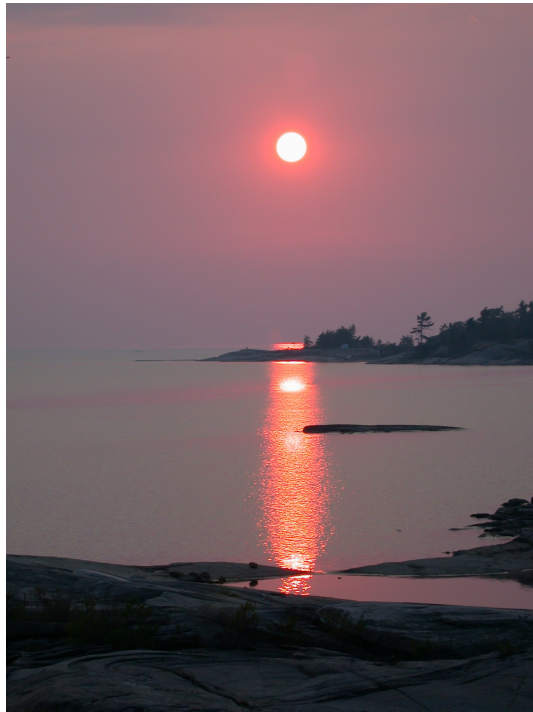
- Physics at the intensity frontier: Violation of fundamental symmetries (time reversal, lepton number, baryon number); neutrino properties ...
- Physics at the cosmic frontier: Particle and nuclear astrophysics; stellar and supernova neutrinos; nature of particle dark matter...
- Physics research carried at both frontiers is complementary to the research at the energy frontier.
- Low-background counting.
- Geophysics, dark life..



## Mixing matrix for three flavors

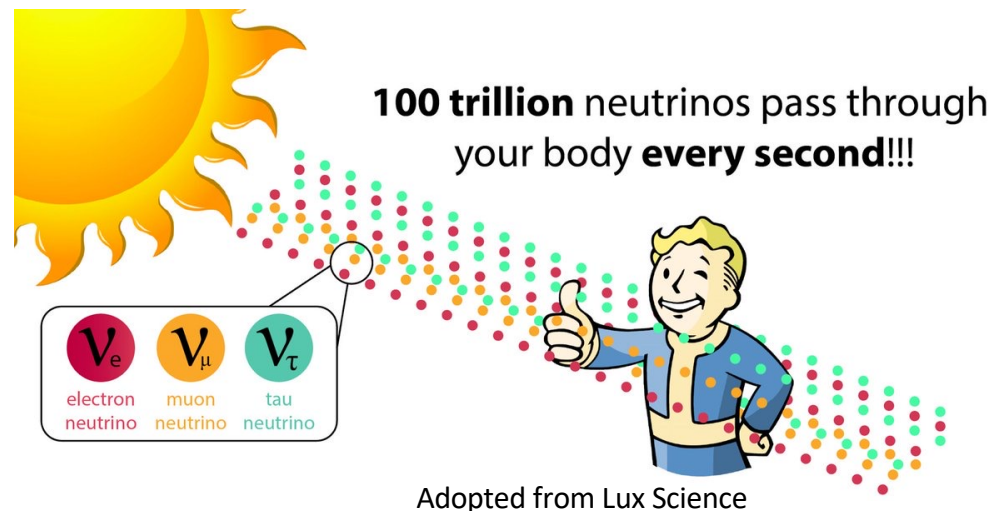
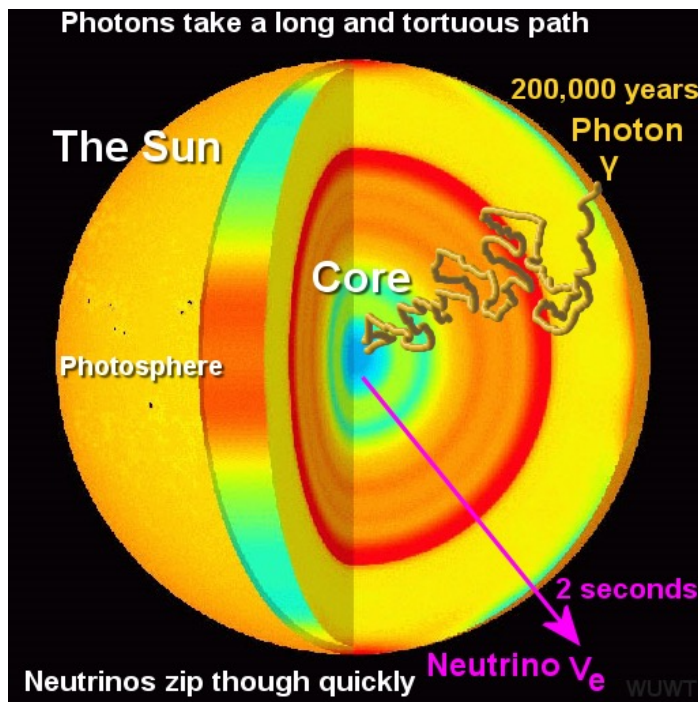
$$U = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{Atmospheric neutrinos}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\phi} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\phi} & 0 & c_{13} \end{pmatrix}}_{\text{Reactor neutrinos}} \\
 \times \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Solar neutrinos}} \underbrace{\begin{pmatrix} e^{i\alpha_1} & 0 & 0 \\ 0 & e^{i\alpha_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Majorana phases}}$$

$$c_{ij} = \cos \theta_{ij} \quad s_{ij} = \sin \theta_{ij}$$

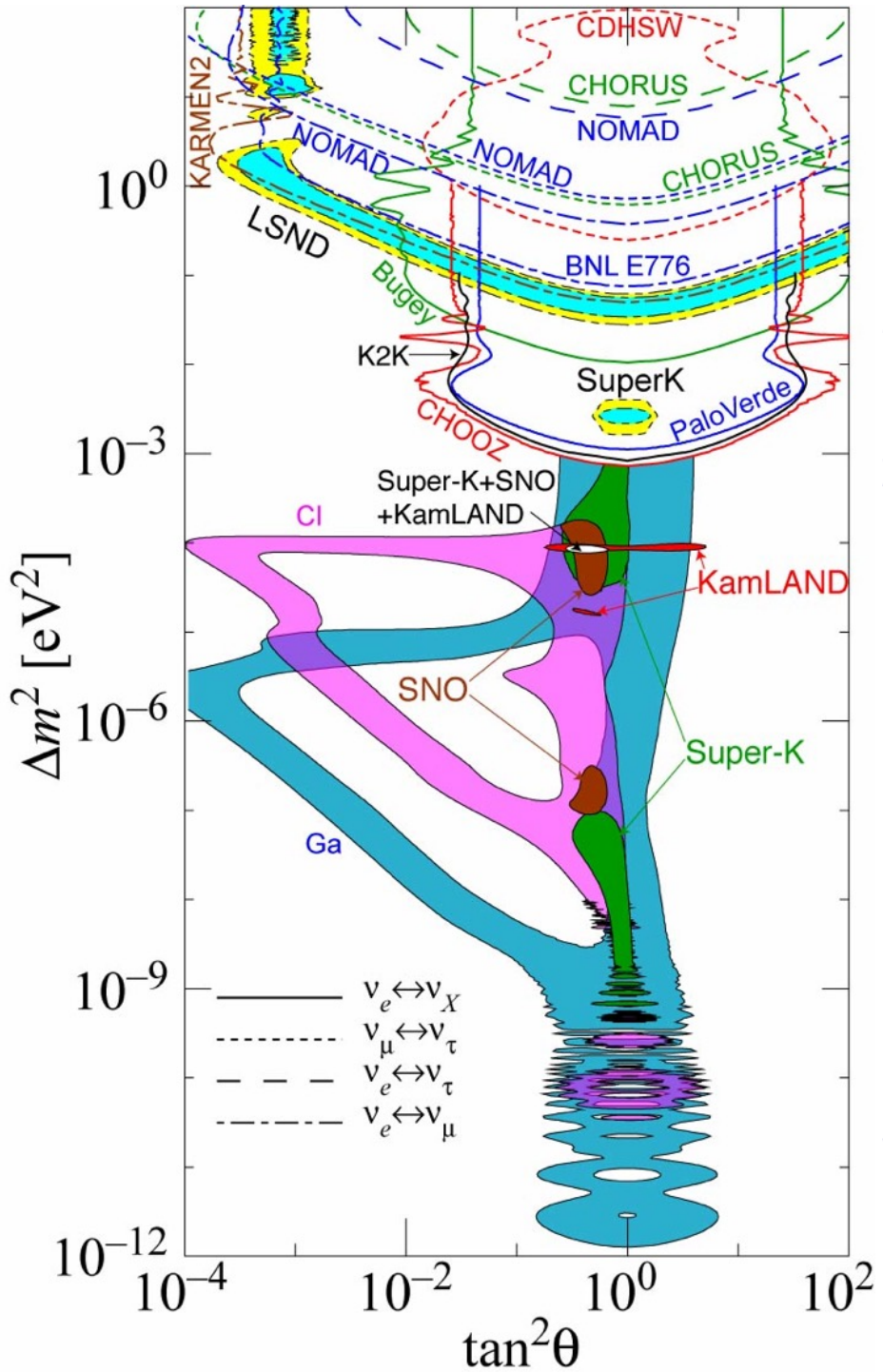


## Sources of neutrinos: Sun

A minor league star (such as our Sun) produces neutrinos mainly through the reaction



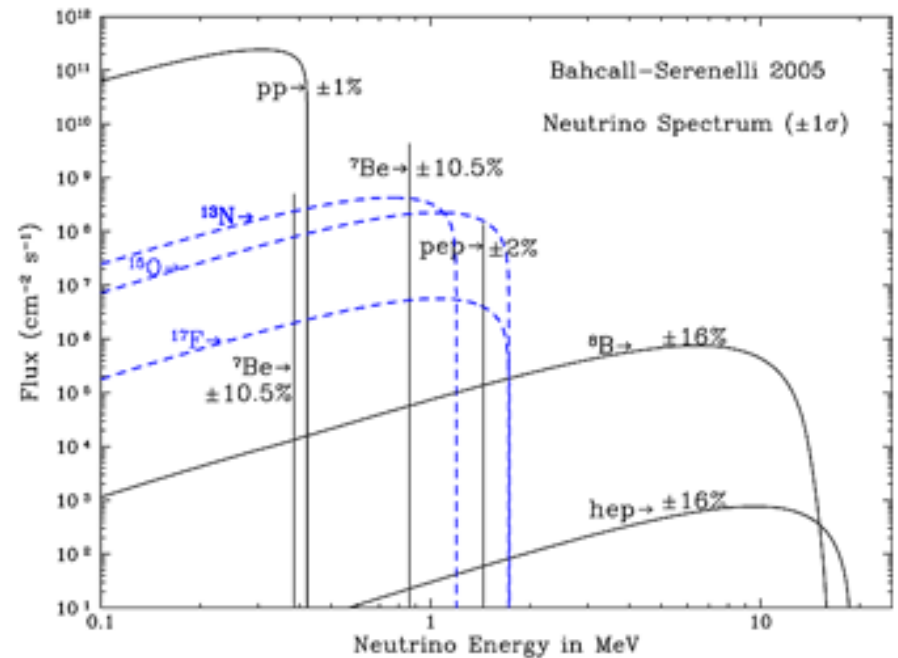
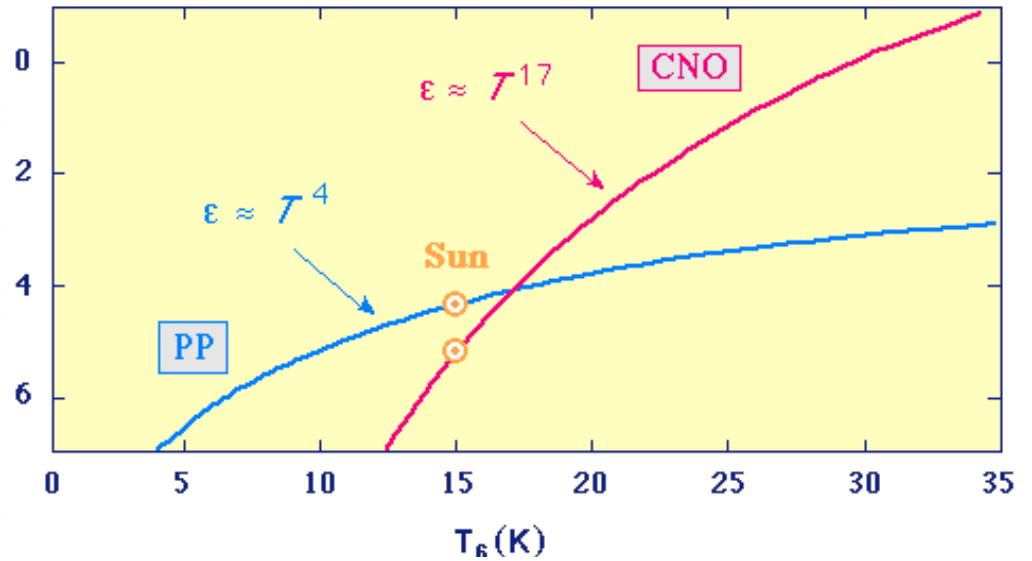
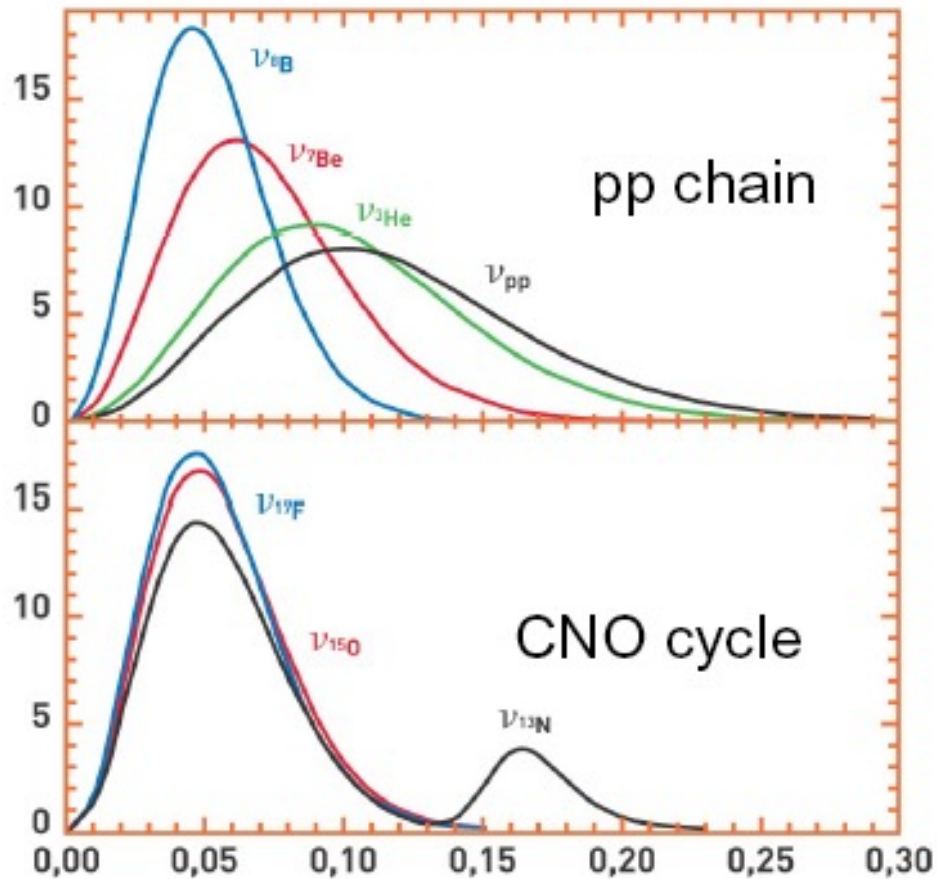




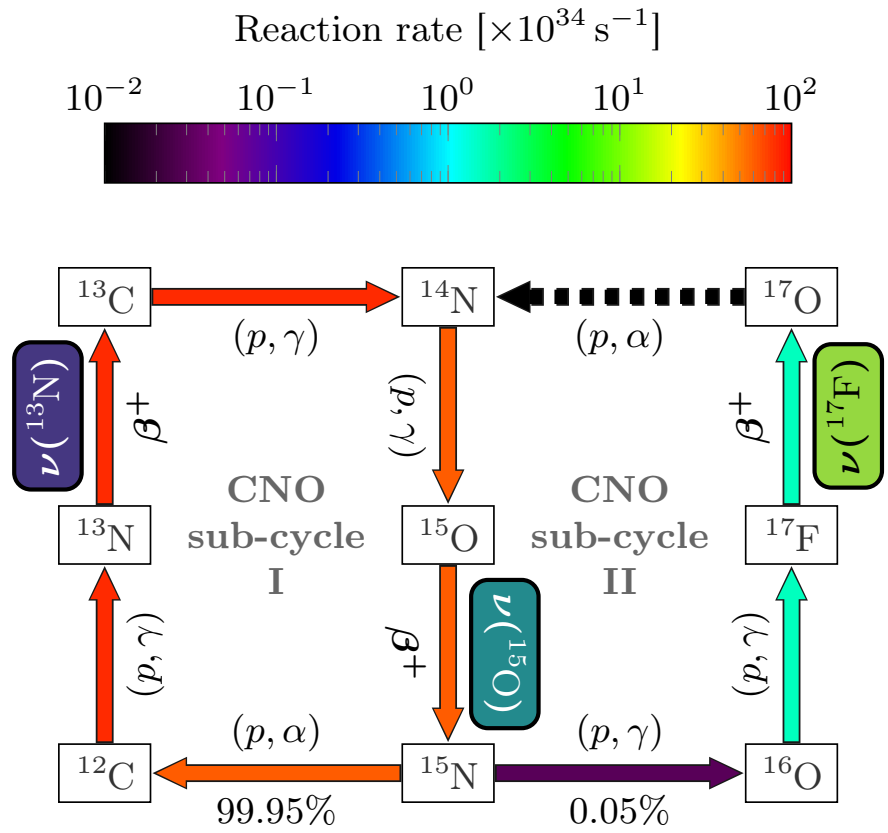
$\theta_{\text{atmospheric}}$  (primarily  $\theta_{23}$ )

