

Supernova Bounds on a Neutrinophilic Dark Sector

Christopher Cappiello

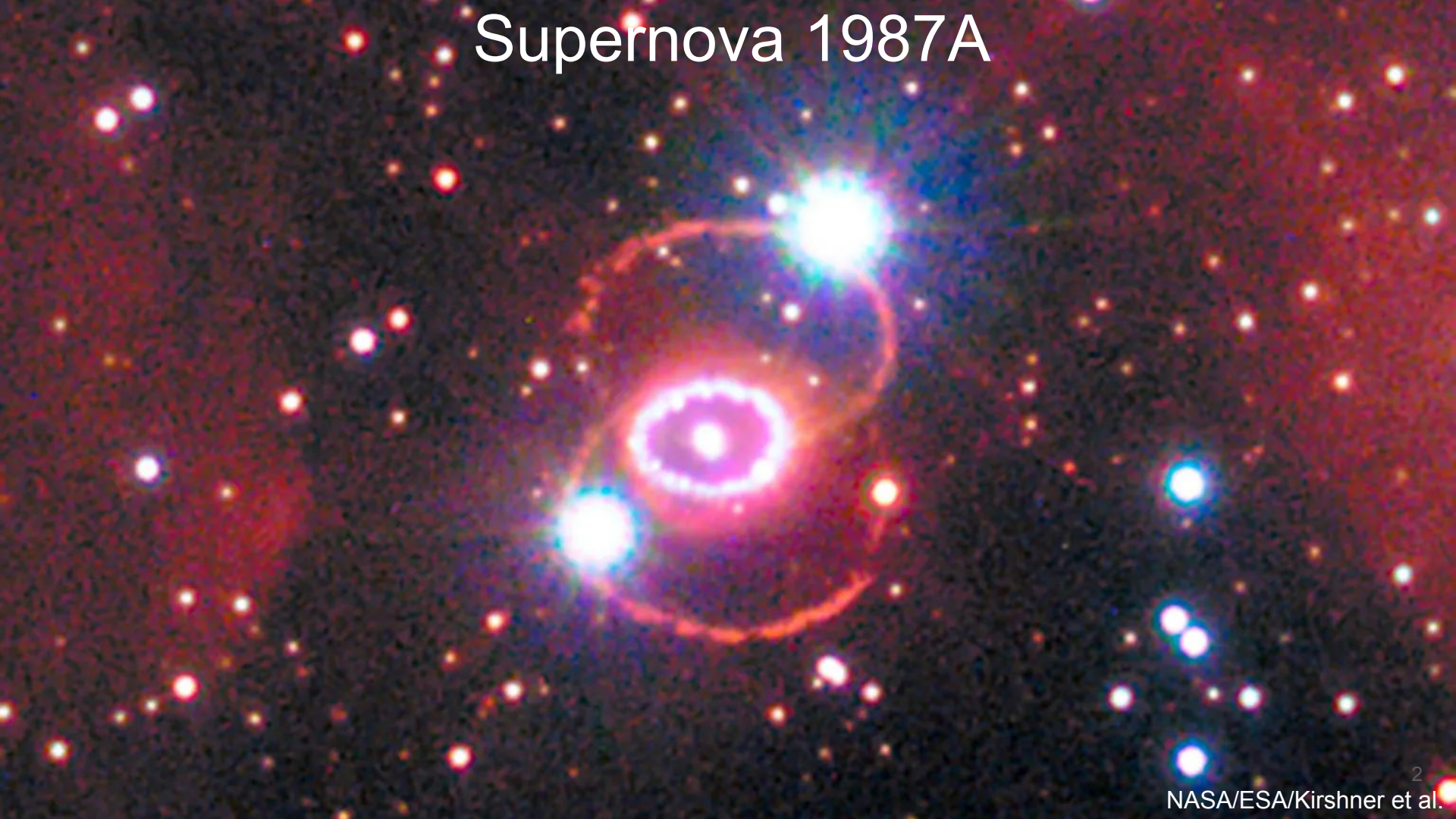
**Based on:
arXiv:2503.09691 + upcoming work
CC, Bhupal Dev, Amol Patwardhan**

CETUP*, July 13, 2026



WashU

Supernova 1987A



Supernova 1987A Constraints on New Particles

Volume 193, number 4

PHYSICS LETTERS B

23 July 1987

CONSTRAINTS ON LIGHT PARTICLES FROM SUPERNOVA SN 1987A

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Received 23 April 1987

We show how the recent observations of neutrinos from SN 1987A can be used to constrain the properties of light particles, including the number of neutrino species N_ν and the coupling strength of a light pseudoscalar boson such as the axion. First we demonstrate the consistency of Kamiokande II and IMB III data with each other and with the standard theory of neutrino emission from a supernova core of radius $R \approx 20$ km and effective temperature $T \approx 5$ to 6 MeV. The total amount of energy emission by neutrinos is compatible with the binding energy of a neutron star only if $N_\nu \leq 6$. The upper limit on the difference between the binding energy and the energy lost via $N_\nu = 3$ neutrino species excludes light pseudoscalar bosons with couplings to electrons in the range $10^{-27} \leq \alpha_e \leq 10^{-24}$, which enables one to extend astrophysical bounds on the axion decay constant to $f_a \geq 4 \times 10^9$ GeV.

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2 MAY 1988

Bounds on Exotic-Particle Interactions from SN1987A

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(Received 26 October 1987)

The observation of a neutrino pulse from the supernova SN1987A constrains the production of light exotic particles in the proto neutron star. We derive a new bound on the axion decay constant, $f_a \geq 10^{10}$ GeV. If right-handed (RH) neutrinos exist, the "RH Fermi constant" is $G_{RH} \leq 10^{-4} G_F$, 2 orders of magnitude below laboratory bounds. The Dirac mass of ν_μ can be constrained below laboratory limits.

PACS numbers: 97.60.Bw, 14.60.Gh, 14.80.Gt, 97.60.Jd

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PHYSICAL REVIEW LETTERS

2 MAY 1988

Axions from SN1987A

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(Received 12 November 1987)

Axion emission from SN1987A by nucleon-nucleon axion bremsstrahlung is considered. On the basis of the neutrino observations the axion luminosity must be $\lesssim 10^{33}$ erg s $^{-1}$. This occurs if (1) axions couple very weakly: $m_a \lesssim 0.75 \times 10^{-3}$ eV; or (2) axions couple strongly enough to be "trapped" and radiated from an "axion sphere" with $T_a \lesssim 8$ MeV: $m_a \gtrsim 2.2$ eV. In general, "axion trapping" occurs for $m_a \gtrsim 0.016$ eV. These mass constraints are at best reliable to within a factor of ≈ 3 .