$\theta_{\text{solar}}$  (primarily  $\theta_{12}$ )

# How much does the CNO cycle contribute in the Sun?

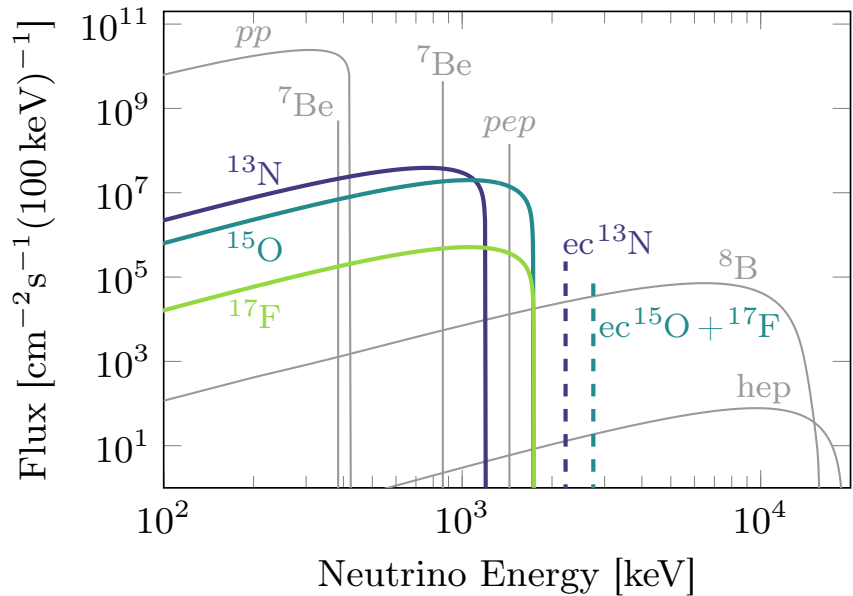


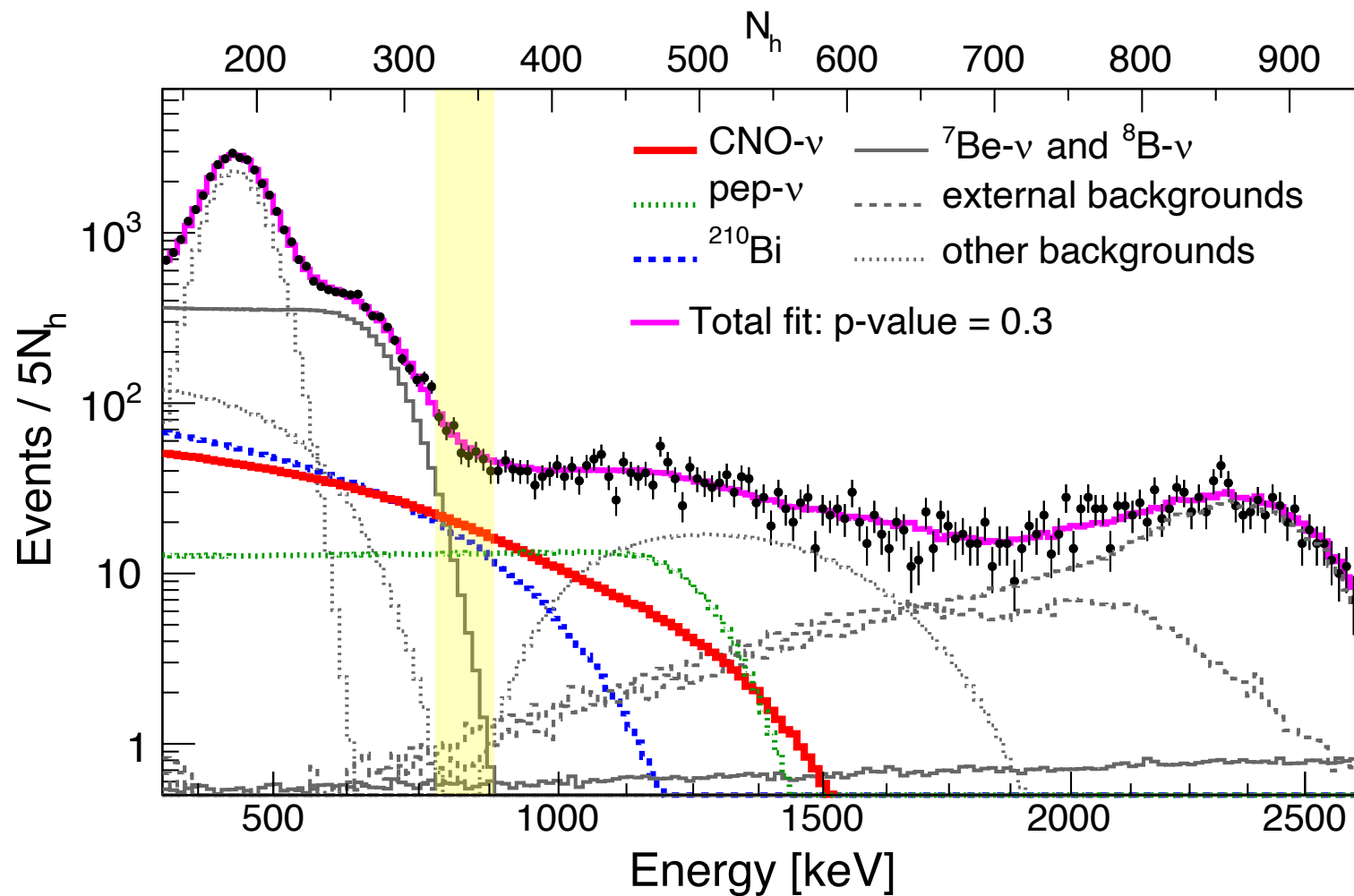
In SSM CNO cycle contribute about 0.8% of the neutrino flux. Data are consistent with this. A more precise measurement of the CNO contribution will provide a test of SSM.



CNO Neutrinos are now detected!

Borexino Collaboration  
2006.15115 [hep-ex]





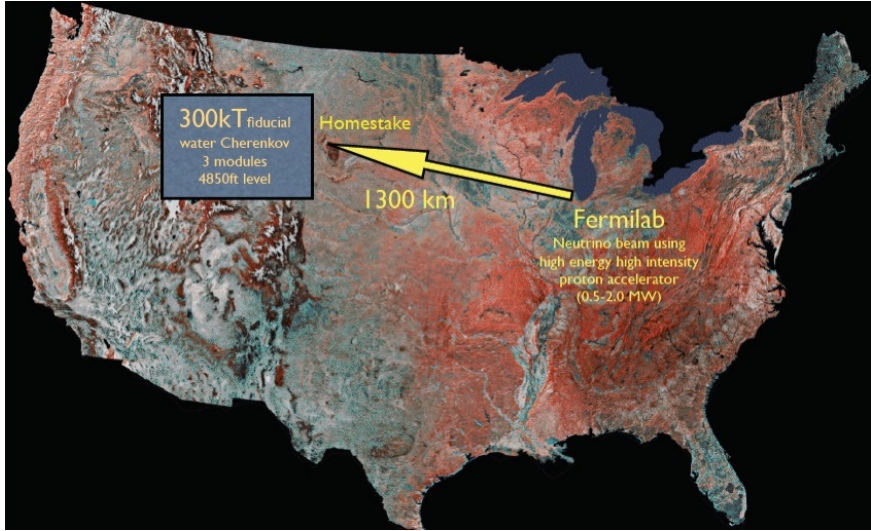
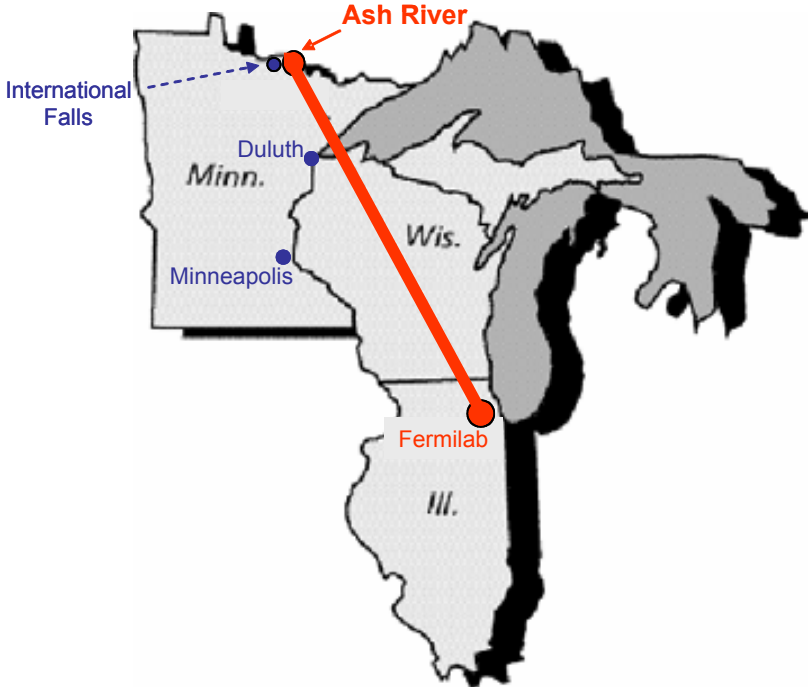
Borexino Collaboration  
 2006.15115 [hep-ex]

# Long-baseline oscillations at GeV energies

$$\text{Osc. max. } L \sim E \quad + \quad \begin{array}{l} \text{Flux at source } \sim E^2 \\ \text{Flux } (L) = \text{Flux } (L=0)/L^2 \end{array}$$

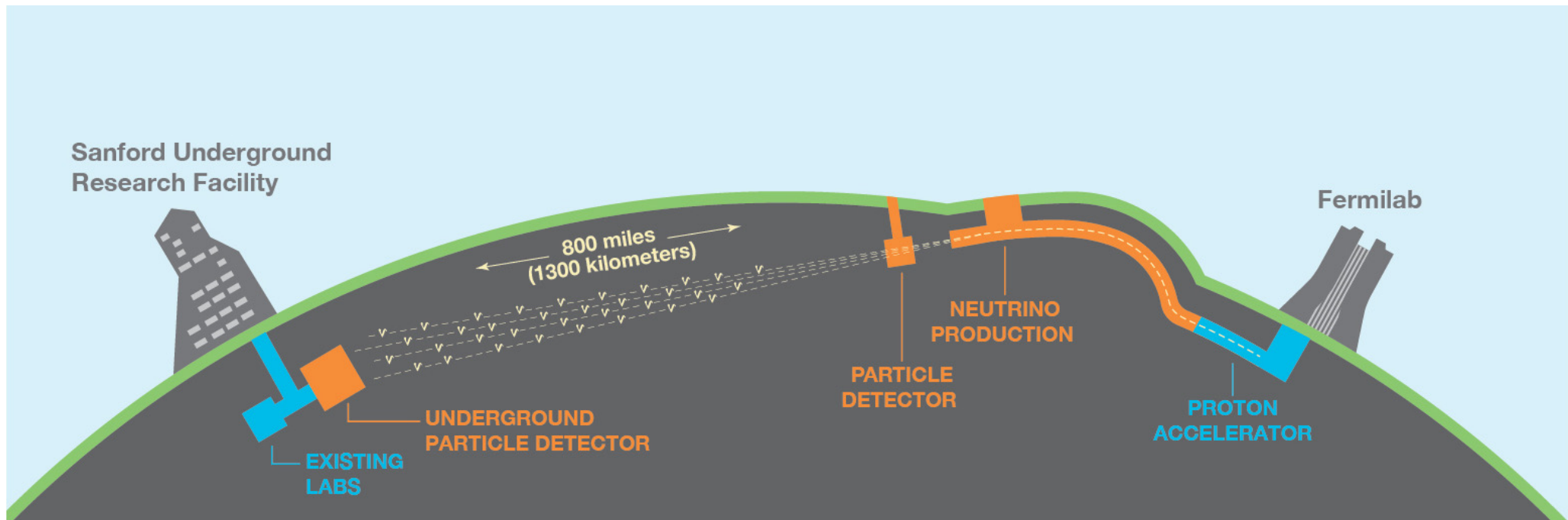
$$\text{Flux } (L) \sim 1 \quad + \quad \sigma \sim E \text{ (DIS)}$$

$$\text{Event rate } \sim E$$

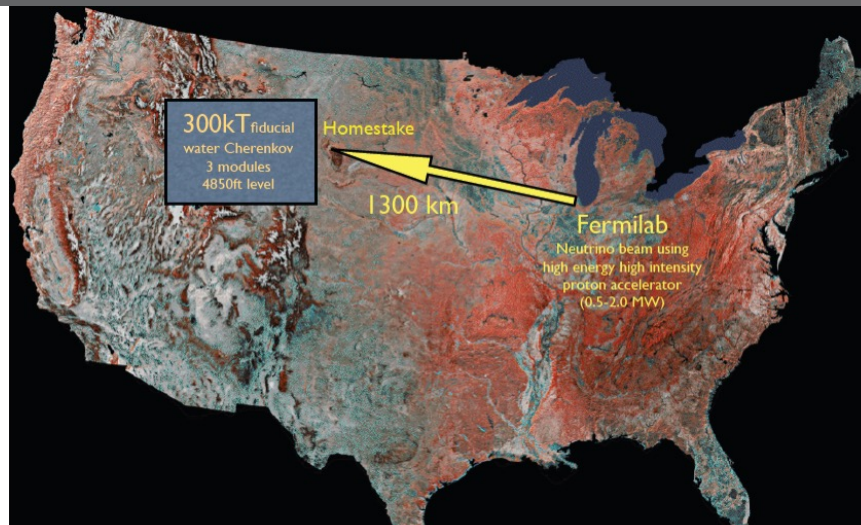




# DEEP UNDERGROUND NEUTRINO EXPERIMENT



The flagship  
experiment...



## Sources of neutrinos: Supernovae



A (hopefully distant enough) core-collapse supernova produces approximately  $10^{58}$  neutrinos in about twenty seconds primarily via

Gravitational binding energy  $\rightarrow \nu_x + \bar{\nu}_x$

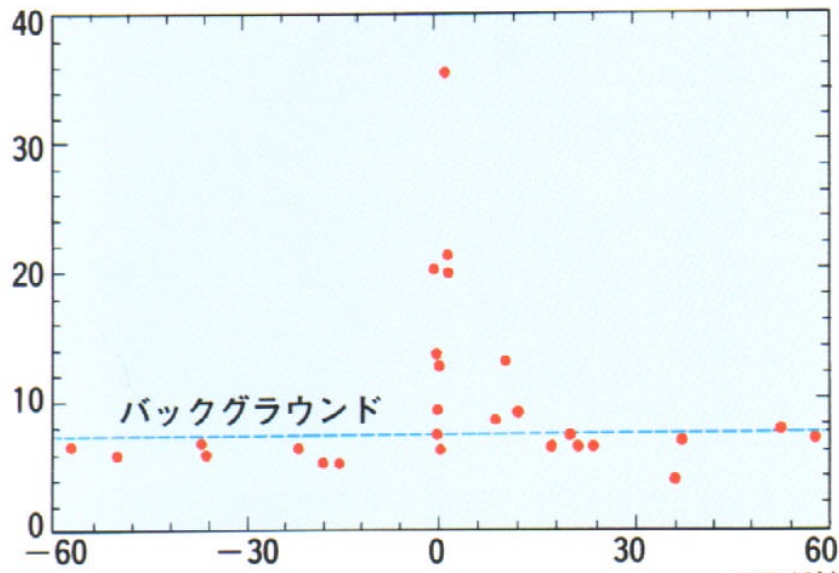
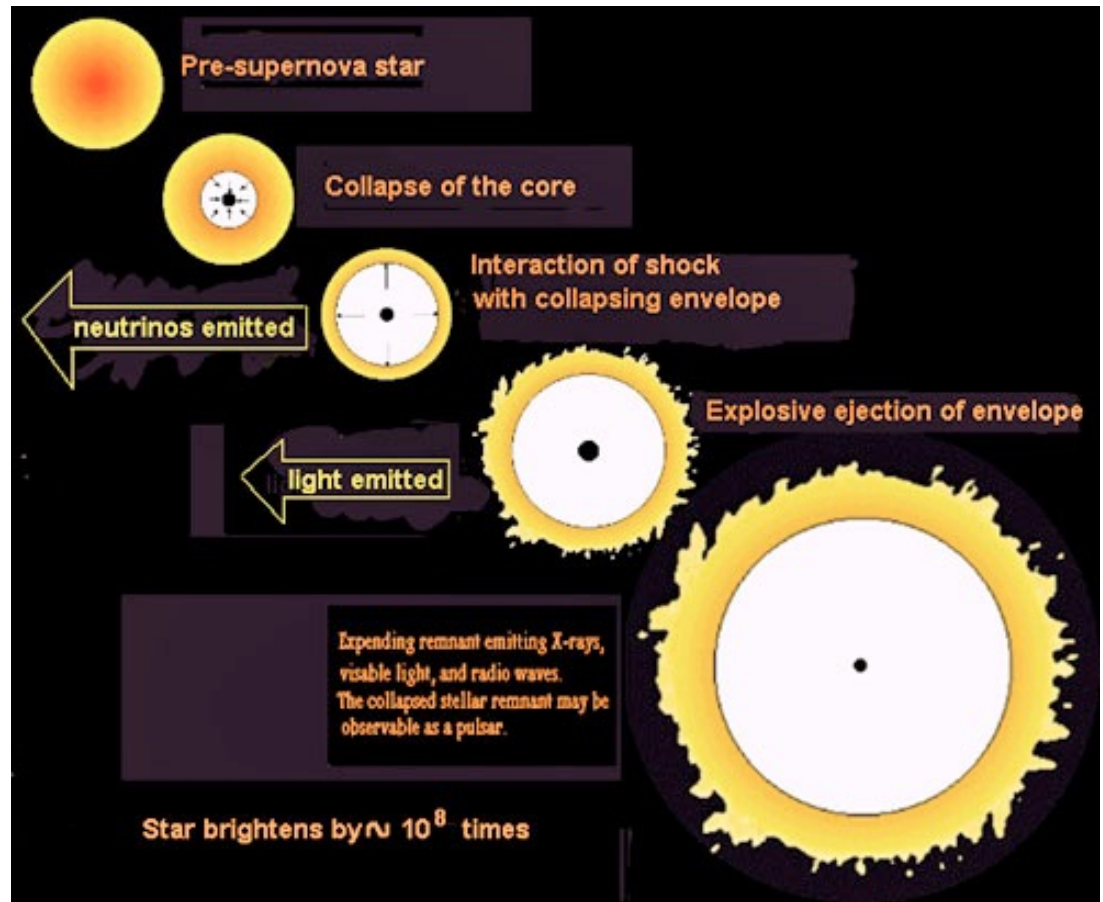
during the cooling and via



during the collapse.

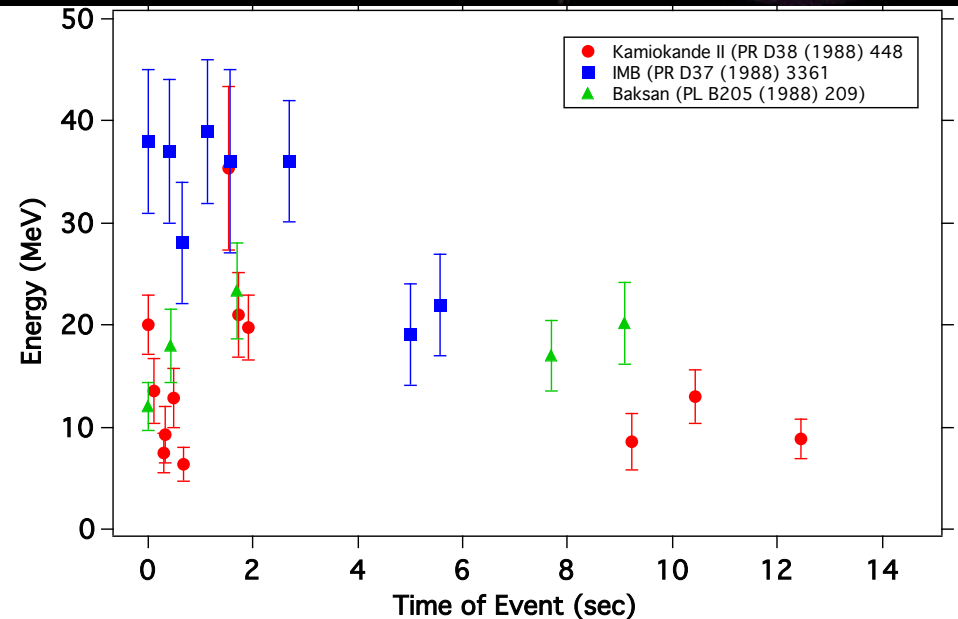
Those neutrinos produced in supernova explosions since the beginning of the Universe still stick around, forming the "Diffuse Supernova Background"

# Neutrinos from core-collapse supernovae



•  $M_{\text{prog}} \geq 8 M_{\text{sun}} \Rightarrow \Delta E \approx 10^{53} \text{ ergs} \approx 10^{59} \text{ MeV}$

• 99% of the energy is carried away by neutrinos and antineutrinos with  $10 \leq E_{\nu} \leq 30 \text{ MeV} \Rightarrow 10^{58}$  neutrinos

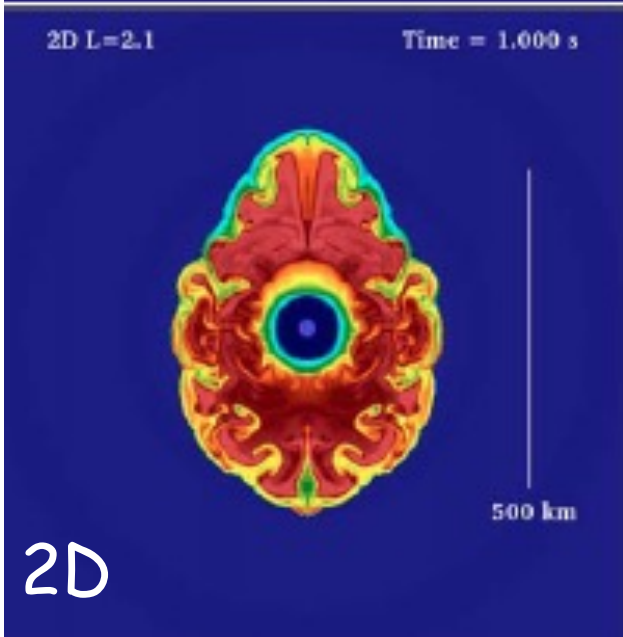
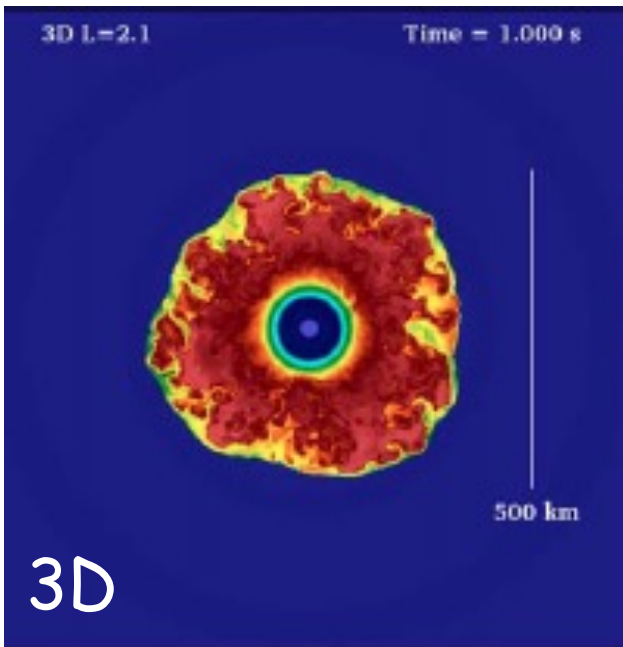




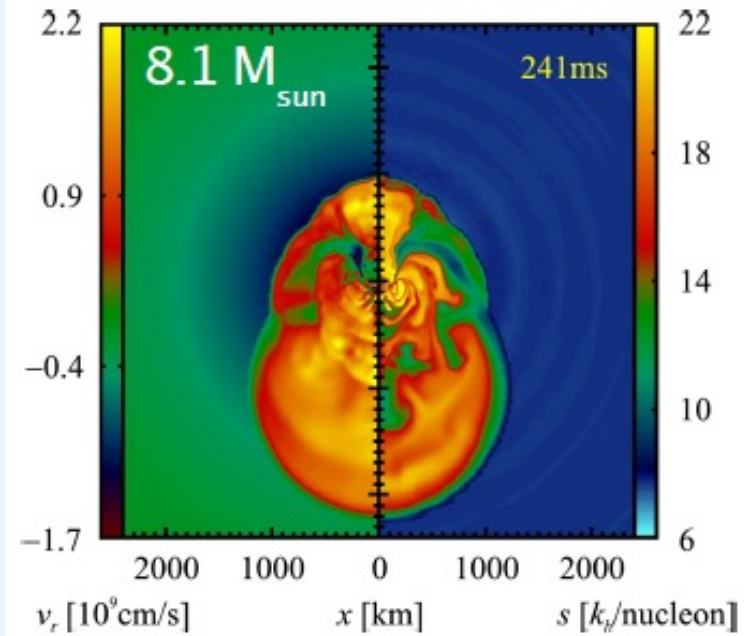
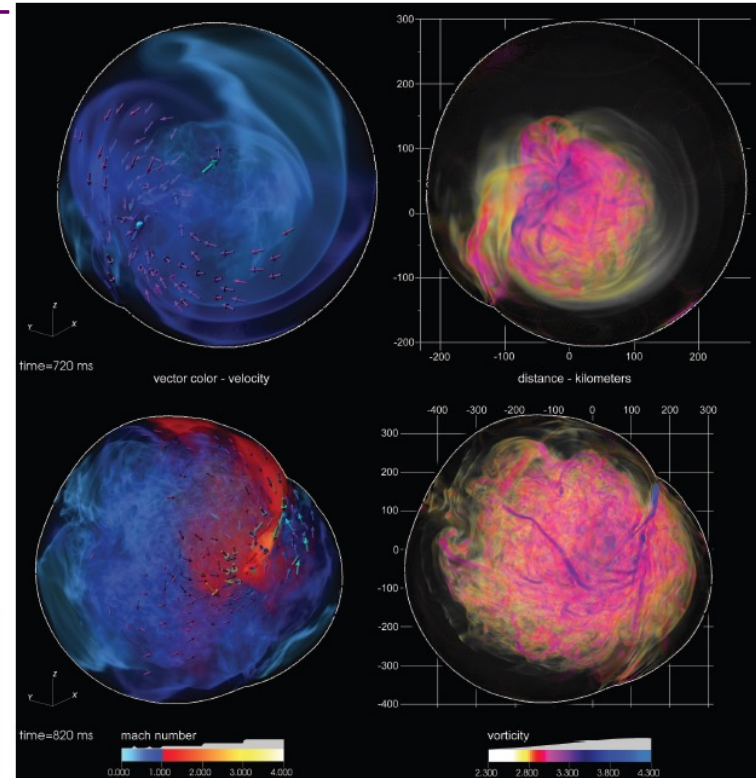
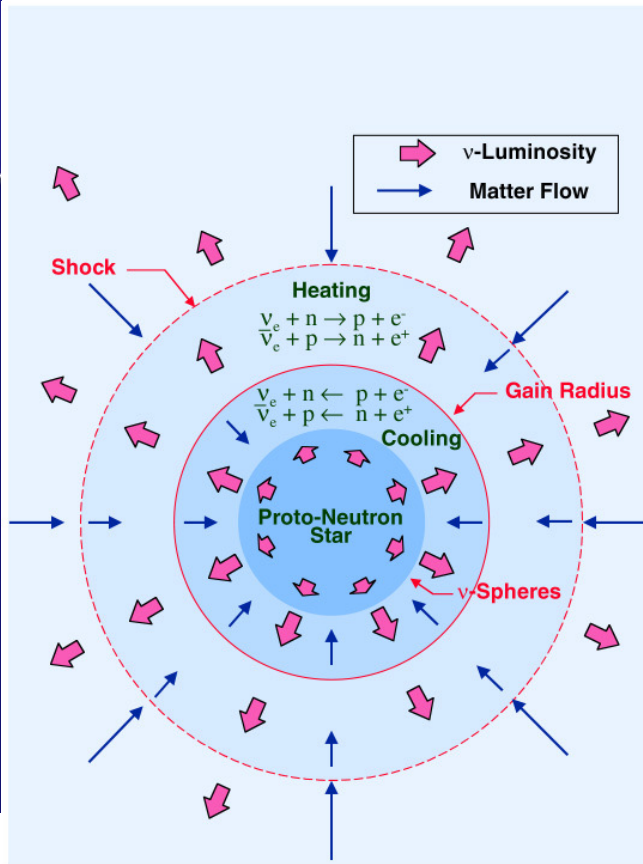


If we want to catch a supernova with neutrinos we'd better know what neutrinos do inside a supernova.

Development of 2D and 3D models for core-collapse supernovae: Complex interplay between turbulence, neutrino physics and thermonuclear reactions.

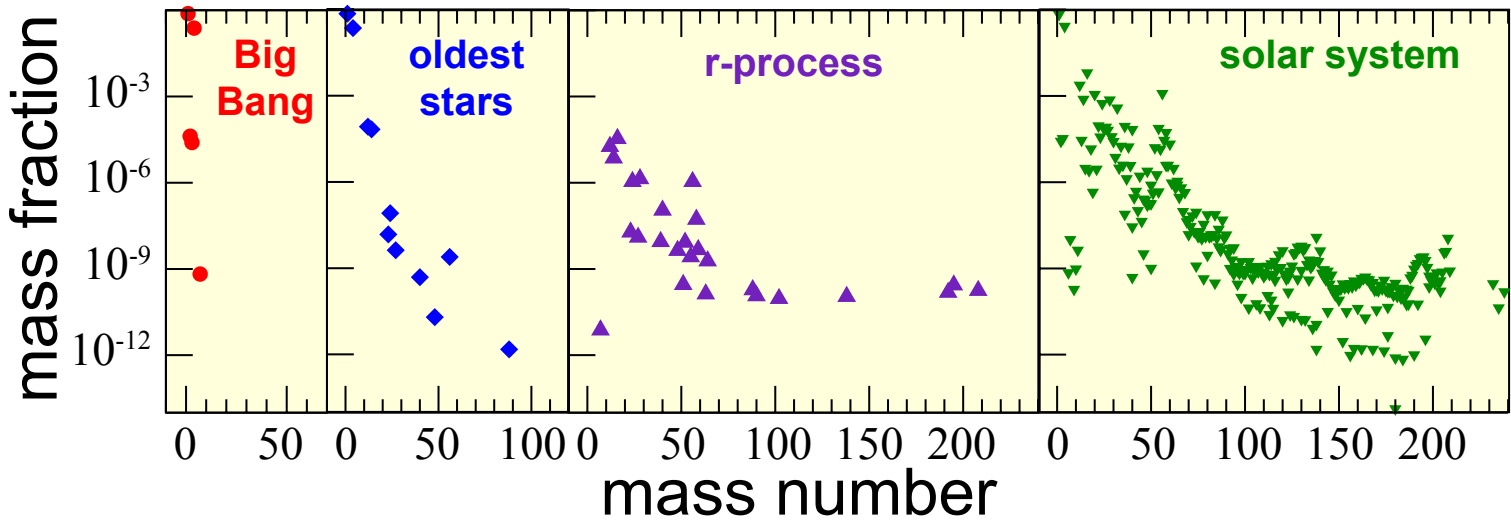


Princeton

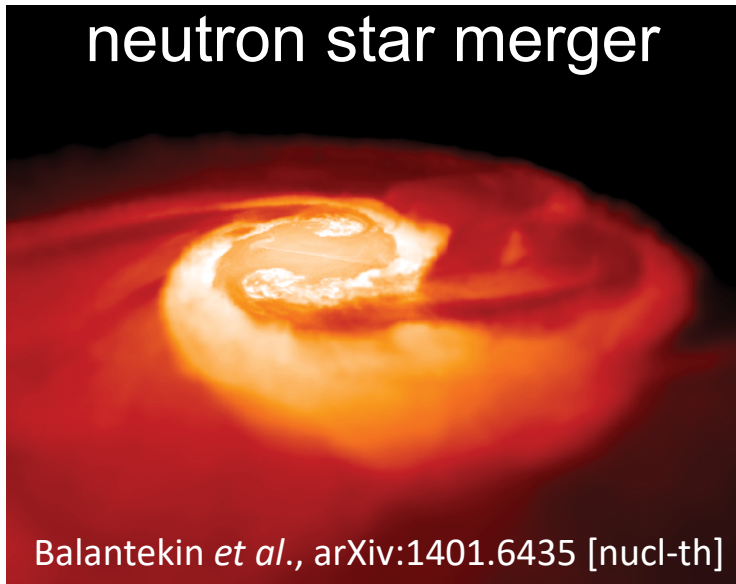
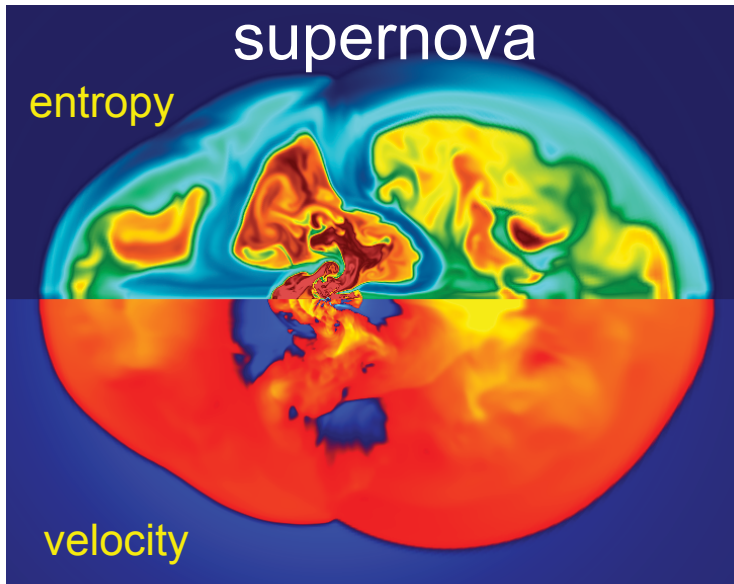