PACS numbers: 97.60.Bw, 14.80.Gt, 97.60.Jd

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24 March 1988

CONSTRAINTS ON AXIONS FROM SN 1987A

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and

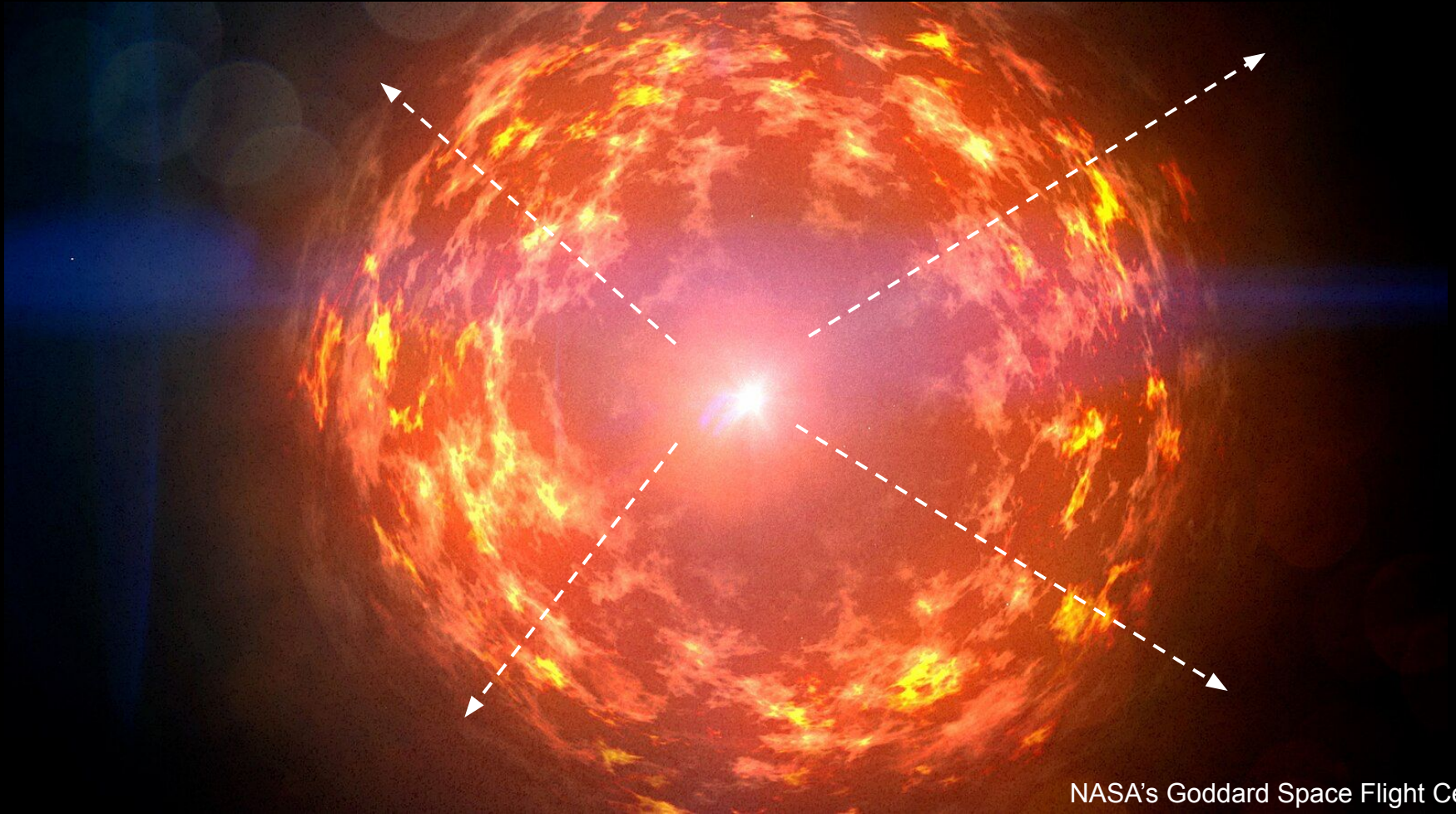
Gary STEIGMAN

Department of Physics, Ohio State University, Columbus, OH 43210-1106, USA

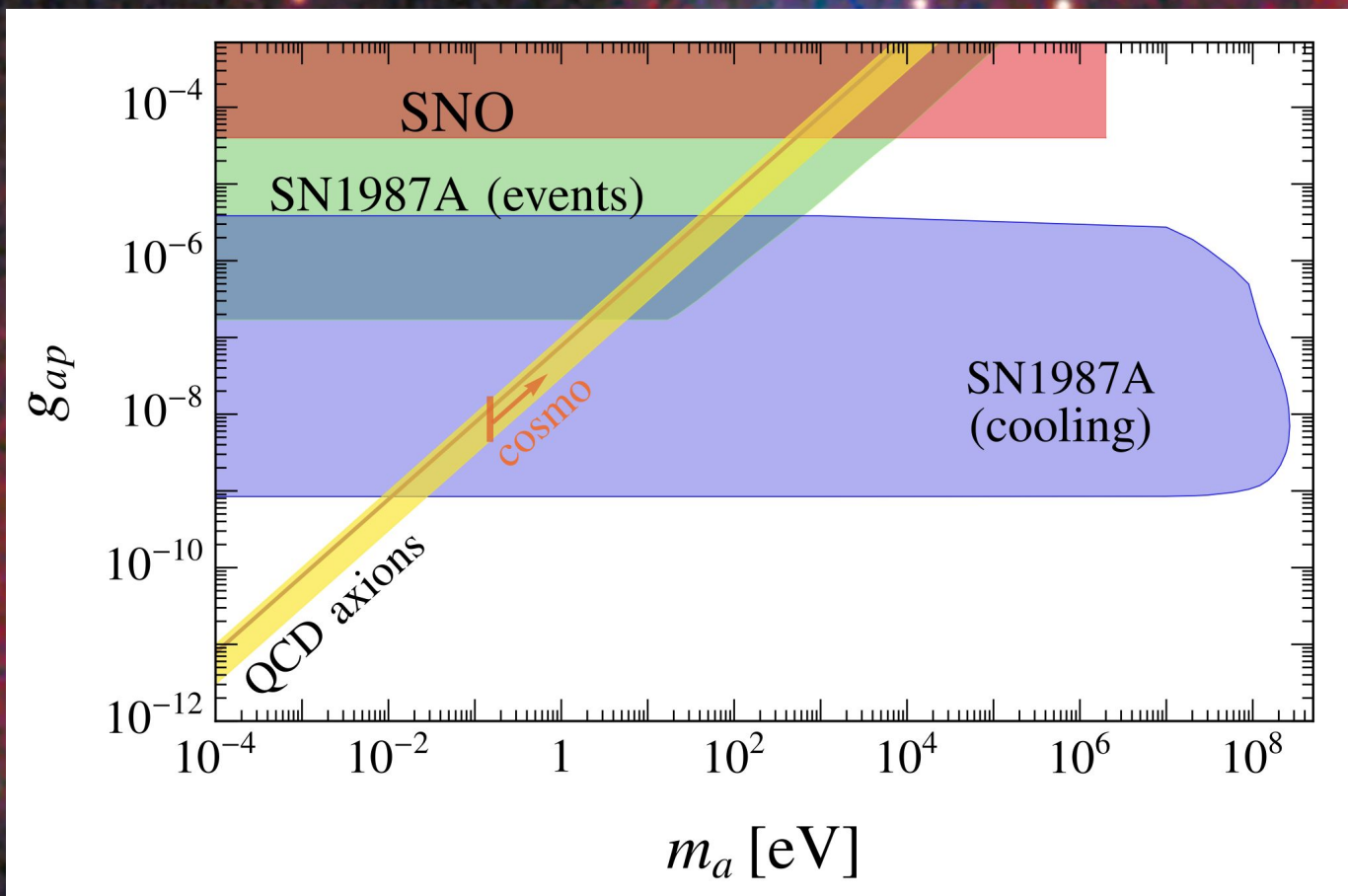
Received 23 December 1987

We argue that the detection of neutrinos from SN 1987A implies that axion emission from the collapsing star must not have been significant. The best axion supernova limits come from nucleon bremsstrahlung processes which involve axion-quark couplings rather than the electron or photon couplings used in obtaining axion limits from red giants. The axion-quark couplings are less sensitive to specific axion model parameters and therefore are capable of yielding a model independent lower bound on f_a . Quantitatively this yields an axion coupling limit of $f_a \gtrsim 8 \times 10^{11}$ GeV (or equivalently $m_a \lesssim 9 \times 10^{-3}$ eV). This approaches the upper bound on $f_a \lesssim 4 \times 10^{12}$ GeV from cosmological density arguments. While the bounds do not exclude the axion, if the axion exists, it must dominate the cosmological mass density. This same argument can also be applied to analogous couplings for any particle with $m \lesssim 50$ MeV.

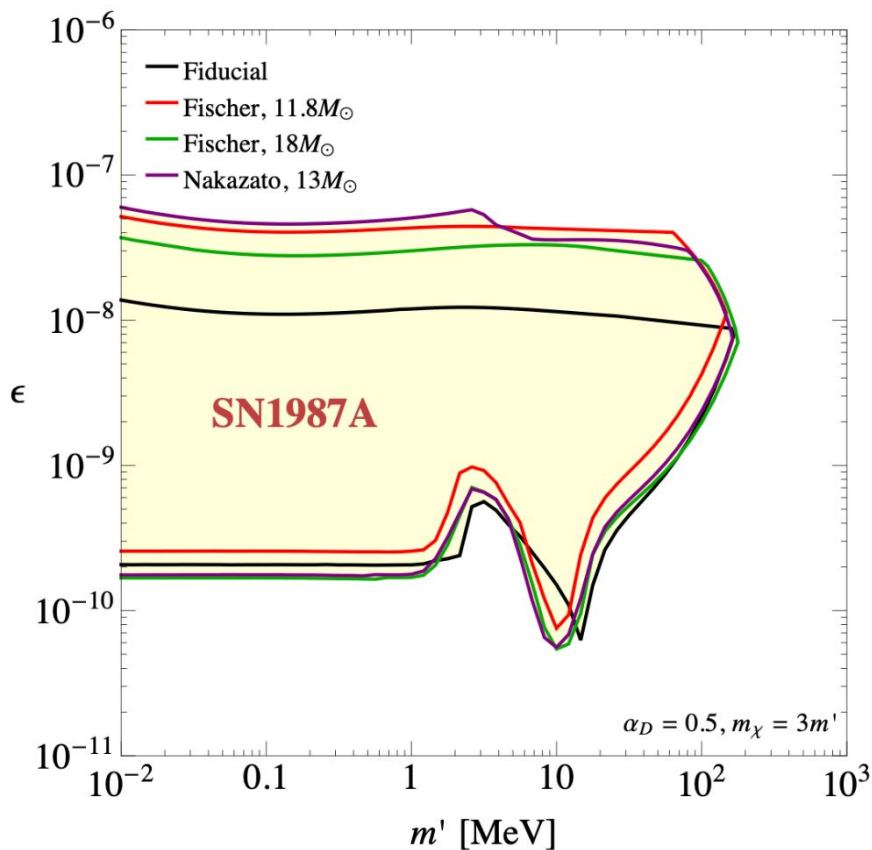
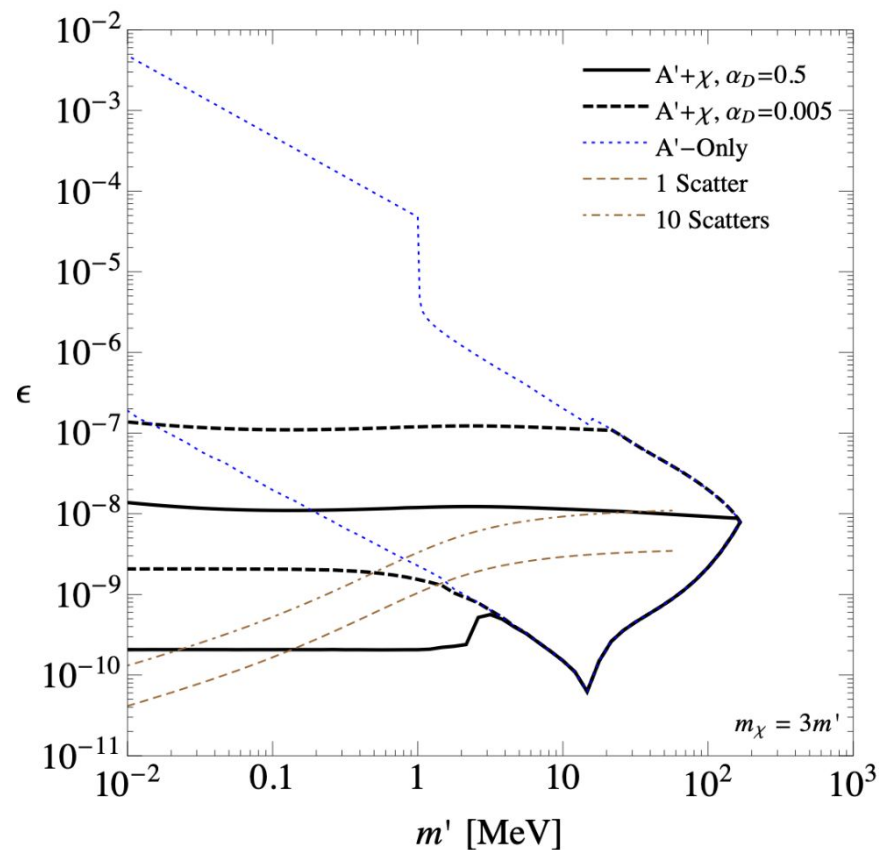
Supernova 1987A Constraints on New Particles



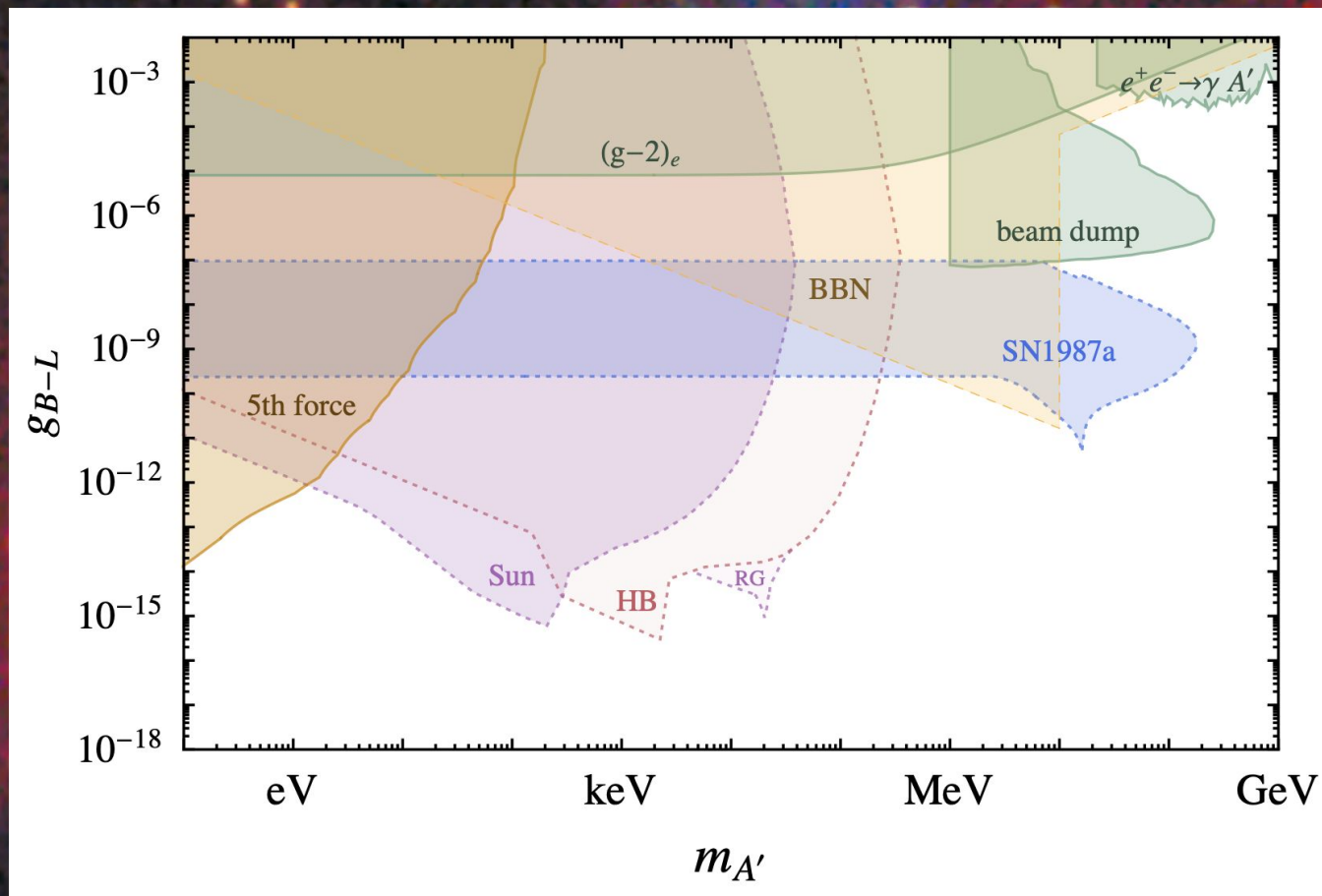
Supernova 1987A Constraints on New Particles



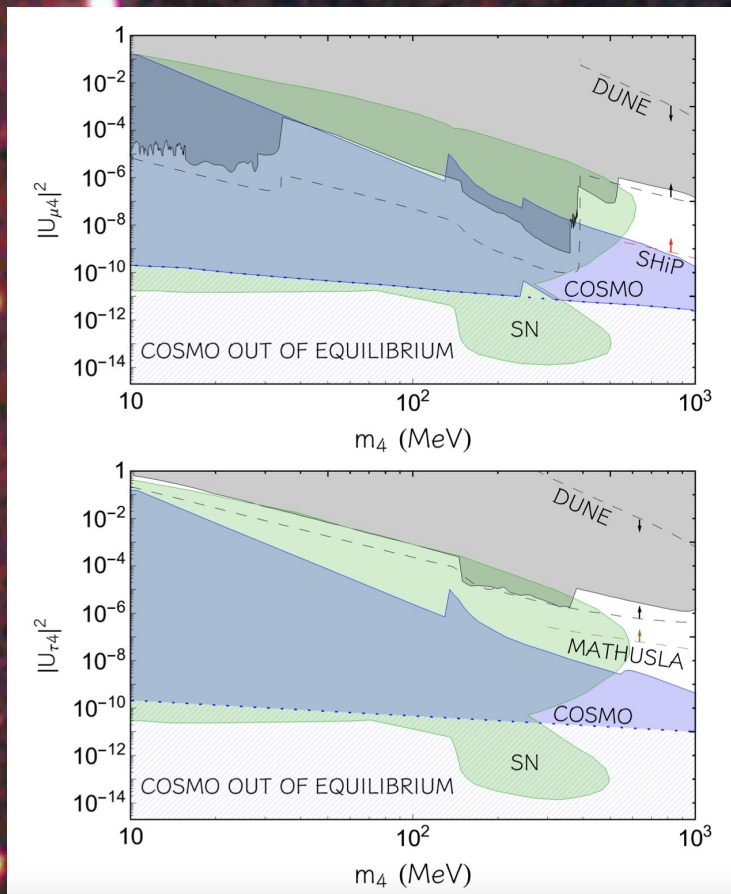
Supernova 1987A Constraints on New Particles



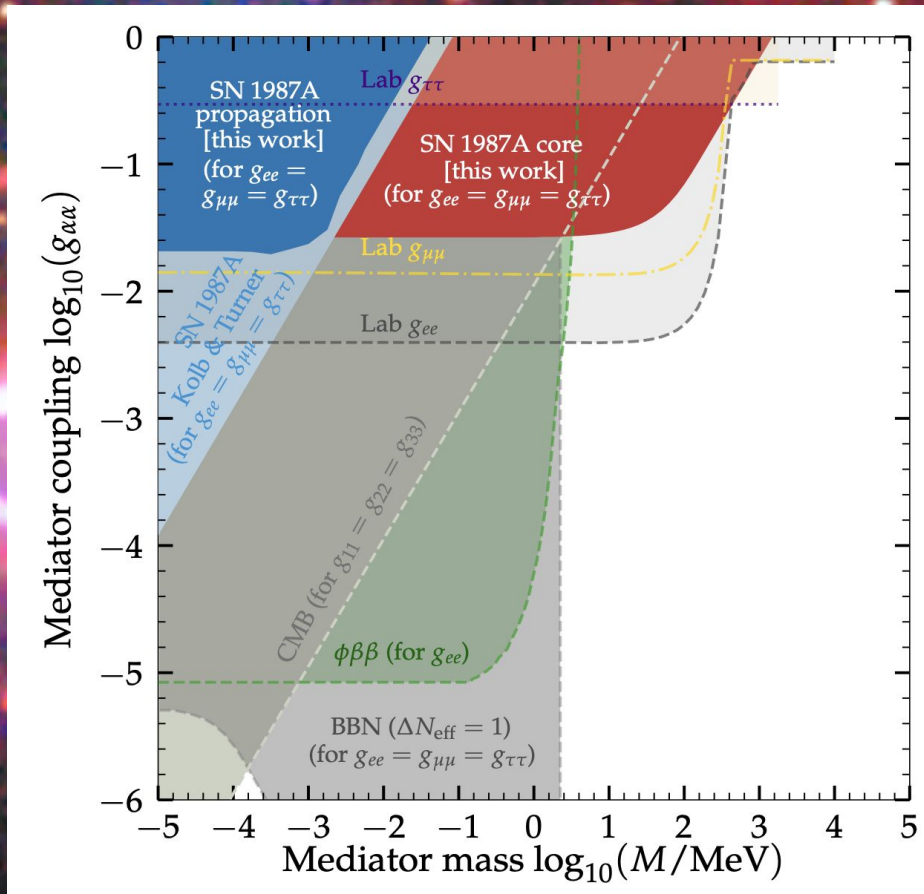
Supernova 1987A Constraints on New Particles