Munich

# The origin of elements

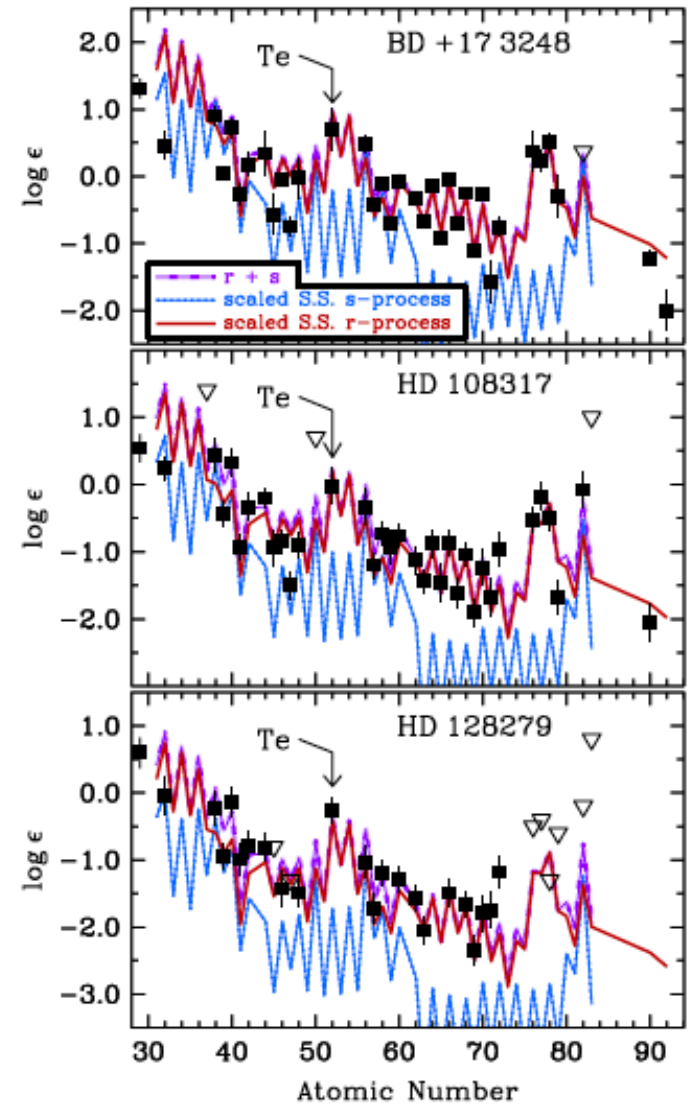
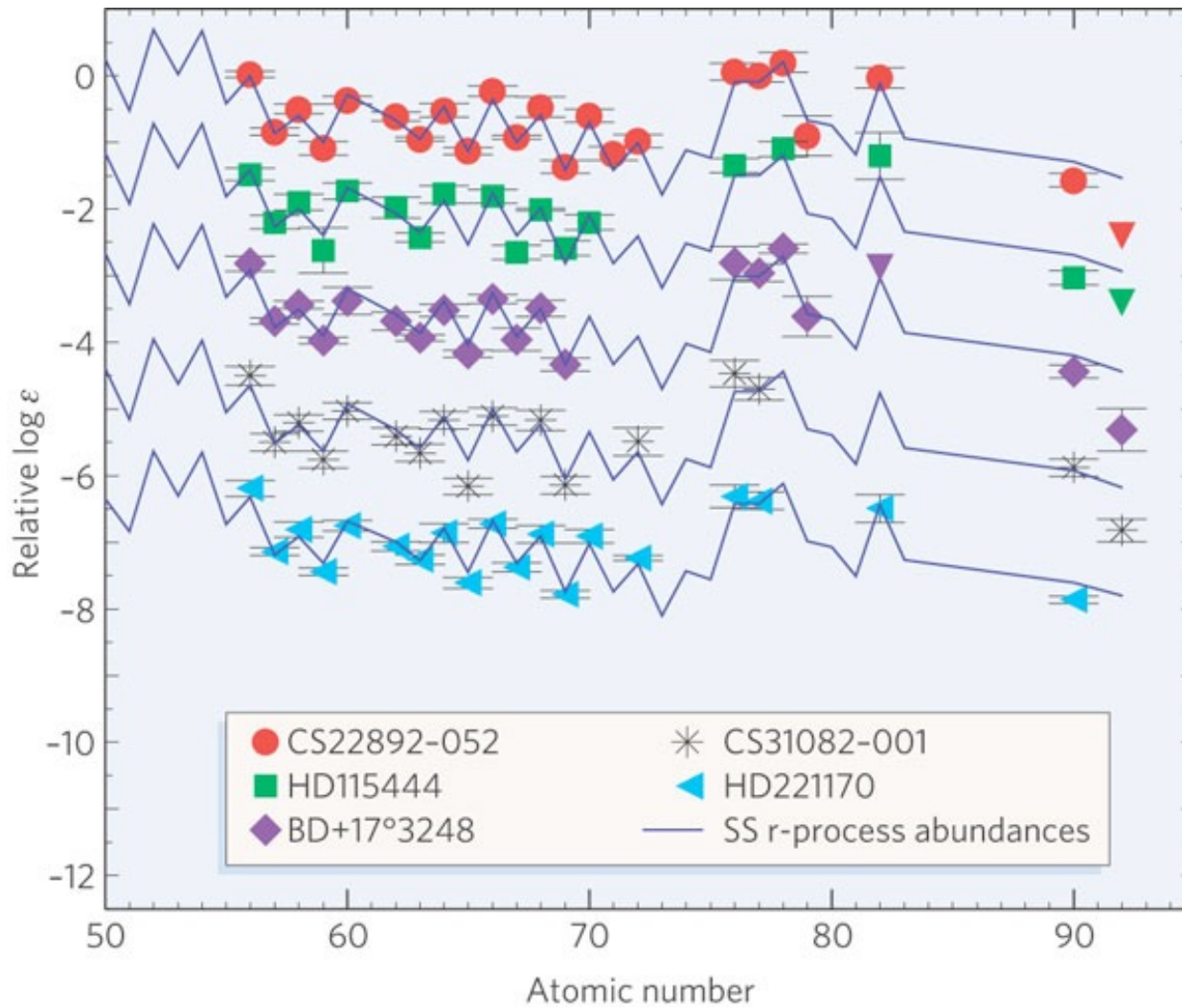


Neutrinos not only play a crucial role in the dynamics of these sites, but they also control the value of the electron fraction, the parameter determining the yields of the r-process.



Possible sites for the r-process

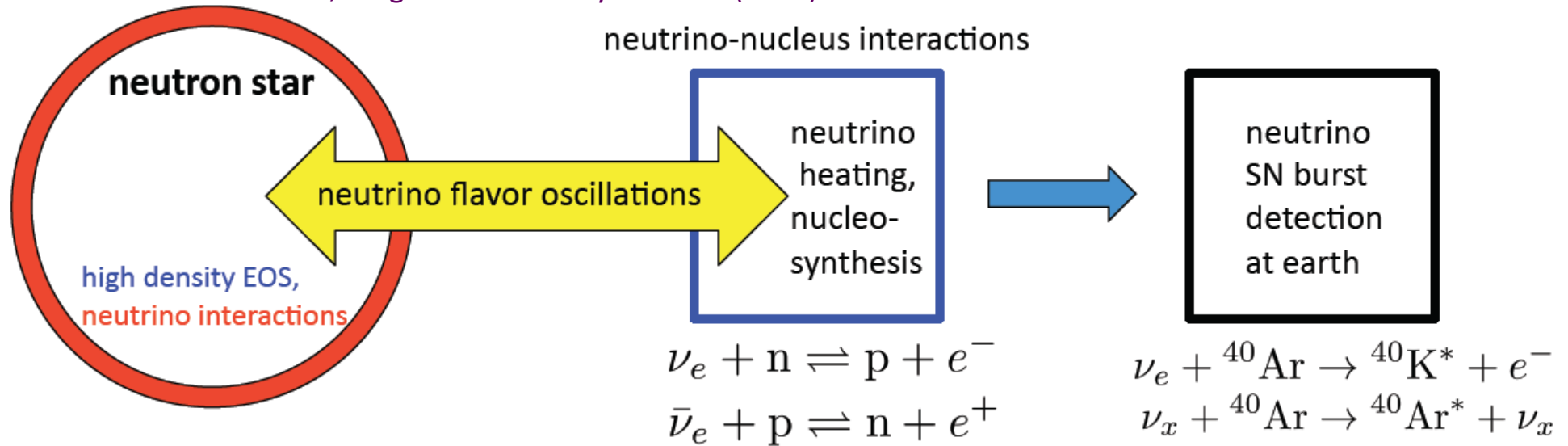
# r-process nucleosynthesis

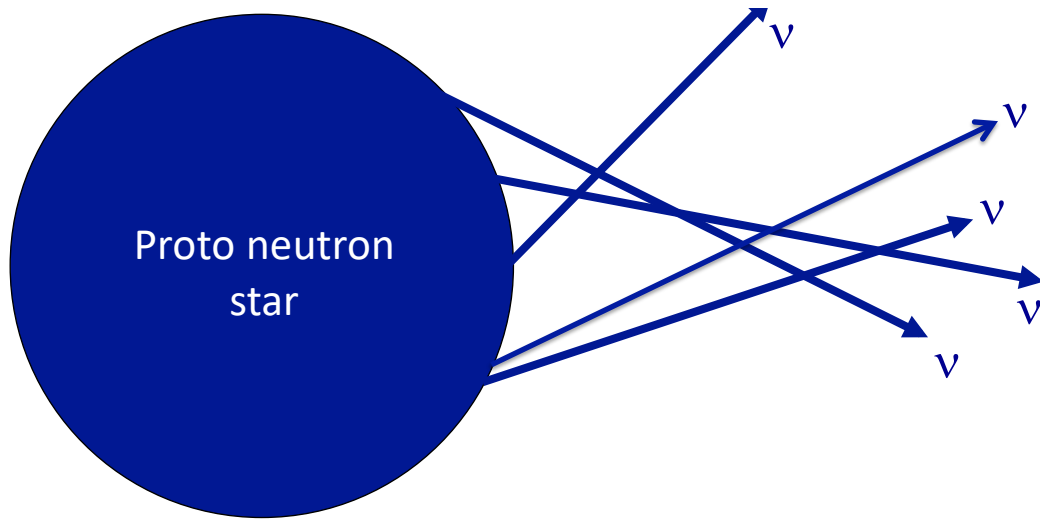


A > 100 abundance pattern fits the solar abundances well.

For example understanding a core-collapse supernova requires answers to a variety of questions some of which need to be answered, both theoretically and experimentally.

Balantekin and Fuller, Prog. Part. Nucl. Phys. **71** 162 (2013)





Energy released in a core-collapse  
SN:  $\Delta E \approx 10^{53}$  ergs  $\approx 10^{59}$  MeV  
99% of this energy is carried away  
by neutrinos and antineutrinos!  
 $\sim 10^{58}$  Neutrinos!  
This necessitates including the  
effects of  $\nu\nu$  interactions!

$$H = \underbrace{\sum a^\dagger a}_{\text{describes neutrino oscillations and interaction with matter (MSW effect)}} + \underbrace{\sum (1 - \cos\theta) a^\dagger a^\dagger a a}_{\text{describes neutrino-neutrino interactions}}$$

The second term makes the physics of a neutrino gas in a core-collapse supernova a very interesting many-body problem, driven by weak interactions.

Neutrino-neutrino interactions lead to novel collective and emergent effects, such as conserved quantities and interesting features in the neutrino energy spectra (spectral "swaps" or "splits").

## Many neutrino system

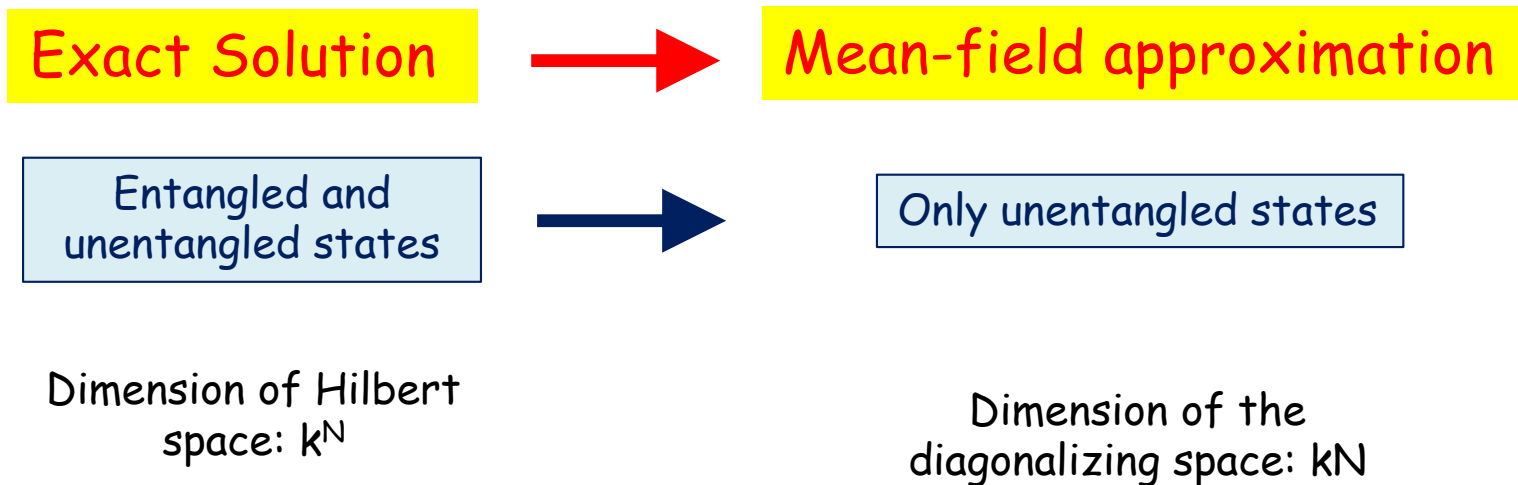
This is the only many-body system driven by the weak interactions:

Table: Many-body systems

<b>Nuclei</b>	Strong	at most $\sim 250$ particles
<b>Condensed matter</b>	E&M	at most $N_A$ particles
<b><math>\nu</math>'s in SN</b>	Weak	$\sim 10^{58}$ particles

Astrophysical extremes allow us to test physics that cannot be tested elsewhere!

A system of  $N$  particles each of which can occupy  $k$  states ( $k$  = number of flavors)



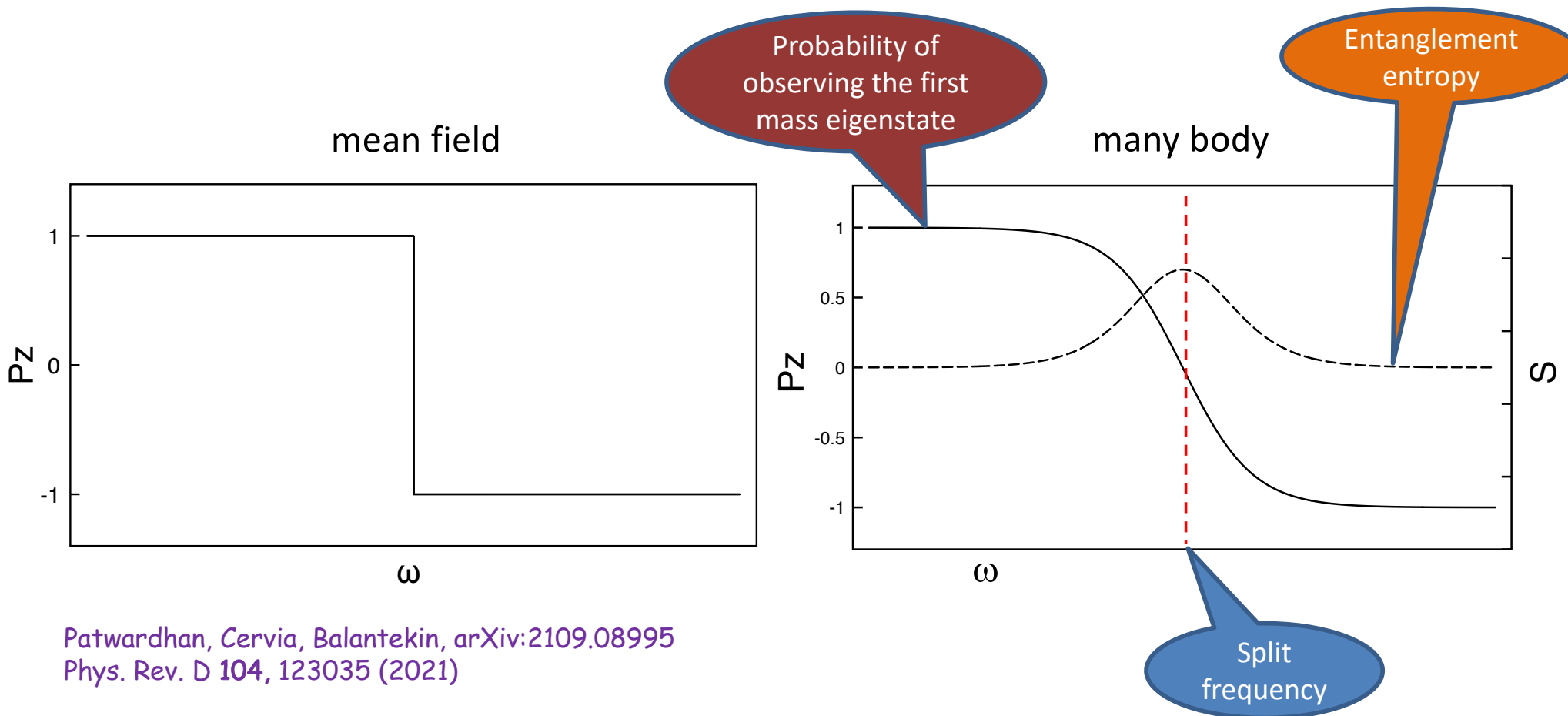
von Neumann entropy

$$S = - \text{Tr} (\rho \log \rho)$$

	Pure State	Mixed State
Density matrix	$\rho^2 = \rho$	$\rho^2 \neq \rho$
Entropy	$S = 0$	$S \neq 0$