Supernova 1987A and Neutrino Physics

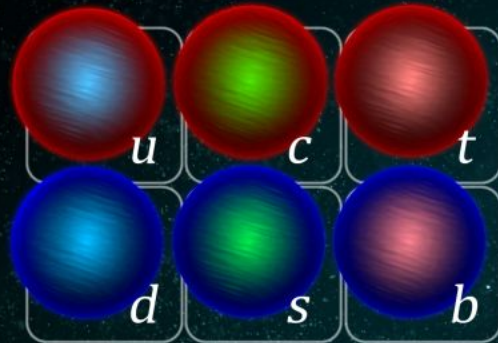


<https://arxiv.org/pdf/2311.00033>

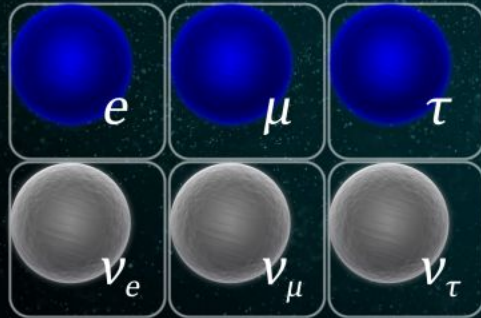


<https://arxiv.org/pdf/1912.09115>

Neutrinophilic Dark Matter



Quarks



Leptons

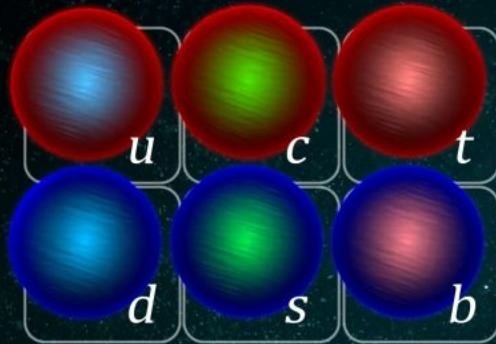


Higgs boson

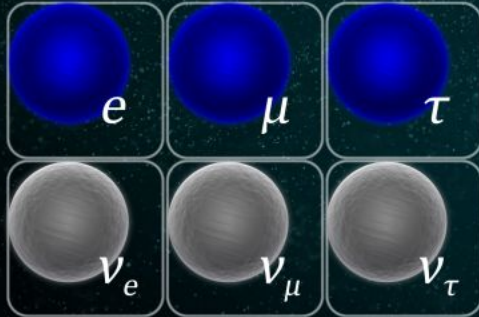


Forces

Neutrinophilic Dark Matter



Quarks



Leptons

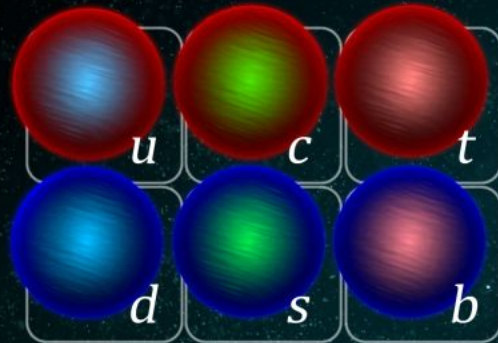


Higgs boson



Forces

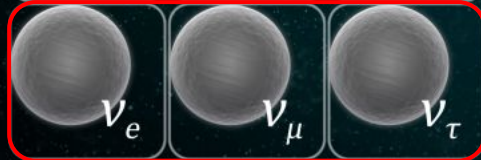
Neutrinophilic Dark Matter



Quarks



Leptons

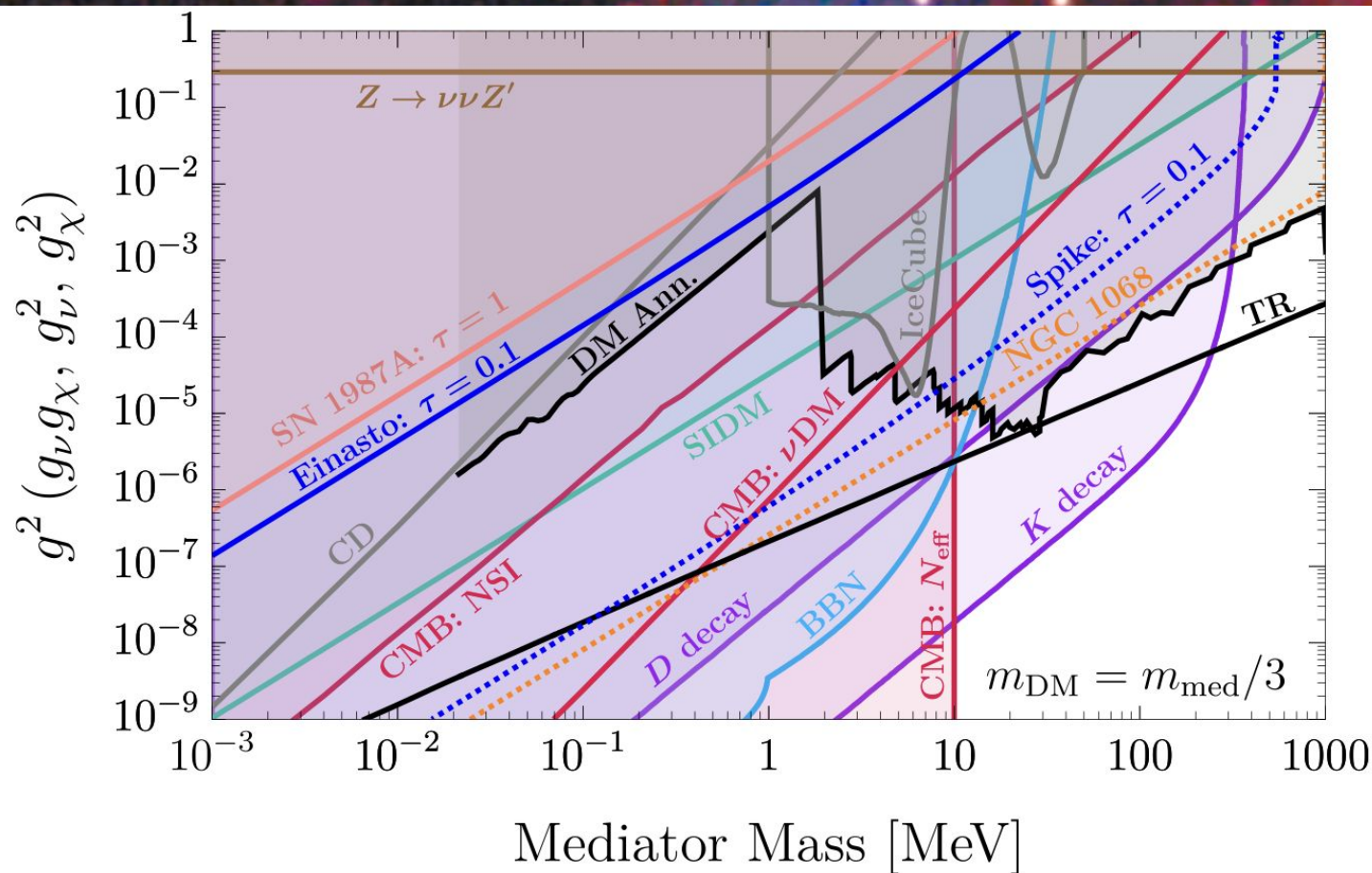


Higgs boson

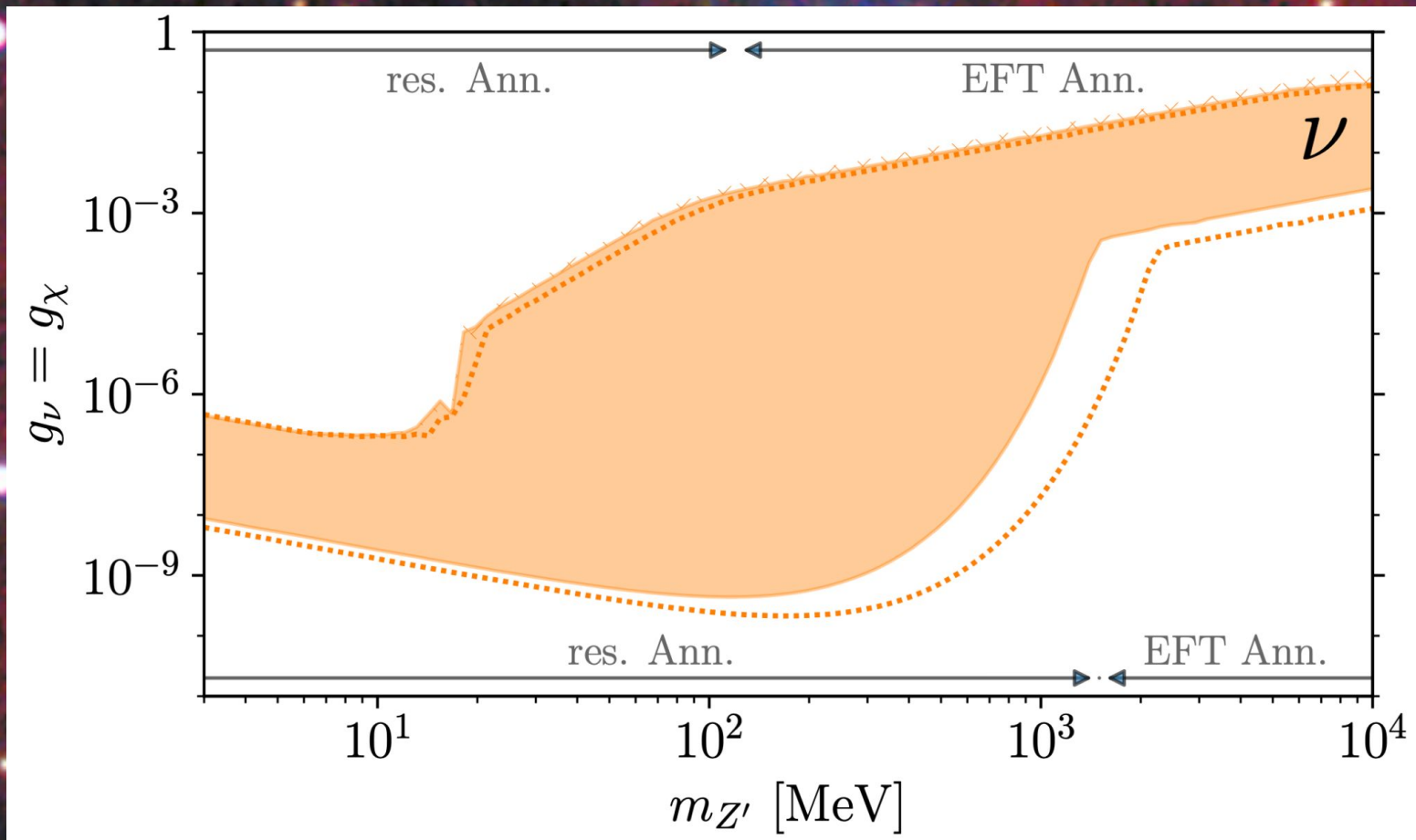


Forces

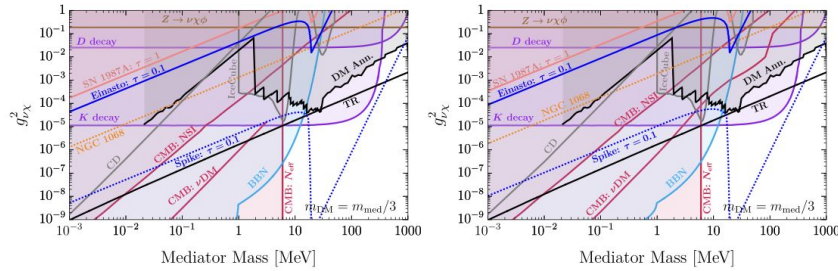
Neutrinophilic Dark Matter



SN Bounds on DM-Neutrino Interactions

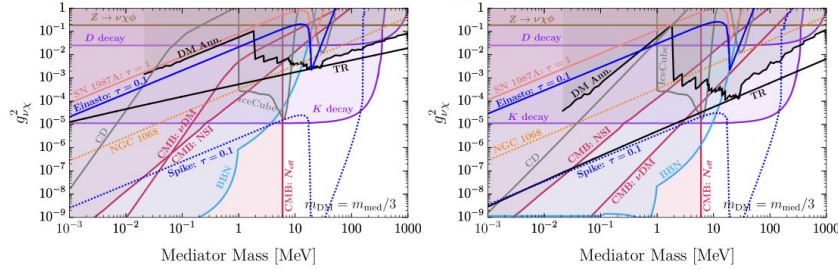


Many Models for DM-Neutrino Interactions!



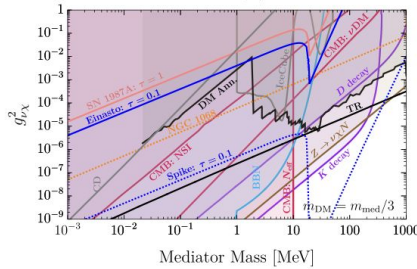
(a) Dirac fermion DM, scalar mediator

(b) Majorana fermion DM, scalar mediator

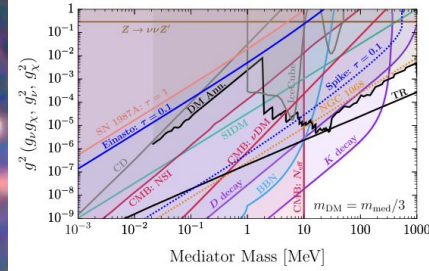


(c) Real Scalar DM, fermion mediator

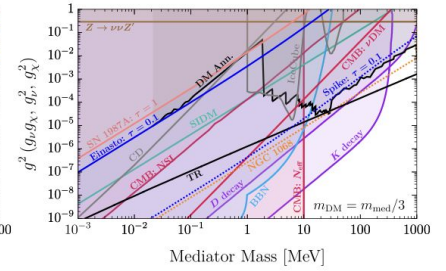
(d) Complex Scalar DM, fermion mediator



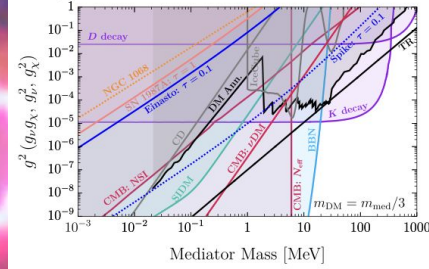
(e) Vector DM, fermionic mediator



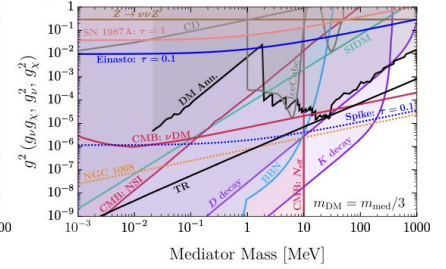
(a) Dirac Fermion DM, vector mediator



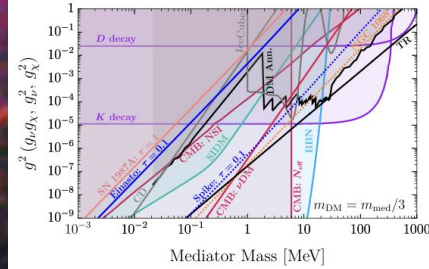
(b) Majorana Fermion DM, vector mediator



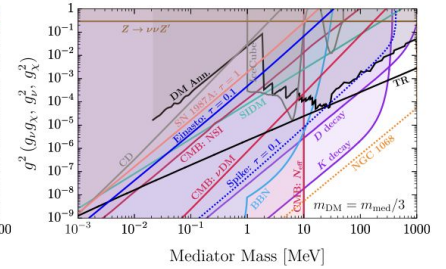
(c) Scalar DM, scalar mediator



(d) Scalar DM, vector mediator.



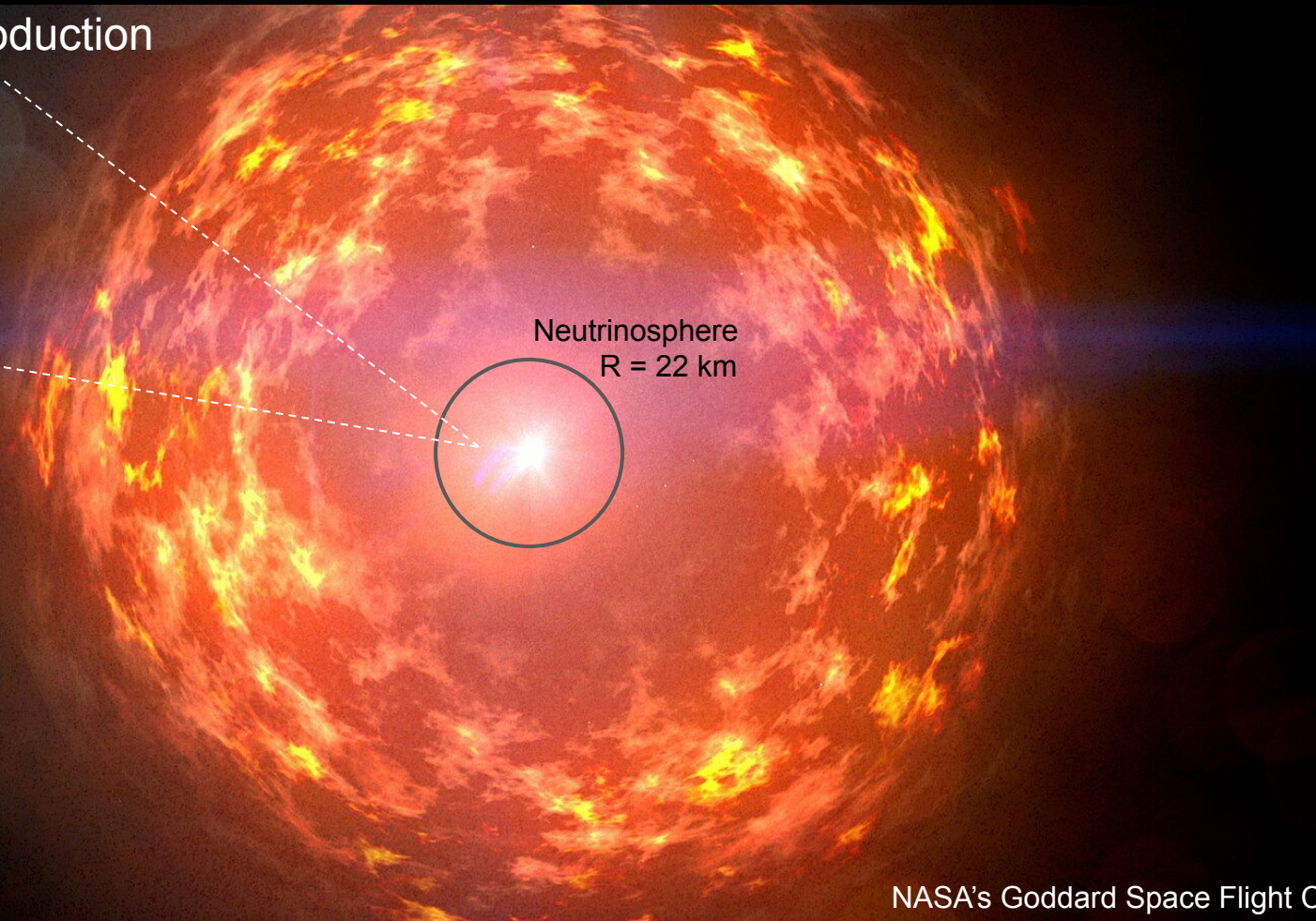
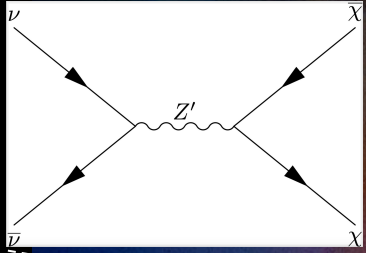
(e) Vector DM, scalar mediator