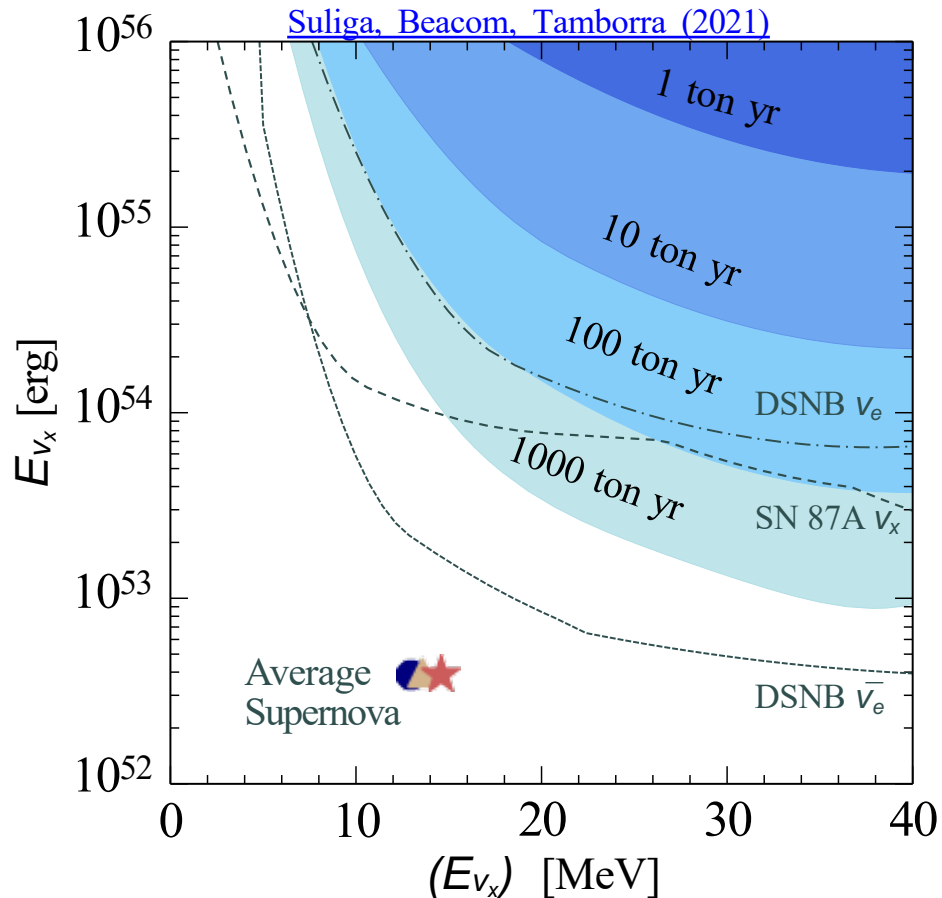


We find that the presence of **spectral splits** is a good **proxy** for deviations from the mean-field results



Patwardhan, Cervia, Balantekin, arXiv:2109.08995  
Phys. Rev. D **104**, 123035 (2021)

# Diffuse supernova neutrino background (DSNB)



- $\bar{\nu}_e$ : soon to be detected by SK + Gd, JUNO
- $\nu_e$ : possibly detectable by DUNE
- $\nu_x$ : CEvNS detectors can improve the existing limits to almost  $\bar{\nu}_e$  level

## Detection of all flavors required to

- rule out potential non-standard scenarios
- bring us closer to understanding the supernova physics

[Guseinov \(1967\)](#), [Totani et al. \(2009\)](#), [Ando, Sato \(2004\)](#), [Lunardini \(2009\)](#), [Beacom \(2010\)](#), [Horiuchi et al. \(2011\)](#), [Lunardini, Tamborra \(2012\)](#), [Møller, Suliga, Tamborra, Denton \(2018\)](#), [Nakazato et al. \(2018\)](#), [Kresse et al. \(2020\)](#) ...



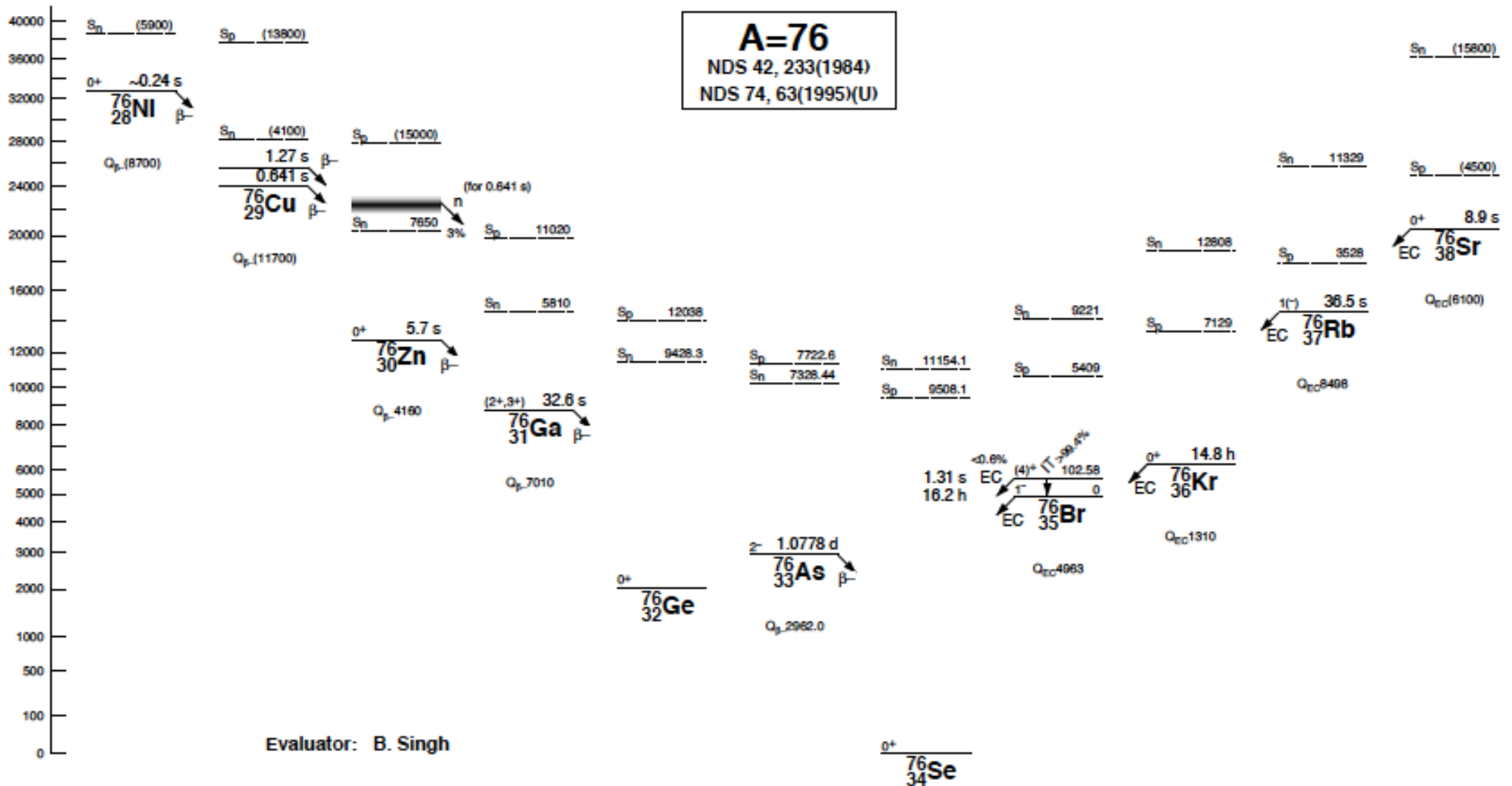
## Double Beta Decay

The second order process, where two neutrinos are emitted, is also possible.

Maria Mayer, 1935

*Maria Goeppert Mayer was awarded the 1963 Nobel for the nuclear shell model, the San Diego Union headline sadly read "San Diego Housewife Wins Nobel Prize".*

# Pairing gives rise to double beta decay:



## Some measurements of $2\beta\beta$ decay

Nucleus	Q-value (MeV)	$T_{1/2}$ (years)
$^{48}\text{Ca}$	4.276	$(3.9\pm 0.7\pm 0.6) \times 10^{19}$
$^{76}\text{Ge}$	2.039	$(1.7\pm 0.2) \times 10^{21}$
$^{82}\text{Se}$	2.992	$(9.6\pm 0.3\pm 1.) \times 10^{19}$
$^{100}\text{Mo}$	3.034	$(7.11\pm 0.02\pm 0.54) \times 10^{18}$
$^{116}\text{Cd}$	2.804	$(2.8\pm 0.1\pm 0.3) \times 10^{19}$
$^{128}\text{Te}$	0.876	$(2.0\pm 0.1) \times 10^{24}$
$^{130}\text{Te}$	2.529	$(7.6\pm 1.5\pm 0.8) \times 10^{20}$
$^{136}\text{Xe}$	2.467	$(1.1) \times 10^{25}$
$^{150}\text{Nd}$	3.368	$(9.2\pm 0.25\pm 0.73) \times 10^{21}$

# Suggestion of neutrinoless double beta decay

Nuovo Cimento, **14**, pp 322-328 (1937)



## SULLA SIMMETRIA TRA PARTICELLE E ANTIPARTICELLE

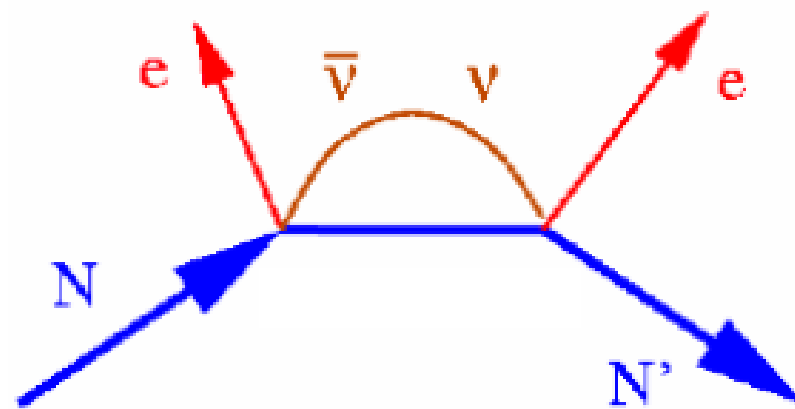
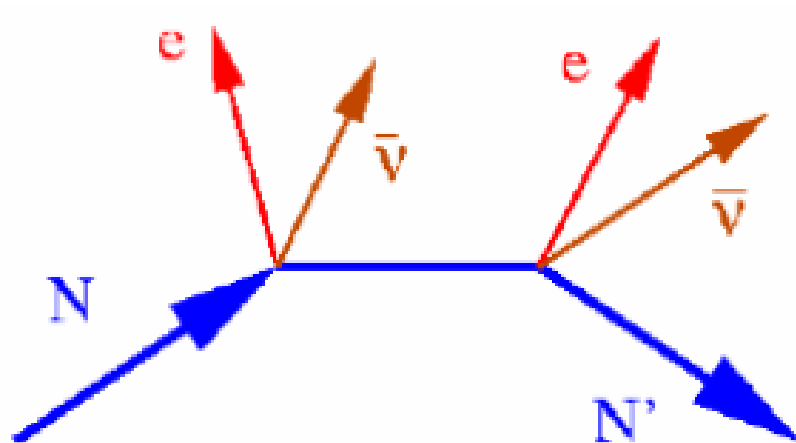
Nota di GIULIO RACAH

*Sunto.* - Si mostra che la simmetria tra particelle e antiparticelle porta alcune modificazioni formali nella teoria di FERMI sulla radioattività  $\beta$ , e che l'identità fisica tra neutrini ed antineutrini porta direttamente alla teoria di E. MAJORANA.

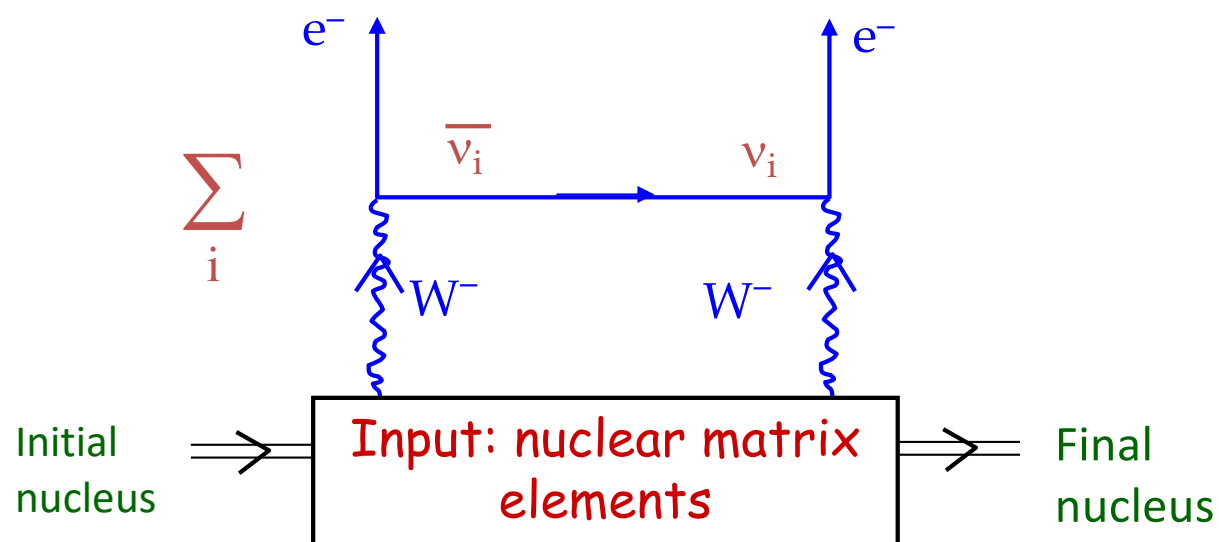
**Summary** - This article shows that the symmetry between particles and antiparticles leads some formal amendments in the theory of Fermi  $\beta$  radioactivity, and that the physical identity between neutrinos and antineutrinos leads directly to the theory of E. Majorana.

$2\nu\beta\beta$

$0\nu\beta\beta$



Majorana nature of the neutrinos permit  
neutrinoless double beta decay:



For Majorana neutrino exchange the leptonic part of the amplitude is given as a sum over mass eigenstates:

$$\mathcal{L}_{\mu\nu} = \sum_i [\bar{e}(x)] \gamma_\mu (1 - \gamma_5) U_{ei} \nu_i(x) \bar{\nu}_i^c U_{ei} \gamma_\nu (1 + \gamma_5) e^c(y)$$

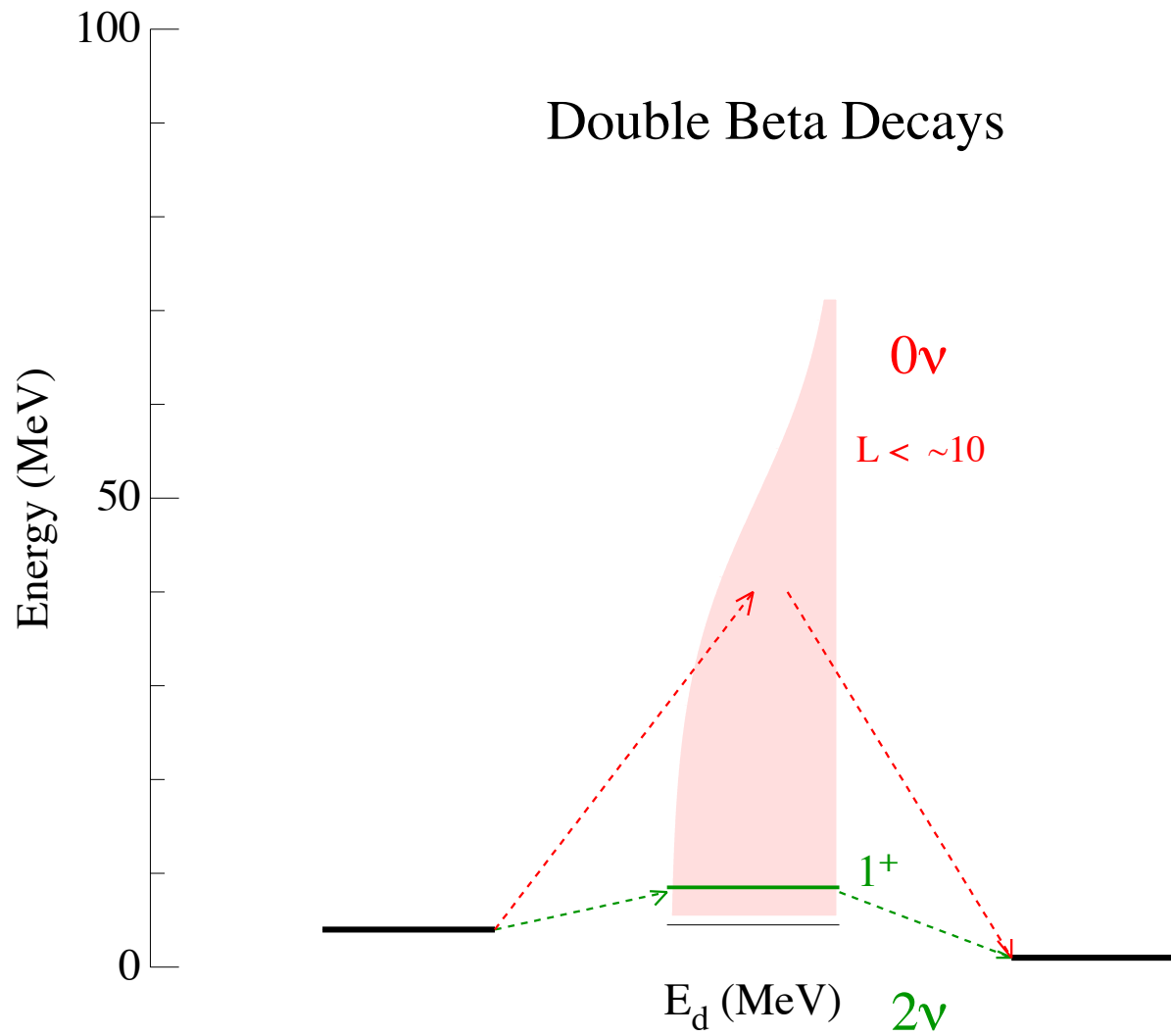
Contracting the two neutrino fields gives  $\frac{q_\mu \gamma^\mu - m_i}{q^2 - m_i^2}$

$q_\mu \gamma^\mu$  term does not contribute to the traces making leptonic tensor proportional to the quantity

$$m_{\beta\beta} = \sum_{i=1}^3 U_{ei}^2 m_i$$



# Double Beta Decays

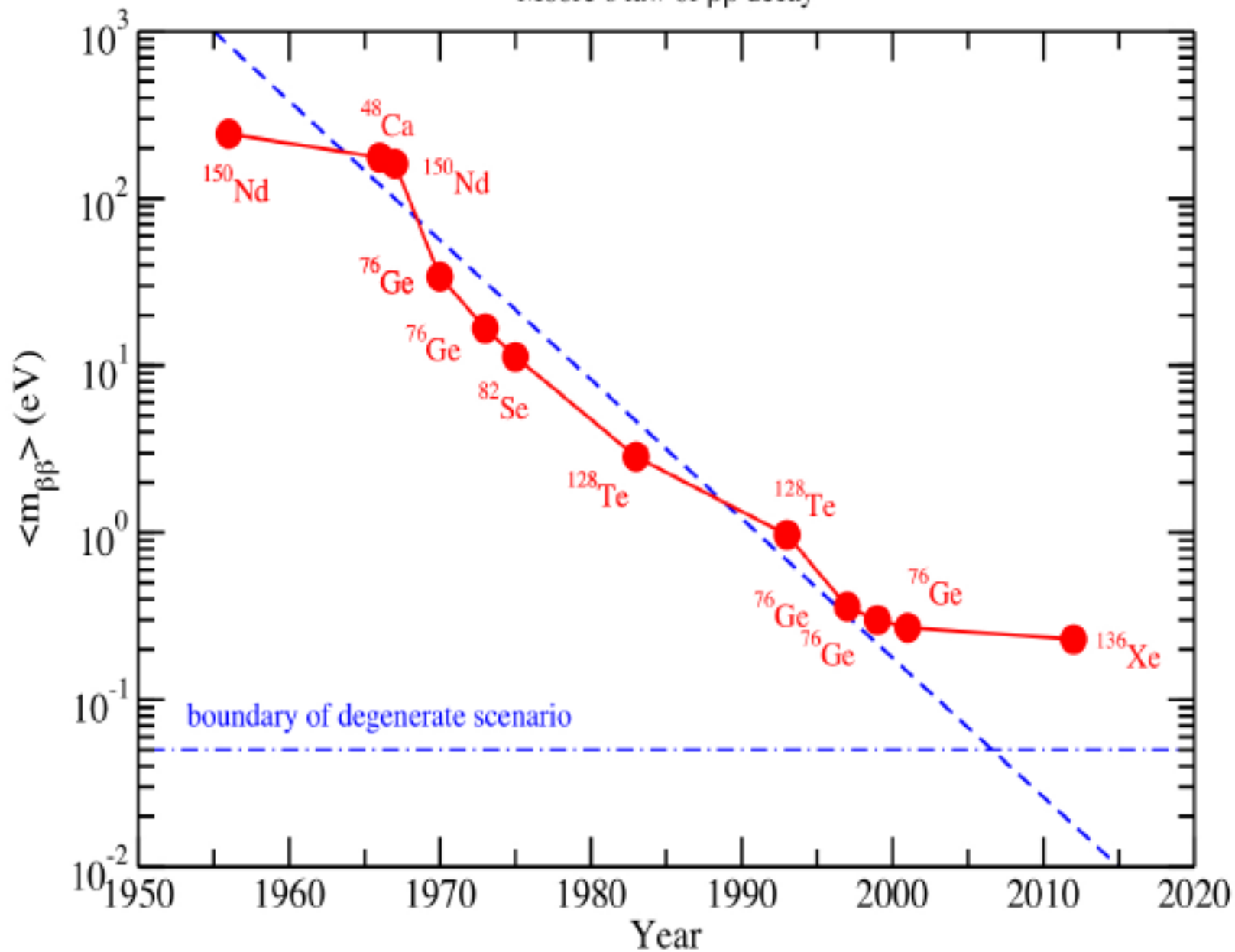


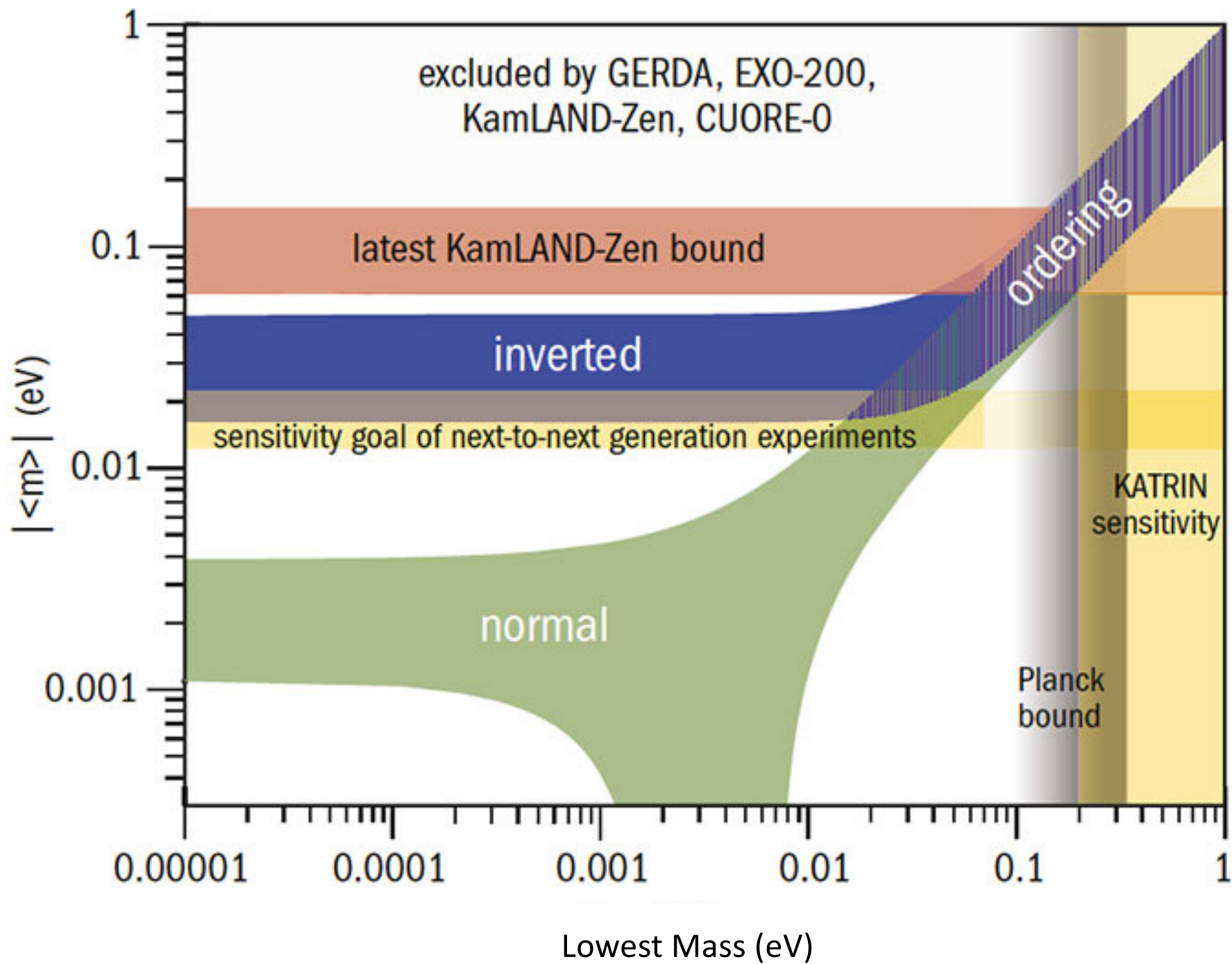
$$\frac{1}{T_{1/2}^{2\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

$$\frac{1}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \left| \sum_i U_{ei}^2 m_i \right|^2$$

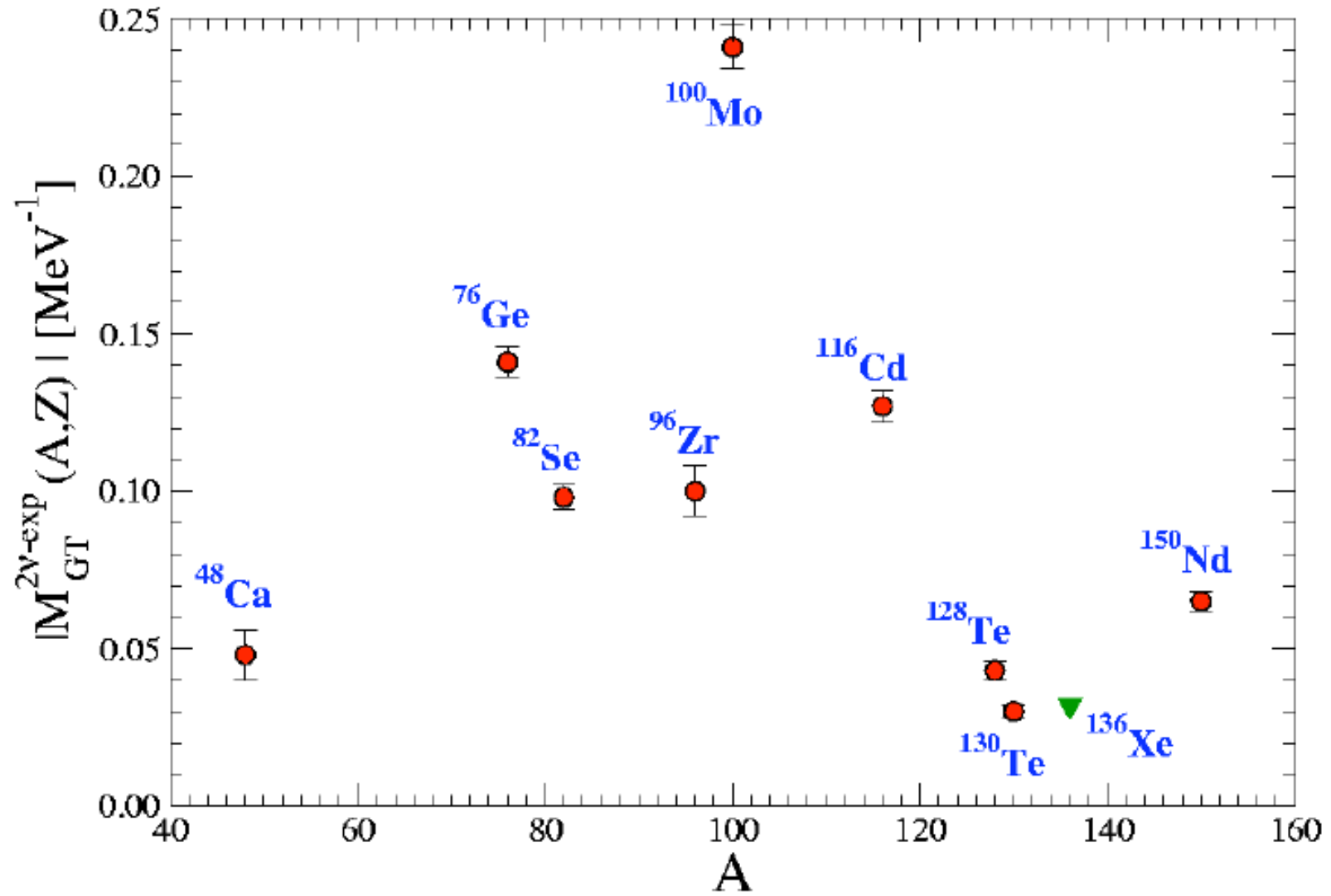
# History of the $0\nu\beta\beta$ decay

Moore's law of  $\beta\beta$  decay





For  $2\nu\beta\beta$  there is a strong shell-model dependence of the matrix elements



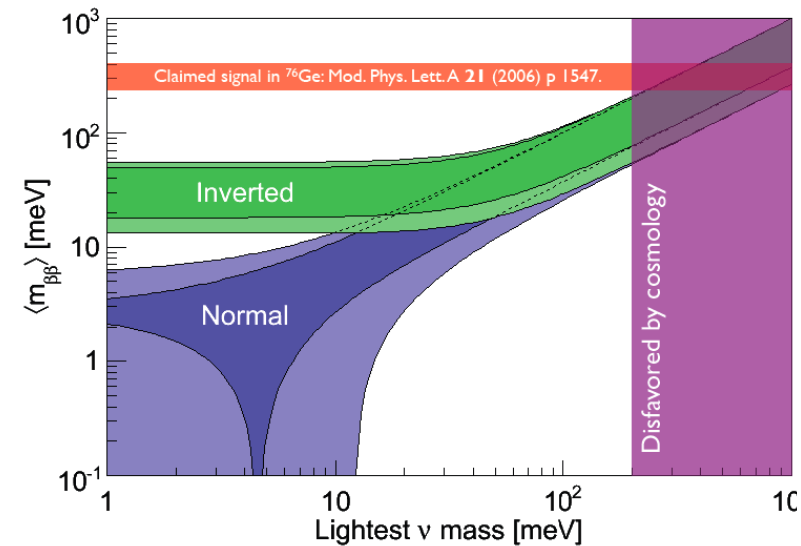
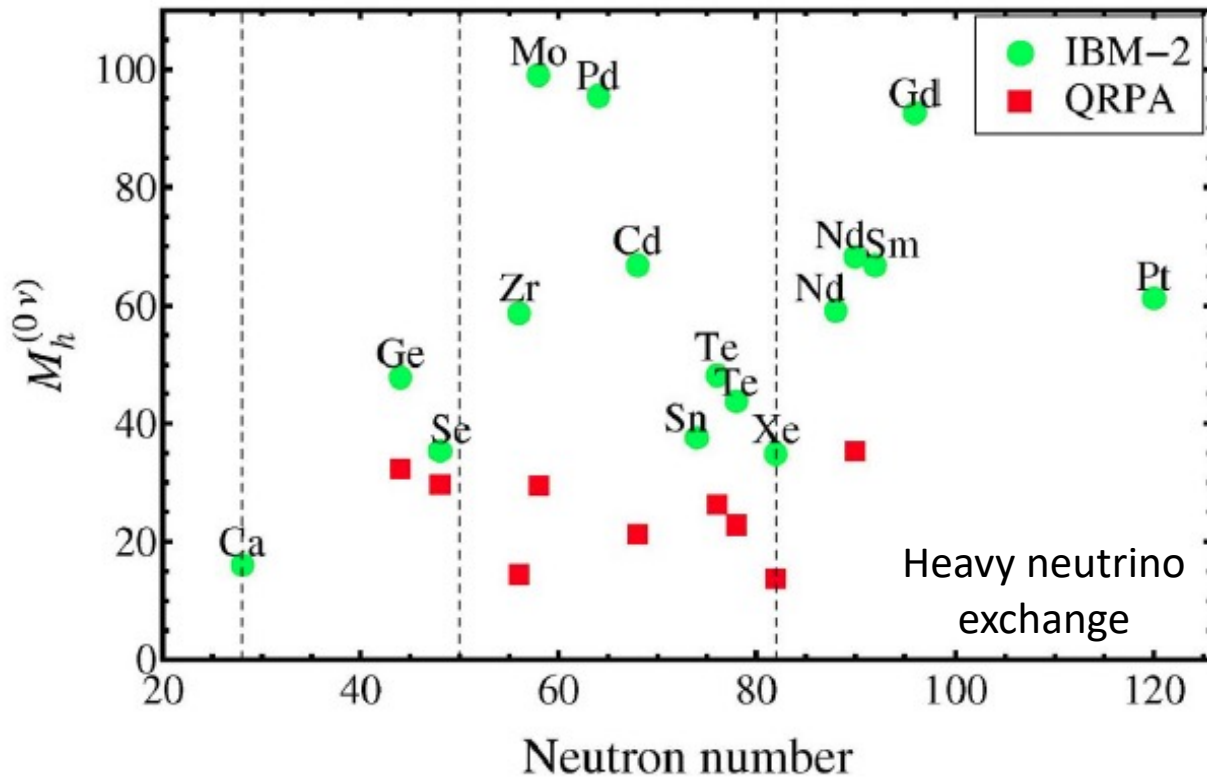
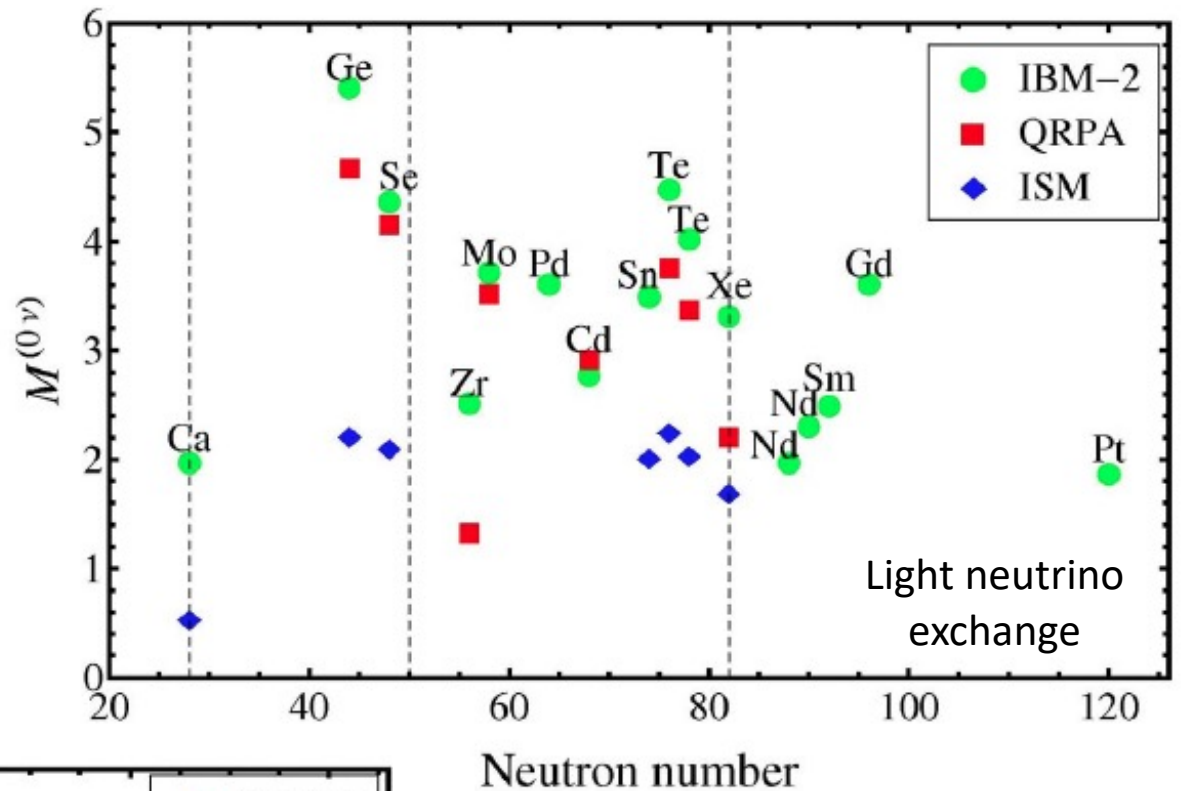
# 0ν double beta decay