(f) Vector DM, vector mediator

Constraining Neutrinophilic Dark Sector

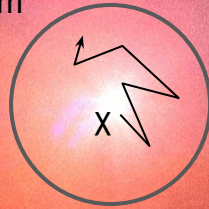
1. Dark Matter Production



Constraining Neutrinophilic Dark Sector

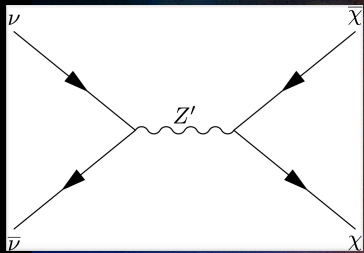
2. Trapping by Neutrinos

Neutrinosphere
 $R = 22 \text{ km}$

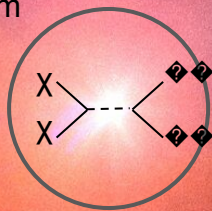


Constraining Neutrinophilic Dark Sector

3. Trapping by Annihilation



Neutrinosphere
 $R = 22 \text{ km}$

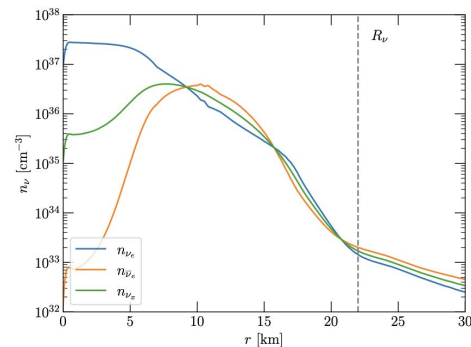
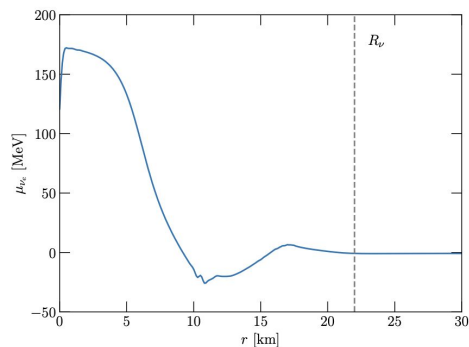
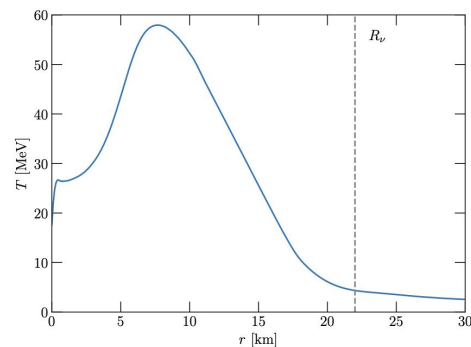
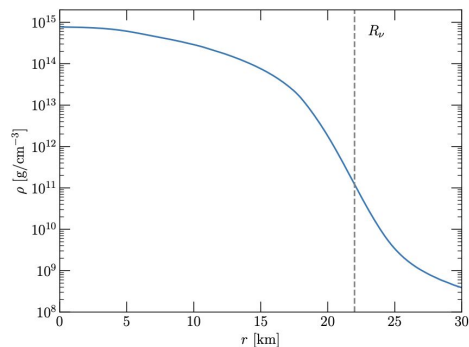


Constraints for a Heavy Mediator

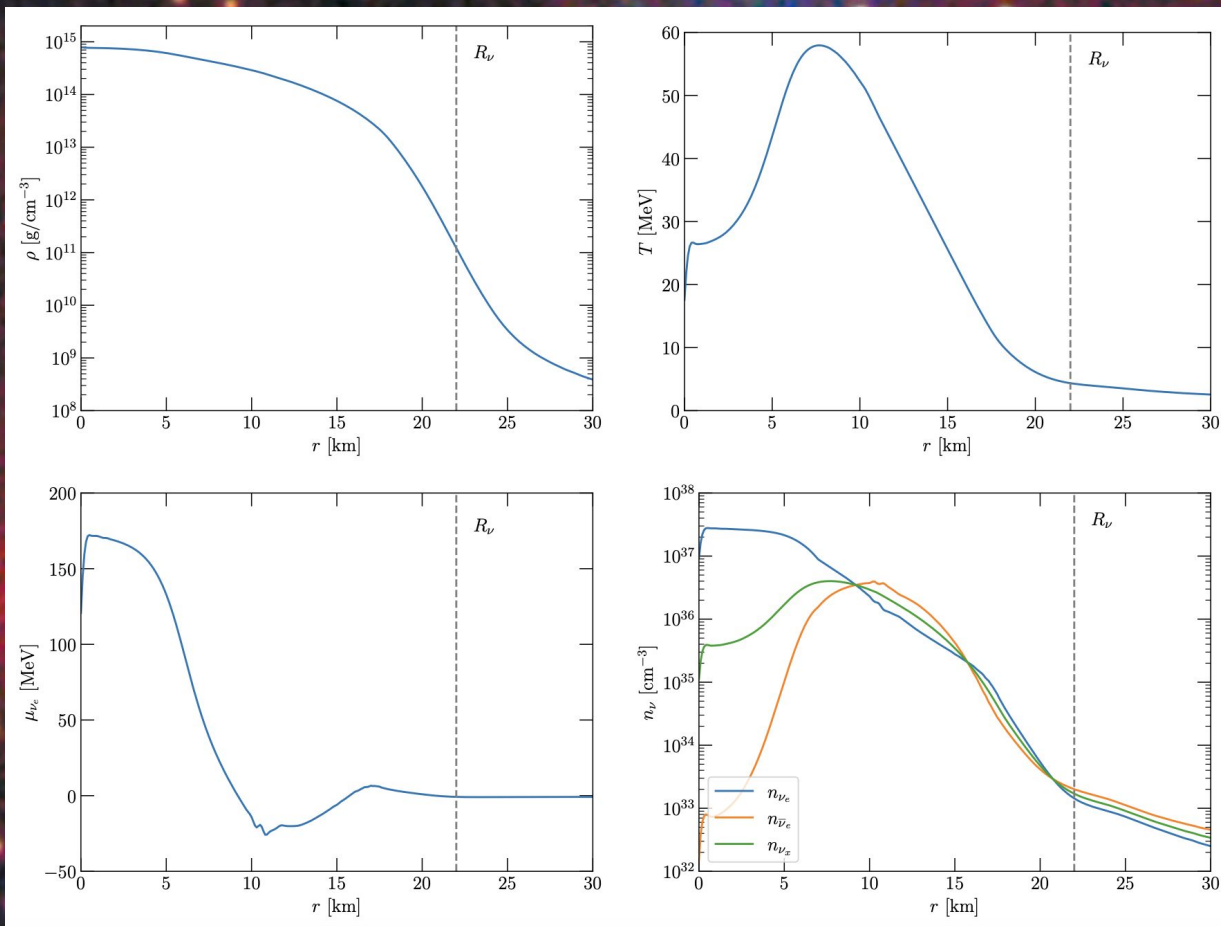
$$-\mathcal{L}_{\text{int}} \supset g_\nu Z'_\mu \bar{\nu} \gamma^\mu P_L \nu + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$

Simulation Inputs: Garching Core
Collapse Supernova Archive

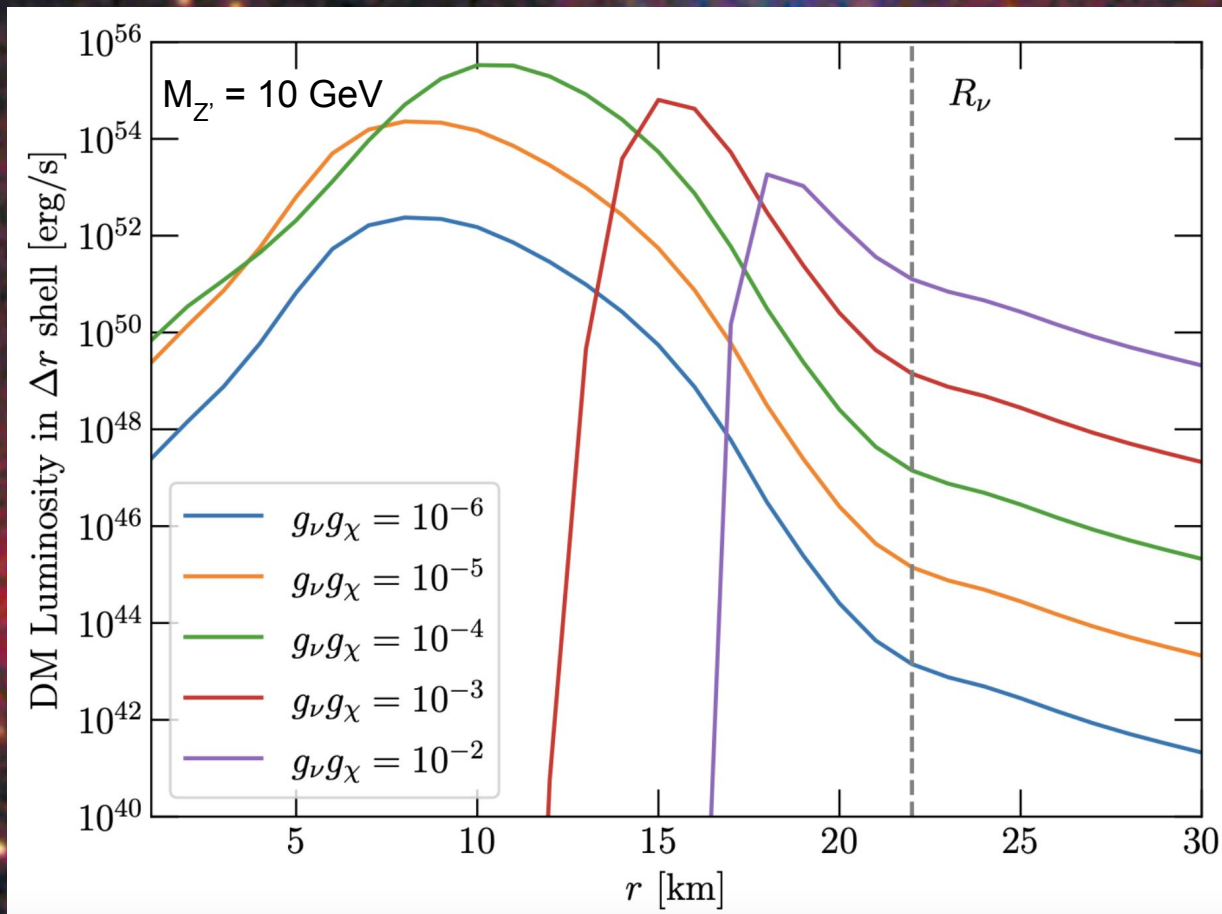
- Credit: Hans-Thomas Janka, Robert Bollig and Lorenz H udepohl
- <https://wwwmpa.mpa-garching.mpg.de/ccsnarchive/archive.html>
- We use a CCSN simulation for 20 solar mass progenitor, snapshot at 1 second post core bounce (PROMETHEUS-VERTEX)