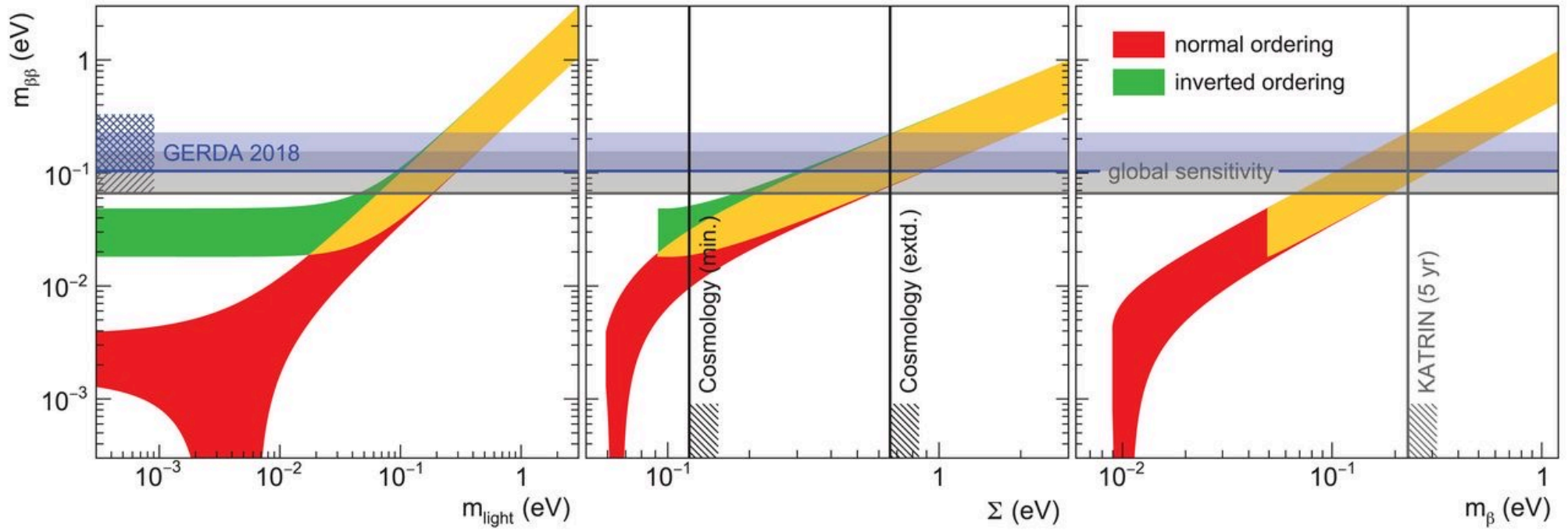
$$(1/T_{1/2}) = G(E,Z) M^2 \langle m_{\beta\beta} \rangle^2$$

$G(E,Z)$  : phase space

$M$  : nuclear matrix element

$$\langle m_{\beta\beta} \rangle = \left| \sum_j |U_{ej}|^2 m_j e^{i\delta(j)} \right|$$





GERDA Collaboration, Science **365**, 1445 (2019)

$$m_{\beta\beta} = \left| \sum_{i=1}^3 U_{ei}^2 m_i \right|$$

$$m_{\beta} = \sqrt{\sum_i |U_{ei}^2| m_i^2}$$

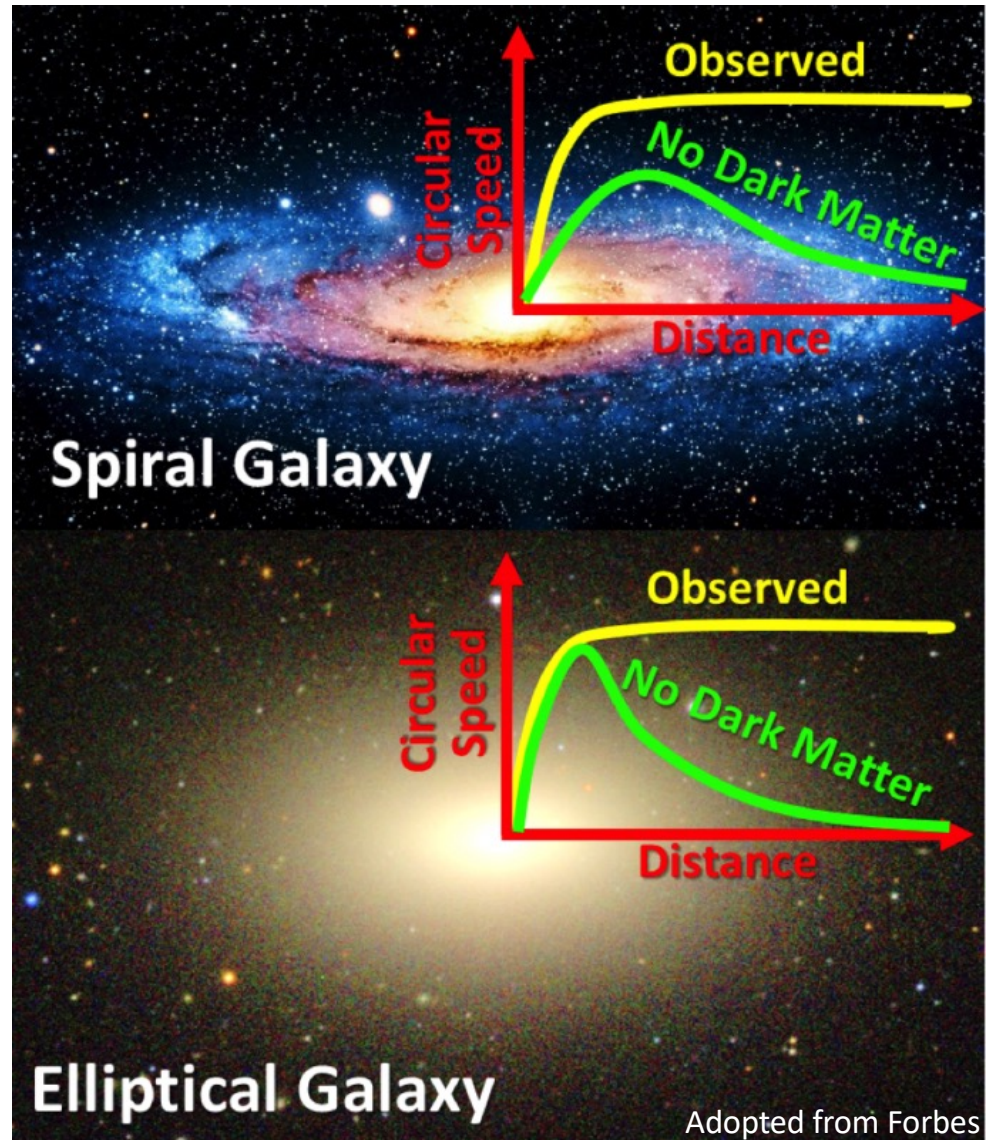
# DARK MATTER



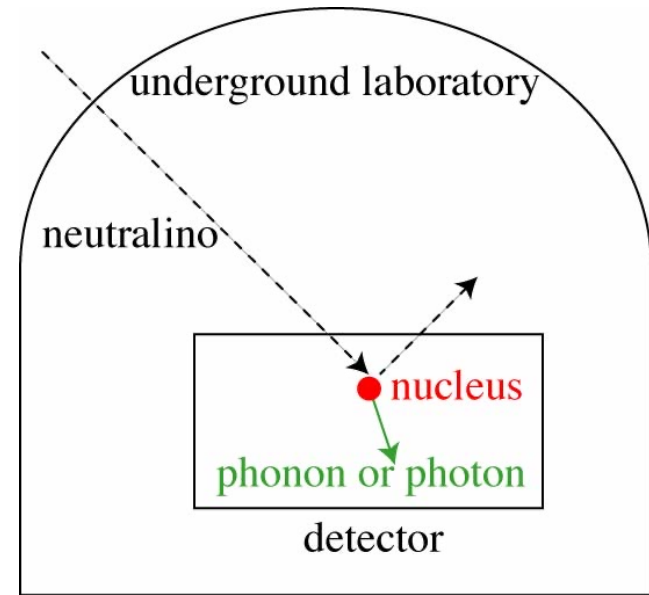
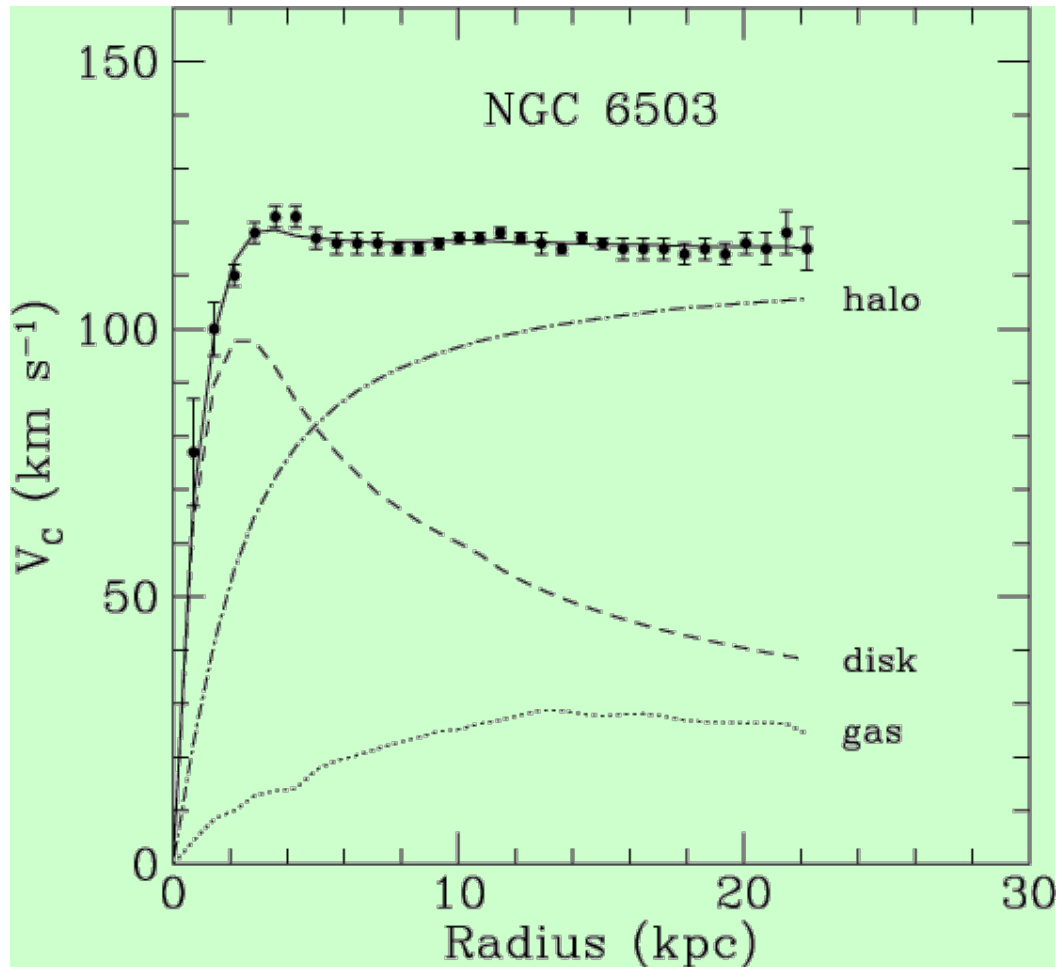
Zwicky



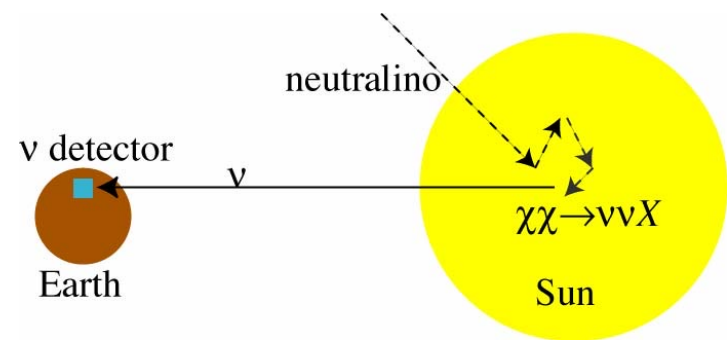
Rubin



# Observing the Dark Matter



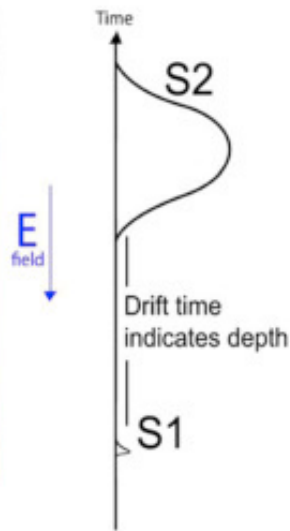
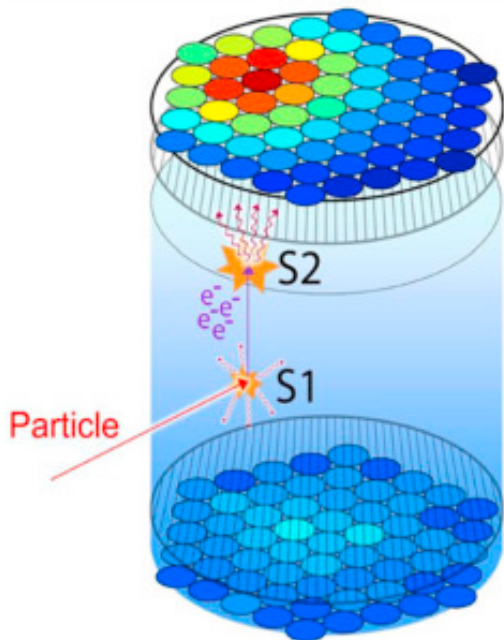
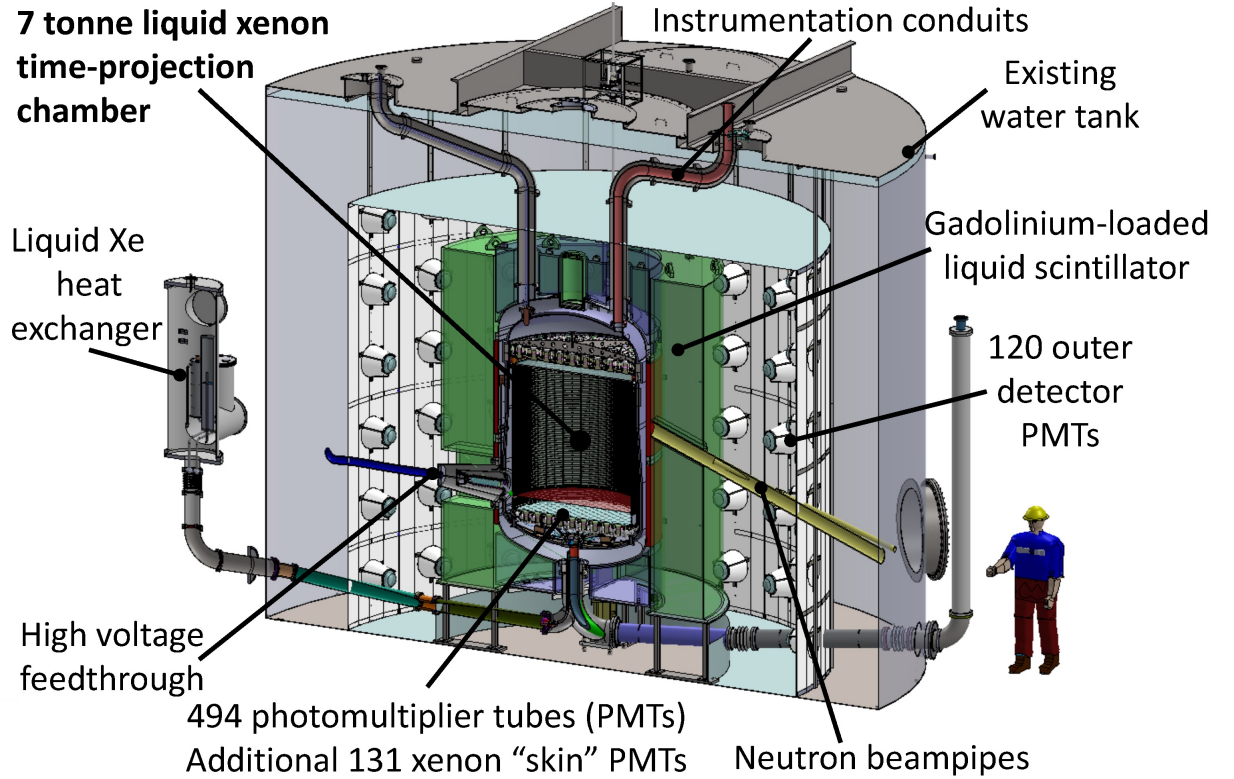
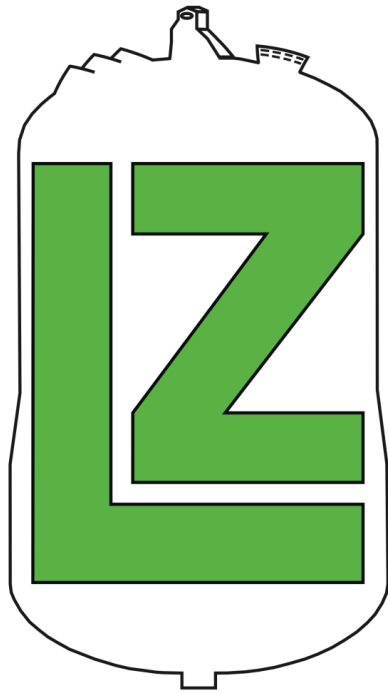
Direct detection





Indirect detection

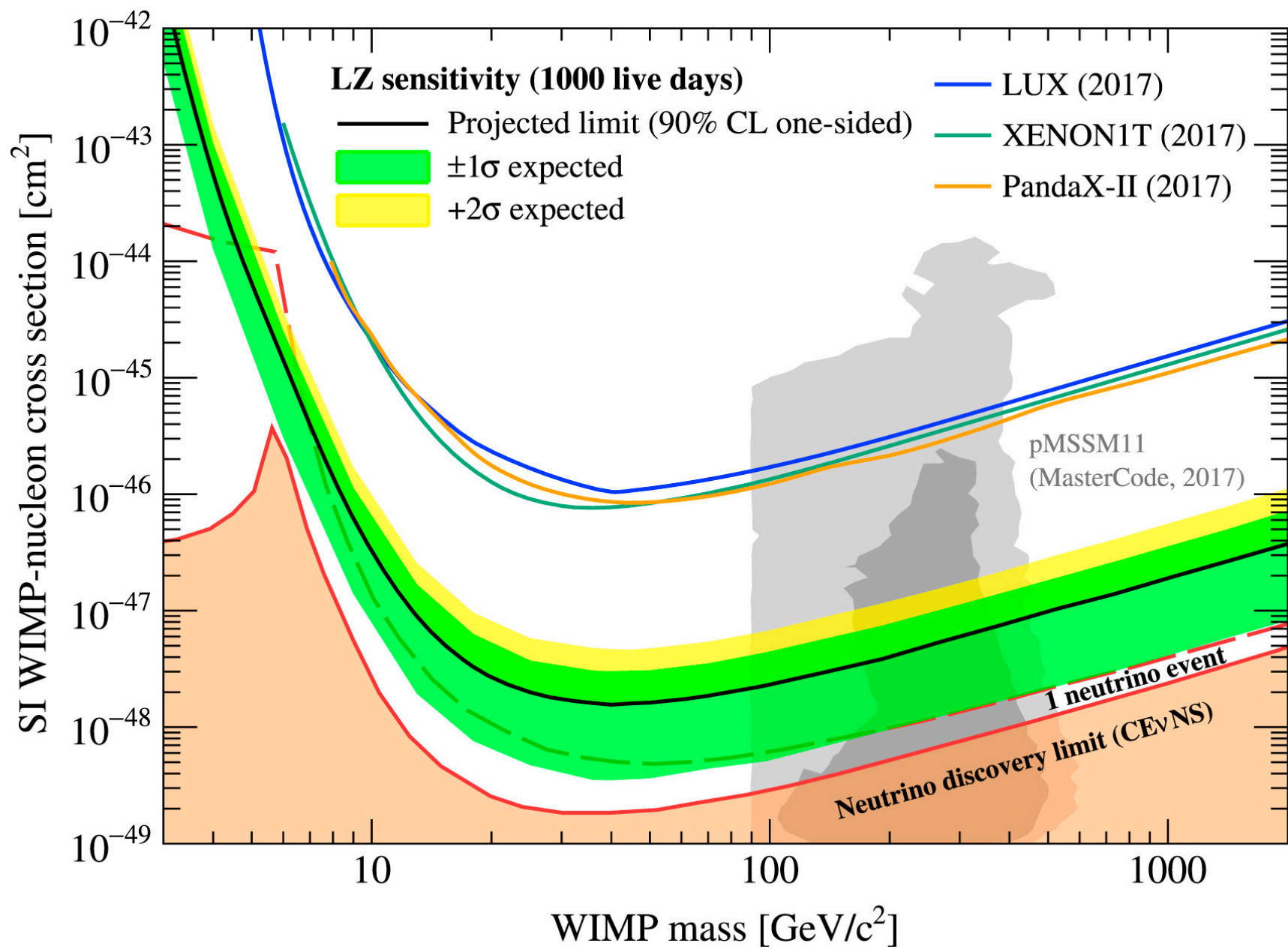


# The LZ Detector



 ionization electrons  
 UV scintillation photons (~175 nm)

*Nucl.Instrum.Meth.A* 936 (2019) 162-165



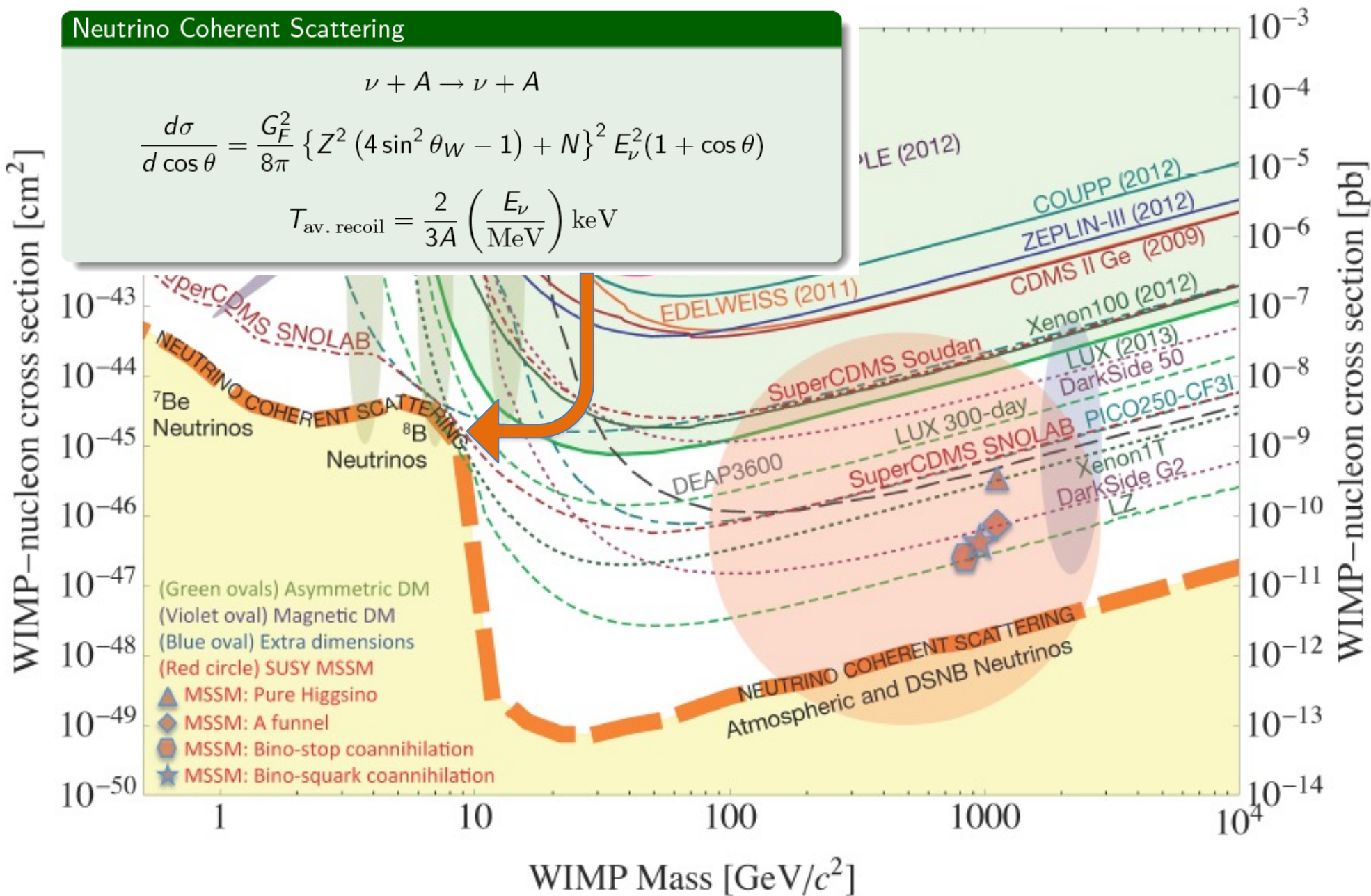
SuperCDMS Soudan Low Threshold  
 (VERSION 10.03/2012)

Neutrino Coherent Scattering

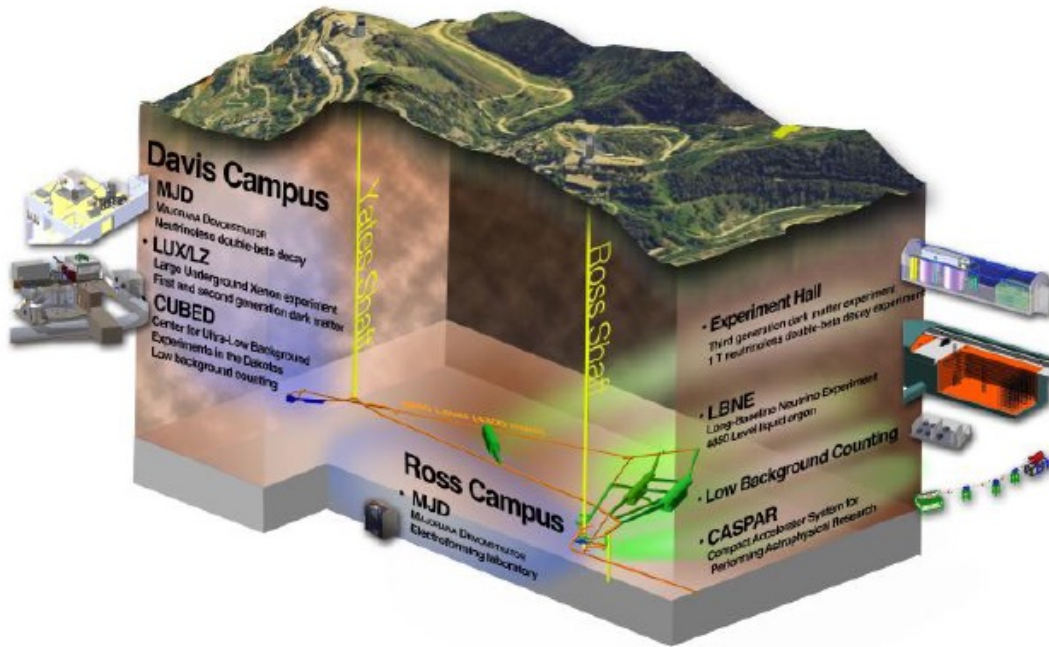
$$\nu + A \rightarrow \nu + A$$

$$\frac{d\sigma}{d\cos\theta} = \frac{G_F^2}{8\pi} \{Z^2 (4\sin^2\theta_W - 1) + N\}^2 E_\nu^2 (1 + \cos\theta)$$

$$T_{\text{av. recoil}} = \frac{2}{3A} \left( \frac{E_\nu}{\text{MeV}} \right) \text{keV}$$



# Compact Accelerator System for Performing Astrophysical Research (CASPAR)

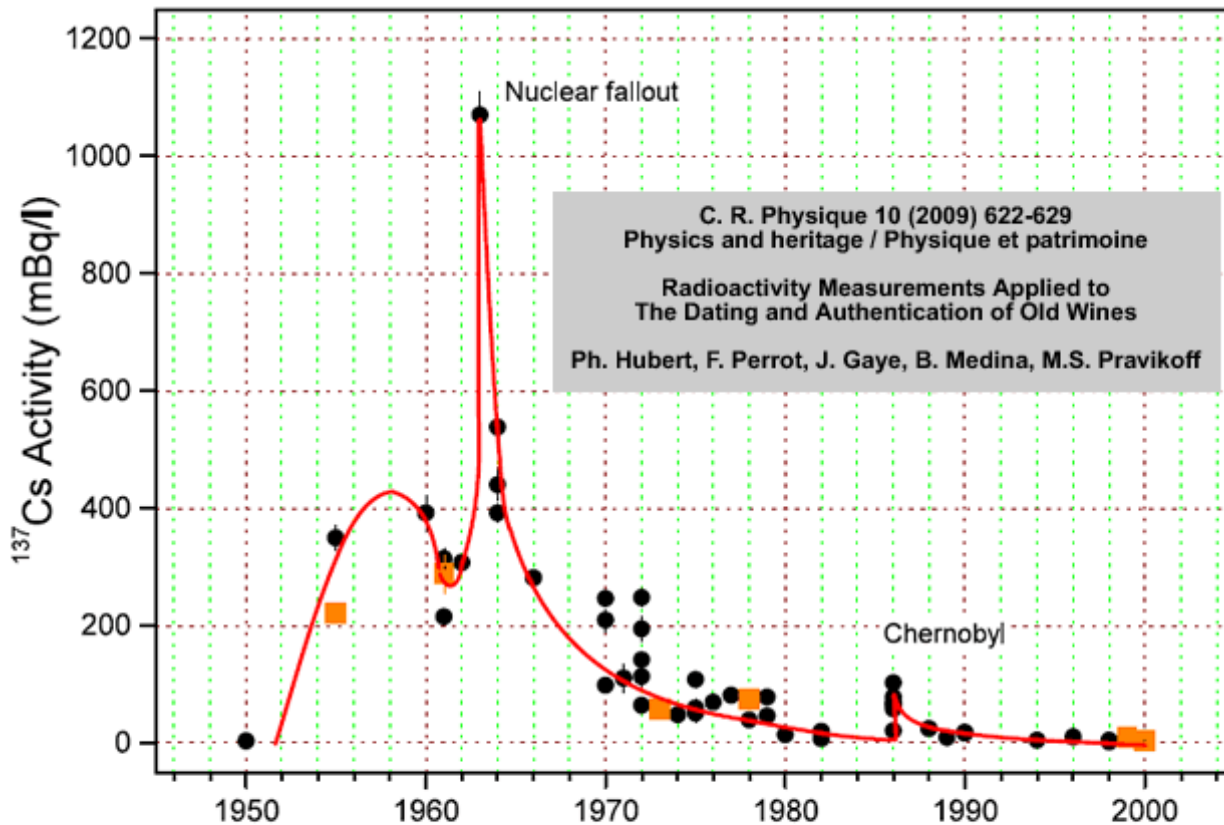


LUNA: Stellar Hydrogen burning  
CASPAR: Helium burning in  
massive Red Giant stars or low-  
mass AGB stars  
 $^{13}\text{C}(\alpha, n)^{16}\text{O}$   
 $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

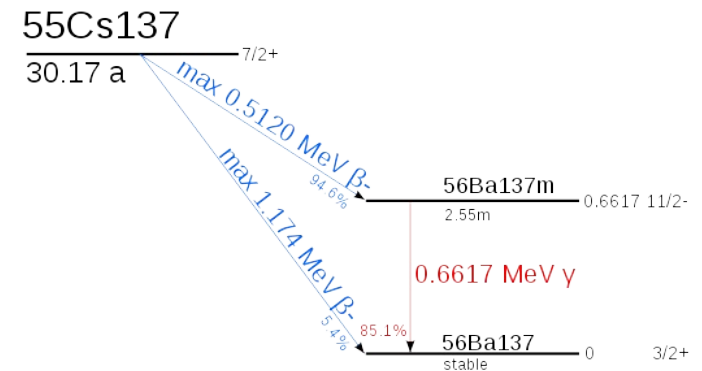


Accelerator vault 2015 after installation

# Sources of neutrinos: Wine



Activity (mBq/l) of the  $^{137}\text{Cs}$  radioactive isotope as a function of the wine vintage. All activities are for Bordeaux wines only, and are normalized to an arbitrary date, January 1st, 2000. The solid circles correspond to measurements after reduction of the wine into ashes, the orange squares correspond to nondestructive measurements, i.e., without opening the bottles. Statistical errors are generally smaller than the dimension of the points.



Thank you