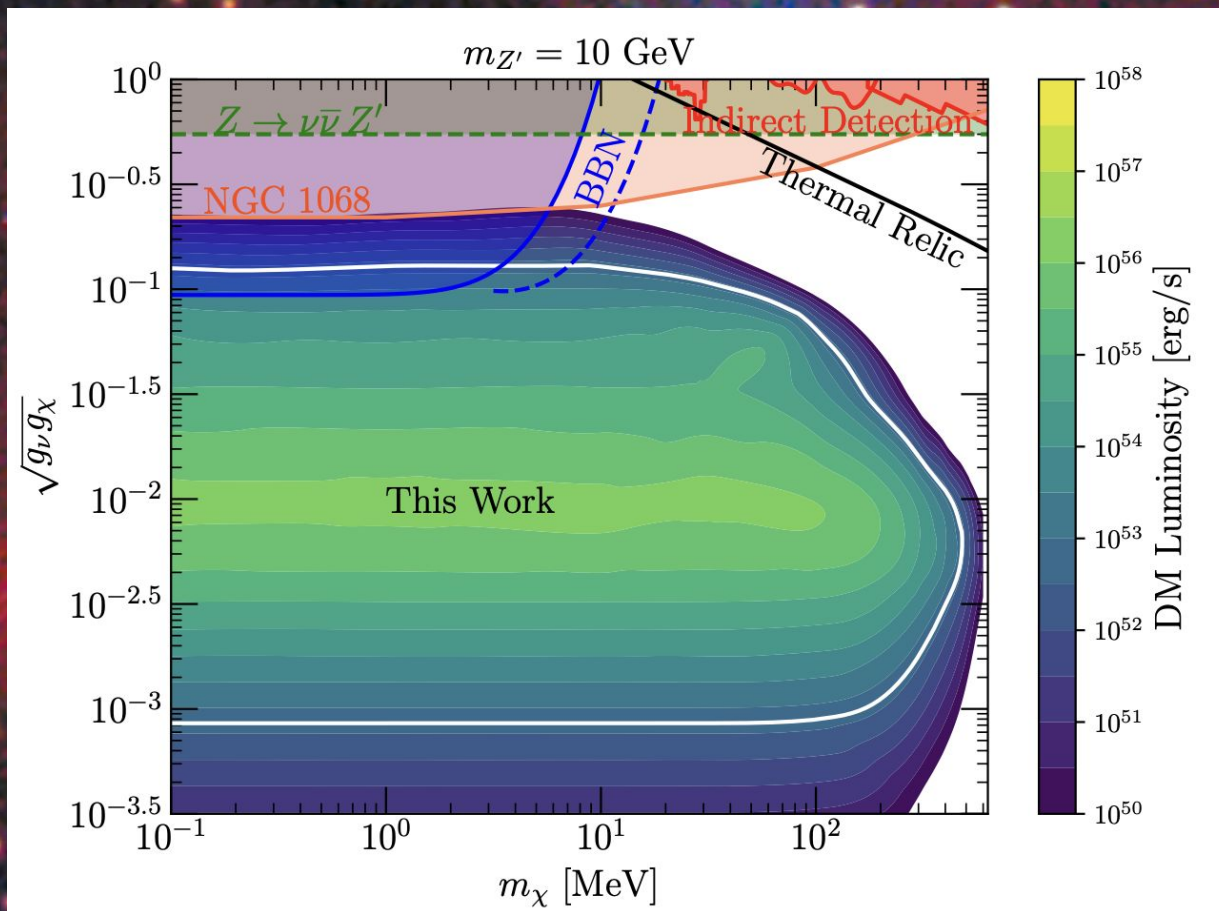
Constraints for a Heavy Mediator



Constraints for a Heavy Mediator



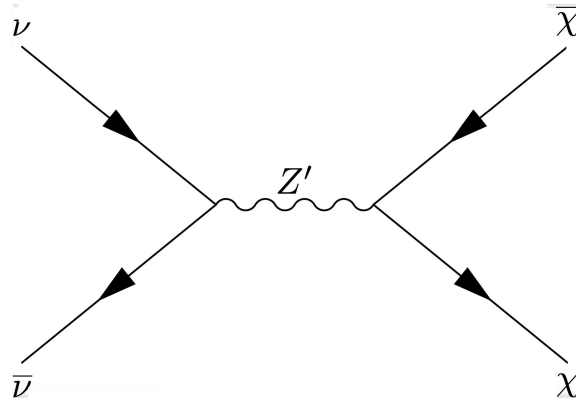
Constraints for a Heavy Mediator



Resonant Dark Matter Production

$$-\mathcal{L}_{\text{int}} \supset g_\nu Z'_\mu \bar{\nu} \gamma^\mu P_L \nu + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$

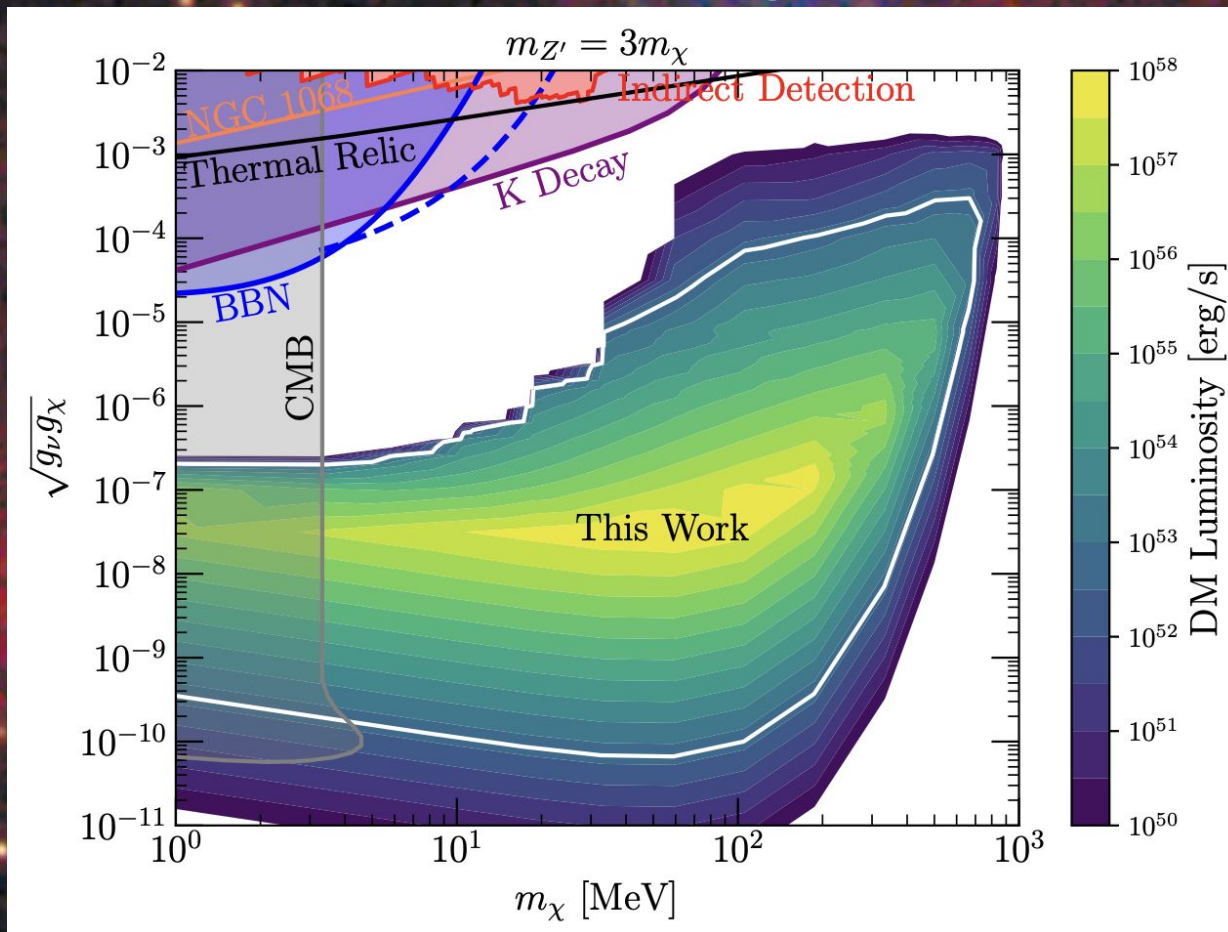
$$\sigma_{\nu_\alpha \bar{\nu}_\alpha \rightarrow \chi \bar{\chi}} = \frac{g_\nu^2 g_\chi^2 (s + 2m_\chi^2)(s - m_\nu^2)}{6\pi s [(s - m_{Z'}^2)^2 + m_{Z'}^2 \Gamma_{Z'}^2]} \sqrt{\frac{s - 4m_\chi^2}{s - 4m_\nu^2}}$$



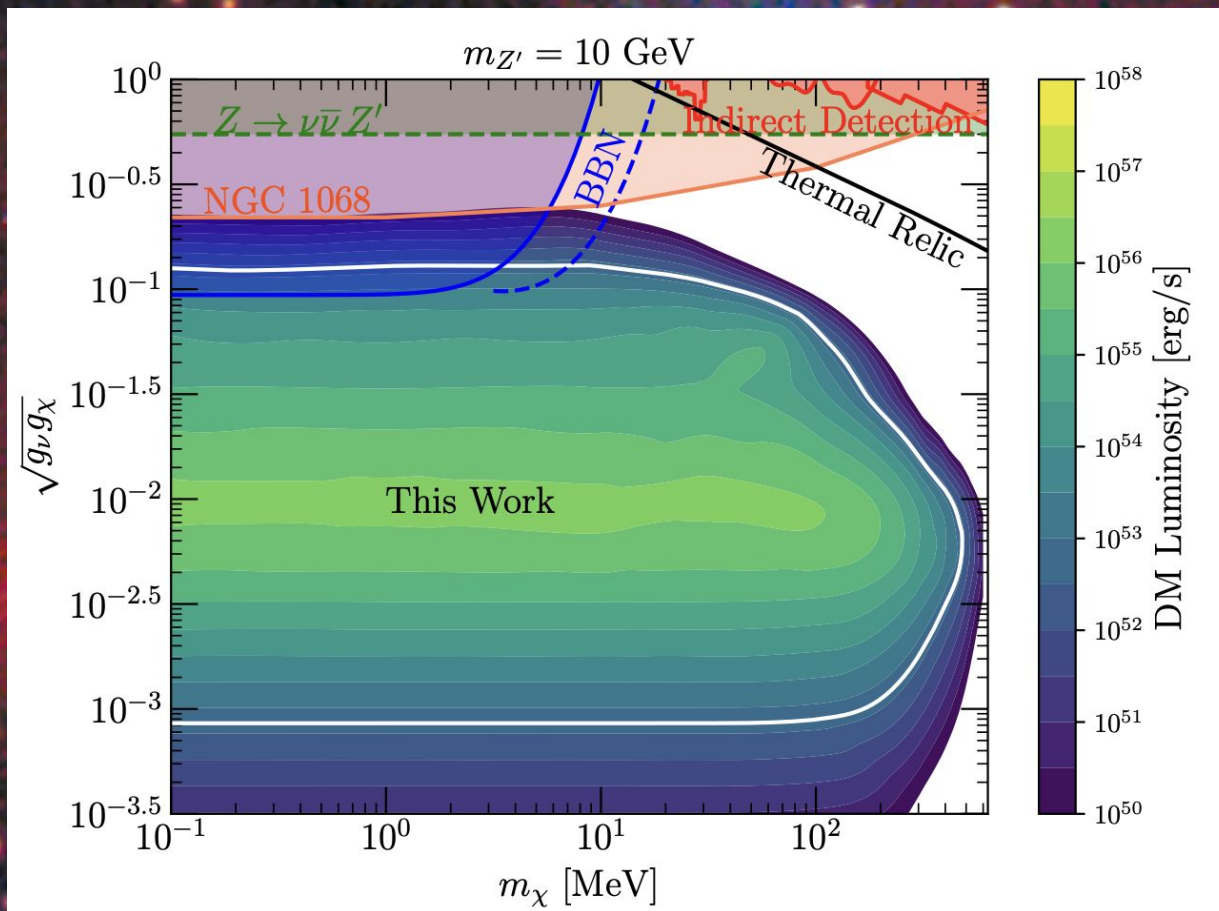
Resonant production

└─→ Enhanced production rate

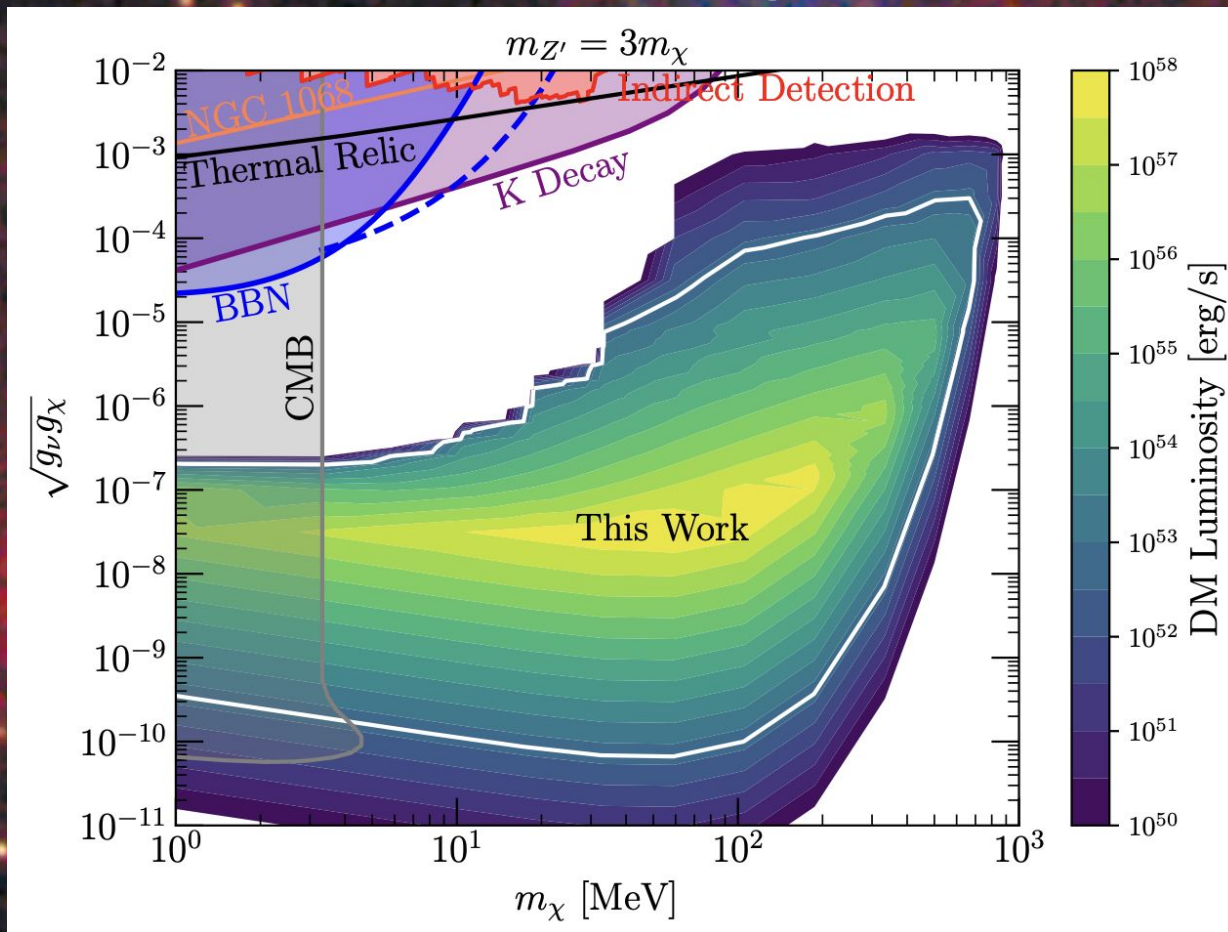
Constraints for a Light Mediator



Constraints for a Heavy Mediator

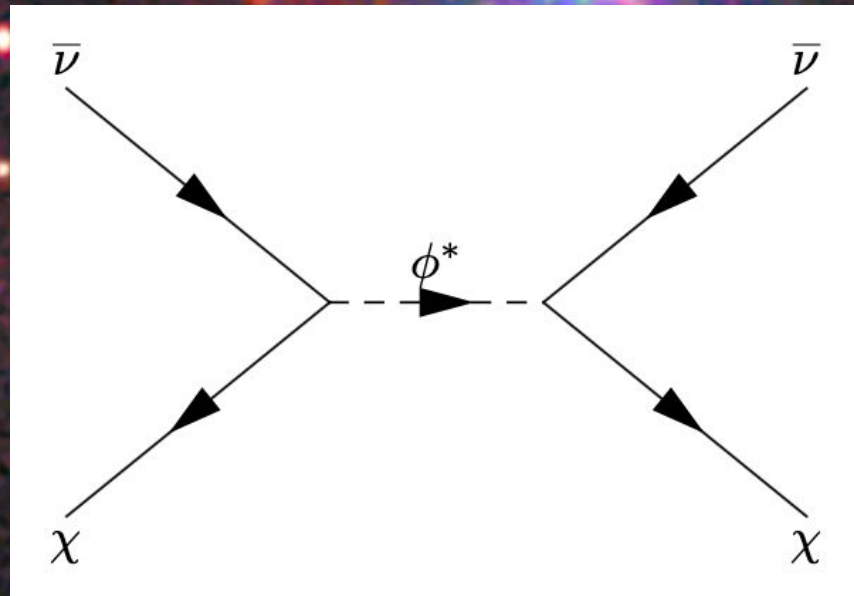


Constraints for a Light Mediator

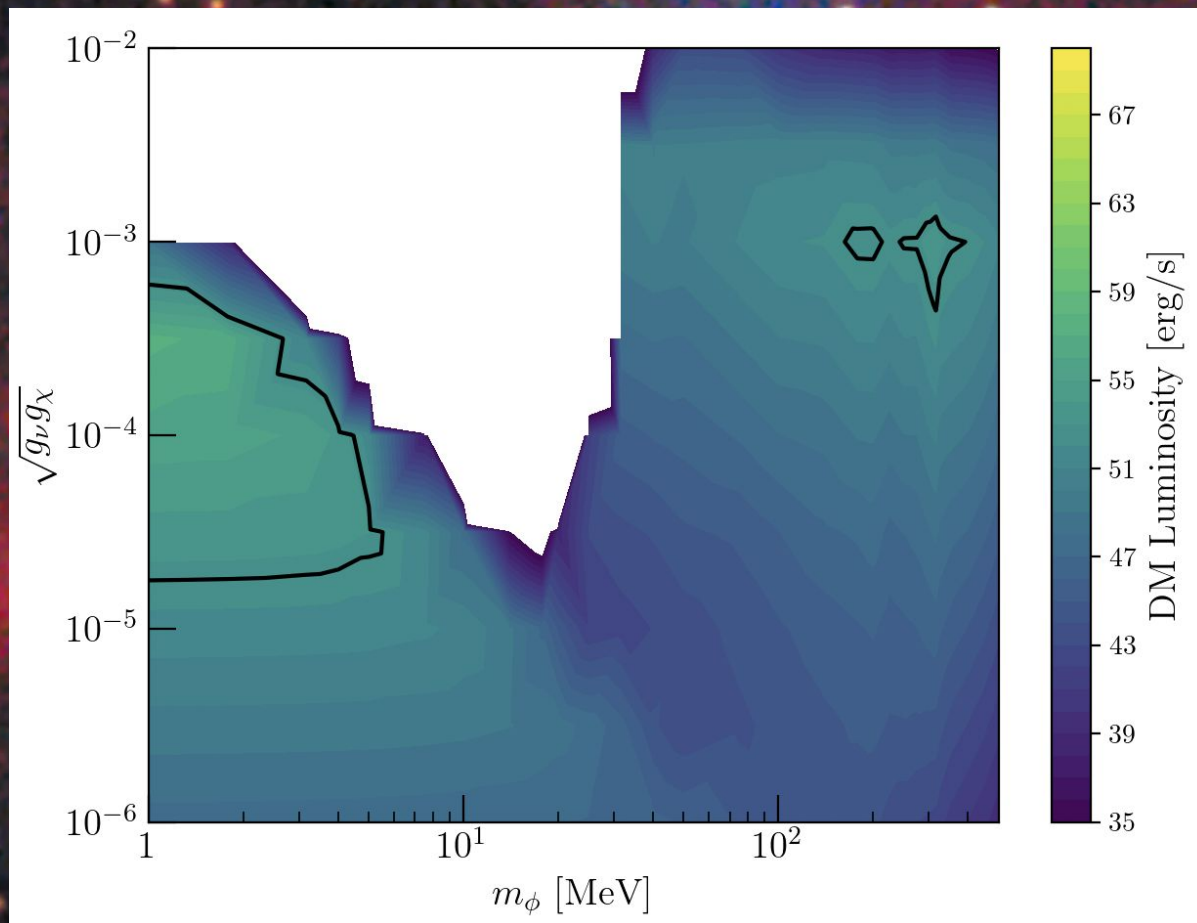


s-channel Models

$$\mathcal{L}_{\text{int}} \supset -g\phi^* \bar{\chi} P_L \nu + \text{h.c.}$$



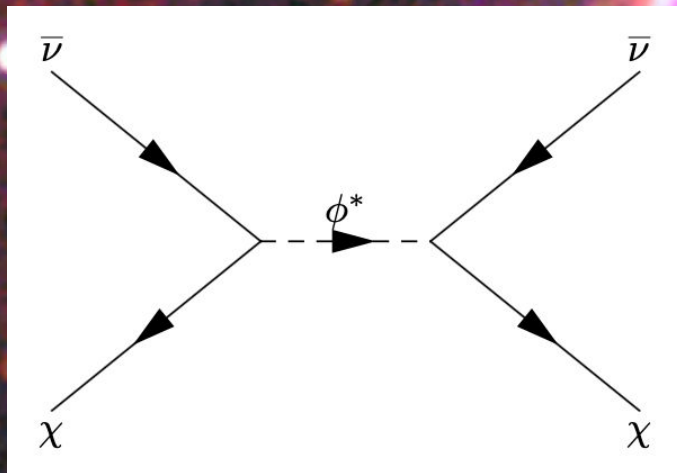
s-channel Models



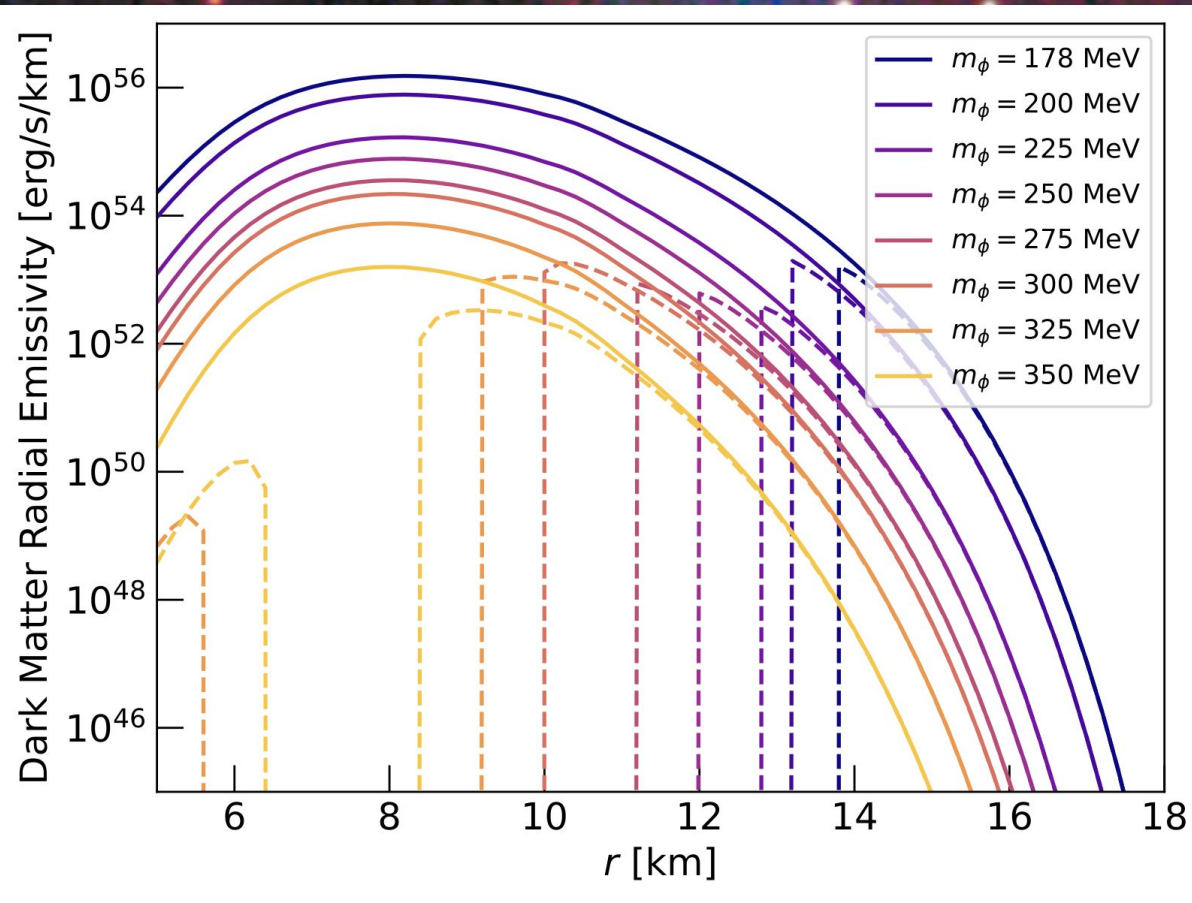
s-channel Models

$$\mathcal{L}_{\text{int}} \supset -g\phi^* \bar{\chi} P_L \nu + \text{h.c.}$$

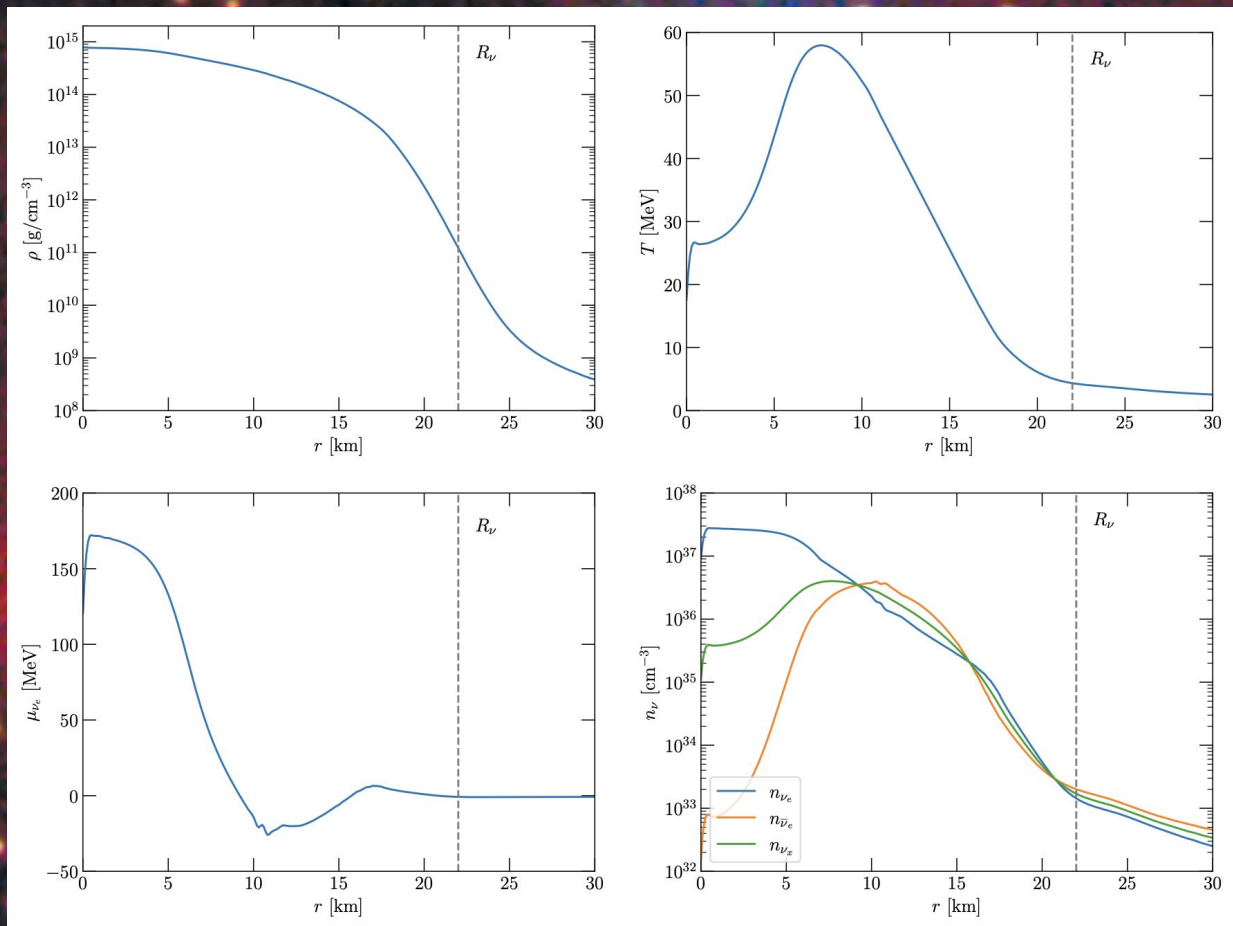
$$\frac{d\sigma_{\bar{\nu}\chi \rightarrow \bar{\nu}\chi}}{dt} = \frac{g^4}{32\pi(s - m_\chi^2)^2} \frac{(s - m_\chi^2)^2}{(s - m_\phi^2)^2 + m_\phi^2 \Gamma_\phi^2}$$



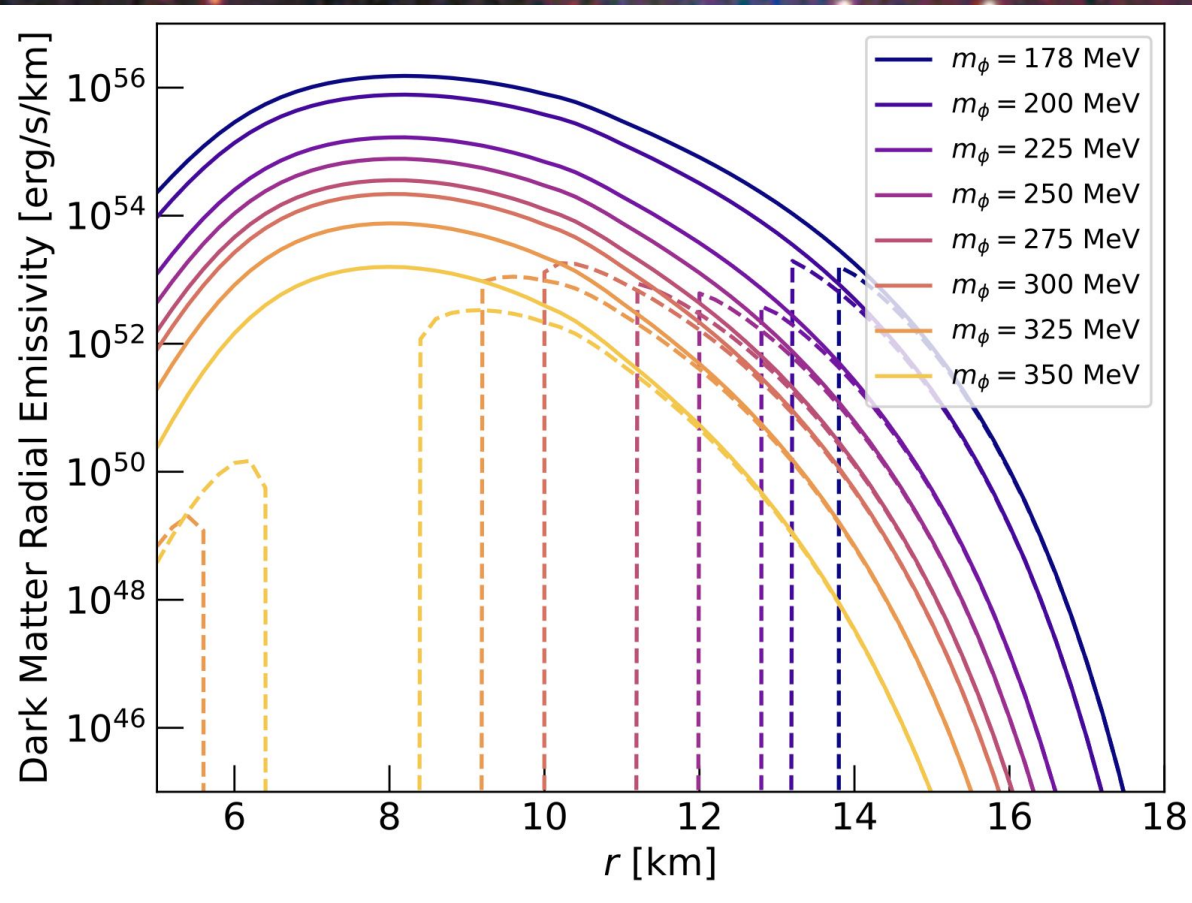
s-channel Models



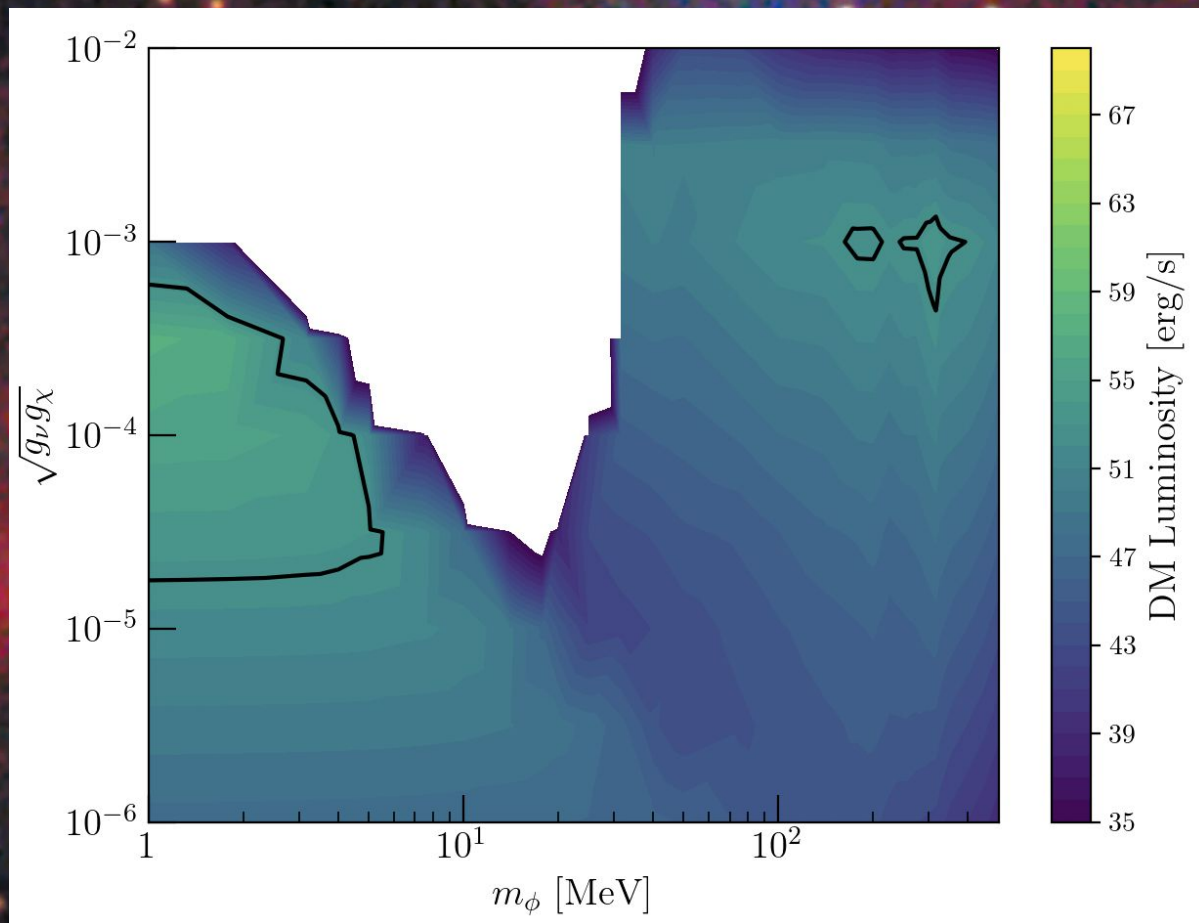
Constraints for a Heavy Mediator



s-channel Models



s-channel Models





Thank you!